

Validation, Vérification, Quantification des erreurs et des incertitudes

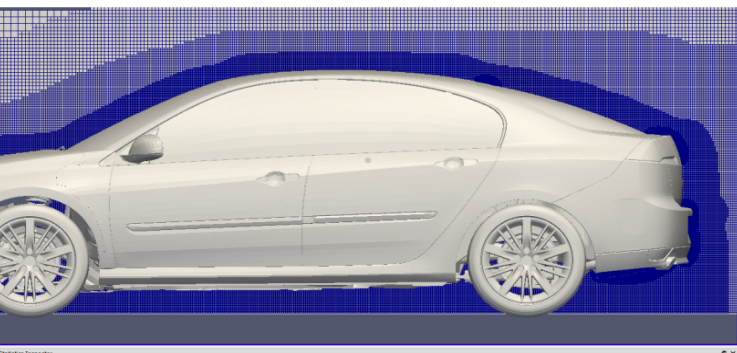
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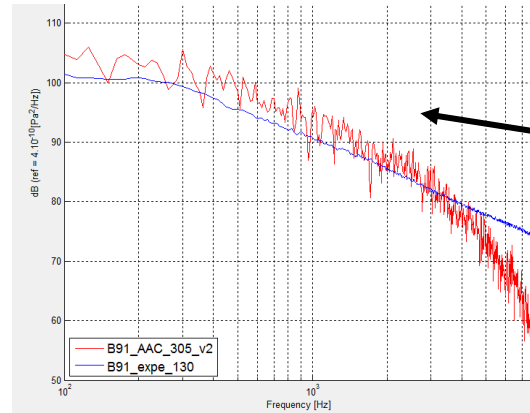
pierre.sagaut@univ-amu.fr

*Ecole Thématique de Simulation Numérique
23-27 avril, 2018*

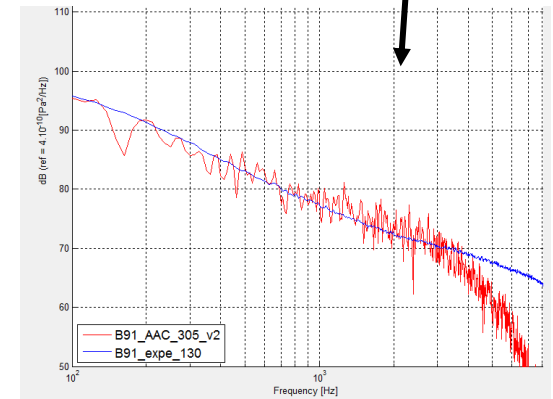
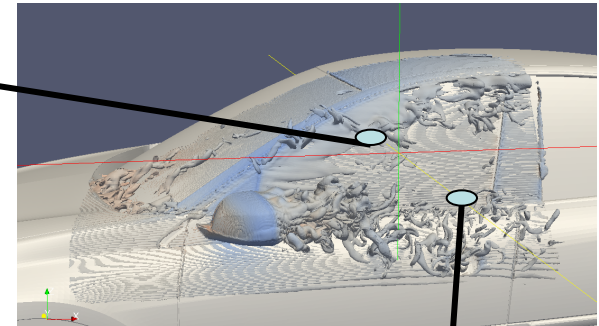


- ☐ Full scale vehicle simulation
- ☐ 186 surfaces (2,3 millions surface triangles)
- ☐ 10 levels of grid refinement, 88.6 millions cells
- ☐ $dx_{min}=1.25mm$
- ☐ 300 000 time steps $\rightarrow$ 0.96 sec
- ☐ $U_0 = 44.4$ m/s
- ☐ Wall Model in first cell LES
- ☐ LBM-ADM model

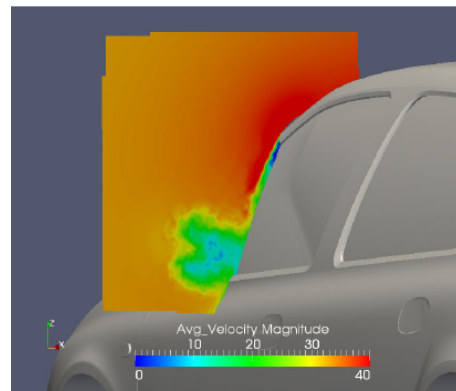
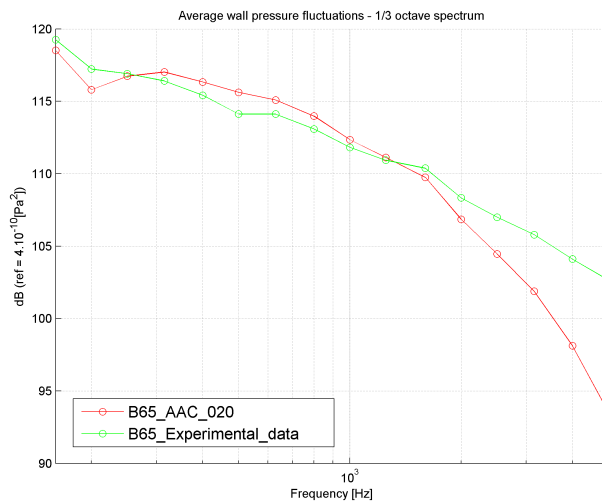
— LaBS
— Measurements

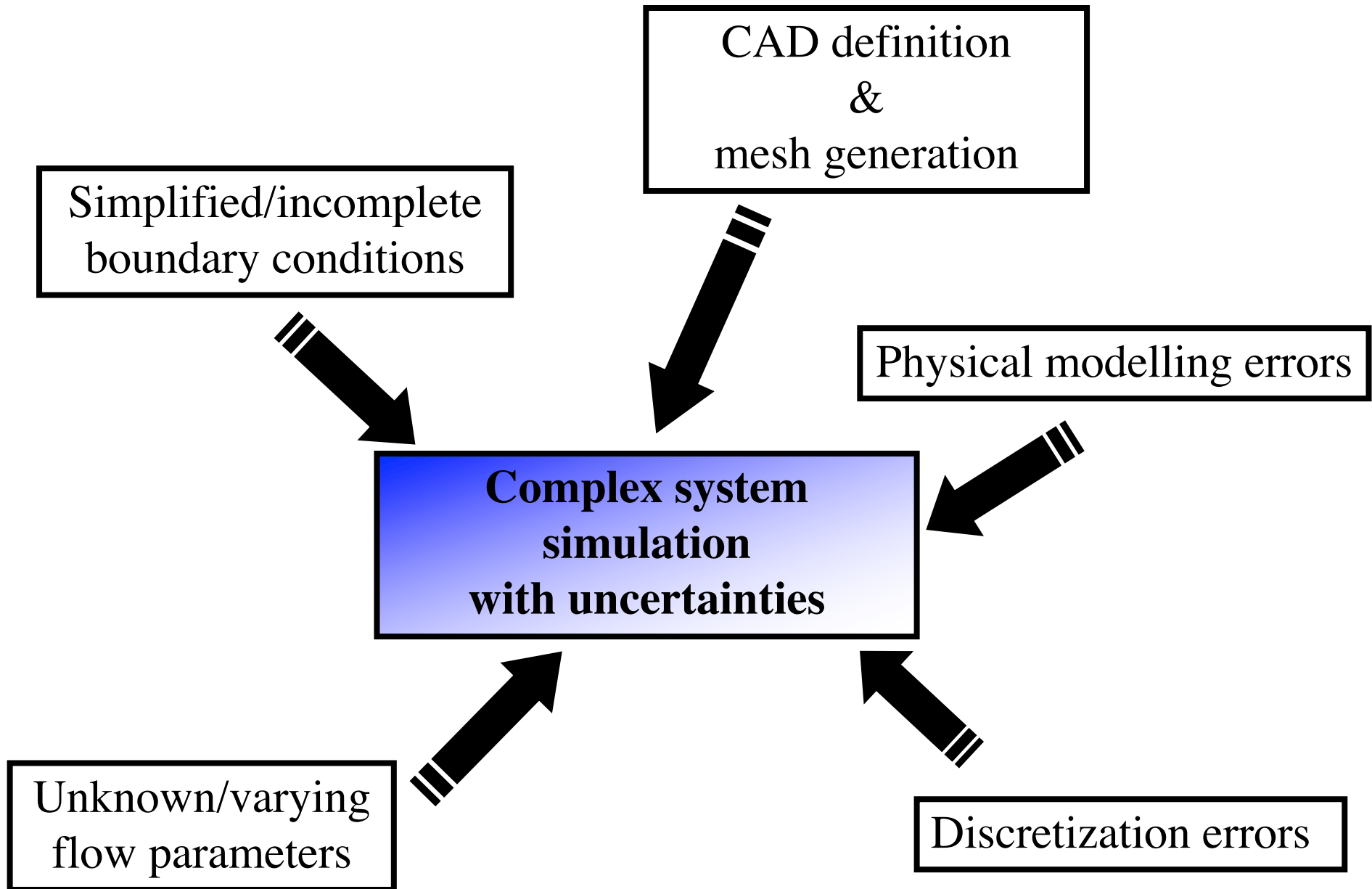


Laguna case : fine band spectra



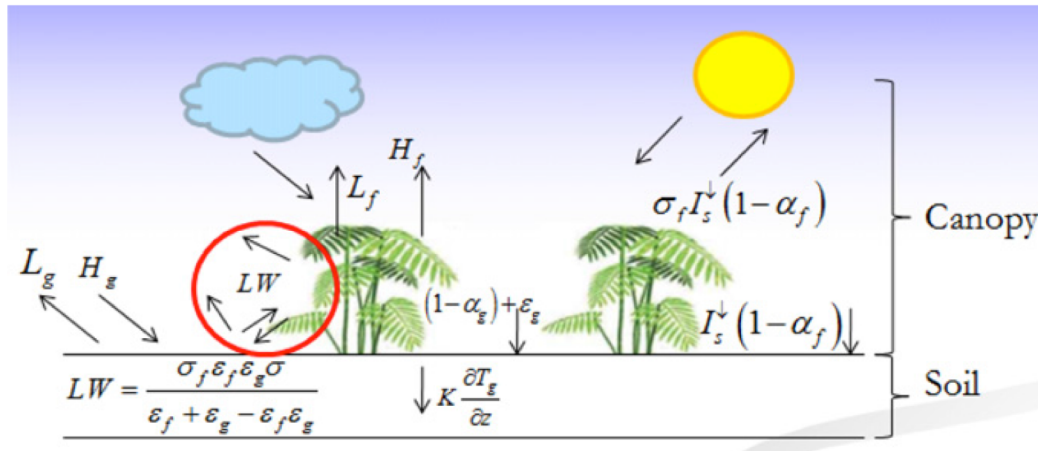
Clio case : third-octave band spectra, averaged on the whole surface of the side window





- **Aleatoric uncertainties** (or irreducible uncertainties):
 - describe physical variations caused by intrinsic randomness in the system and its environment
 - can be characterized in terms of probability distributions and covariance matrices
- **Epistemic uncertainty**:
 - caused by a certain lack of knowledge (structural uncertainty in the model form or insufficient measurement data to quantify the value of an input parameter)
 - not probabilistic in nature
 - can be better described using intervals
 - can be reduced either by increasing model fidelity or by performing additional experiments
- **Numerical errors**

Example: Green Roof Models



16 physical/free parameters

Fig. 1. Main heat flow on a green roof. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

Table 1

Different inputs data for the numerical model.

LAI	Leaf area index (m^2/m^2)
ϵ_f	Foliage emissivity
ϵ_g	Soil emissivity
α_f	Albedo for foliage
α_g	Albedo for soil
ρ_f	air density in the foliage (kg/m^3)
σ_f	Foliage density of (%)
T_a	Air temperature
RH	Air temperature
I_{ir}	Long-wave radiation (W/m^2)
I_s	Short-wave radiation (W/m^2)
W	Wind speed (m/s)
K	Thermal conductivity of the dry growing media ($\text{W}/\text{m}^{-1} \cdot \text{K}^{-1}$)
Δ	Soil thickness (m)
P_a	Atmospheric pressure (Pa)
P	Precipitation (mm)
D	Drained water (mm)

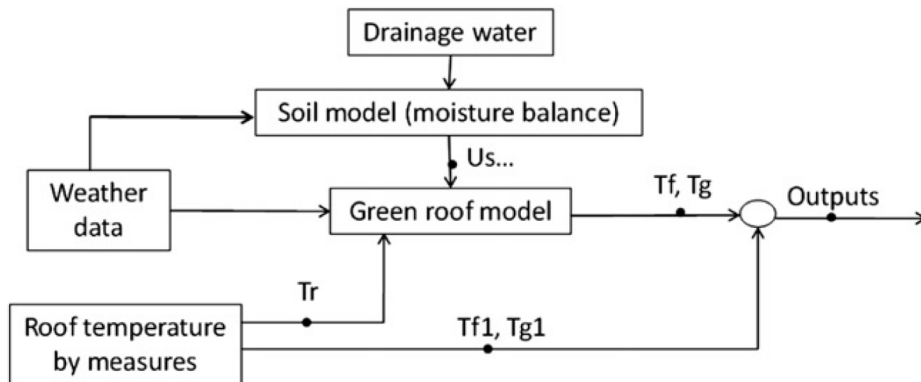
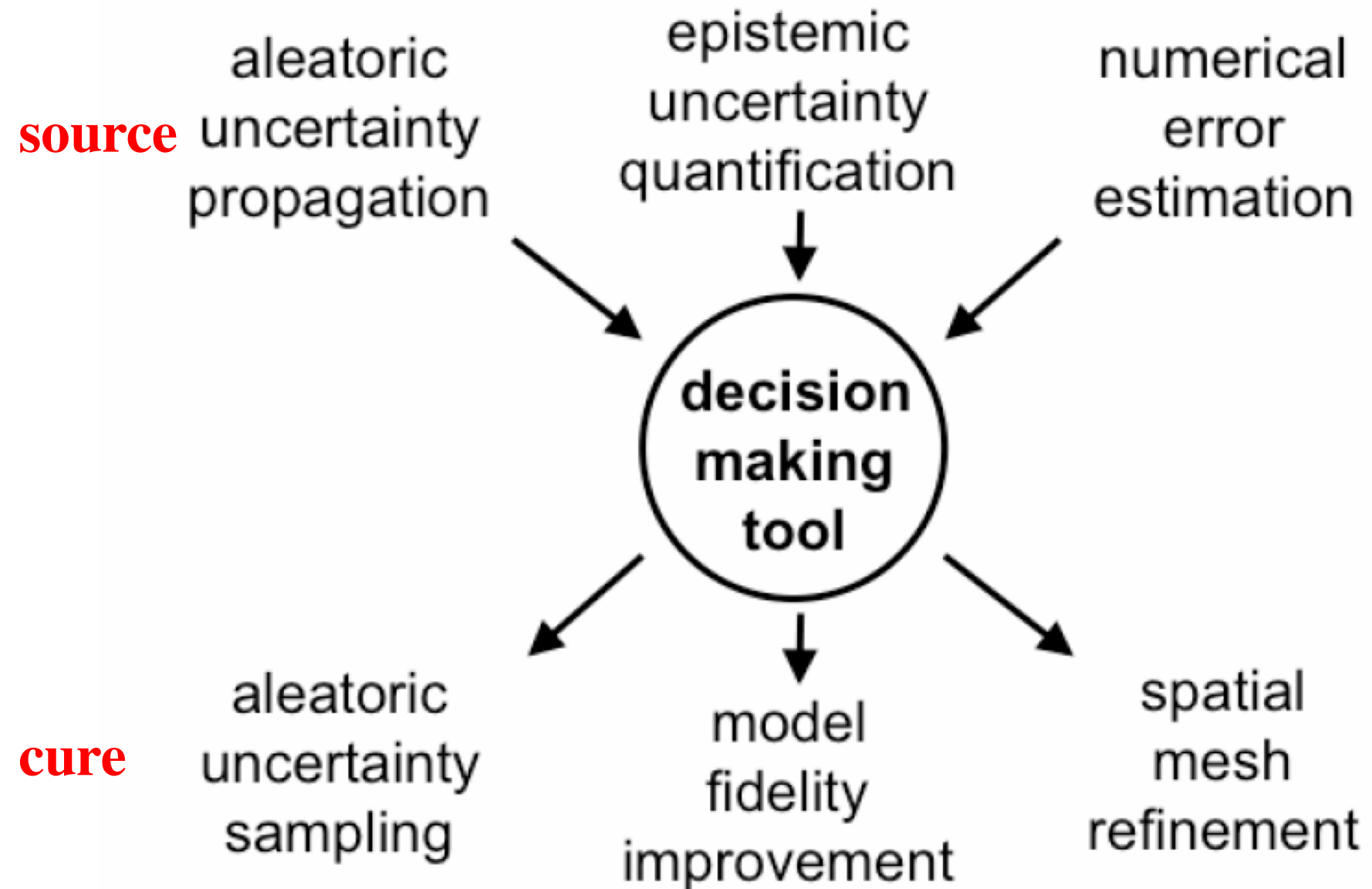
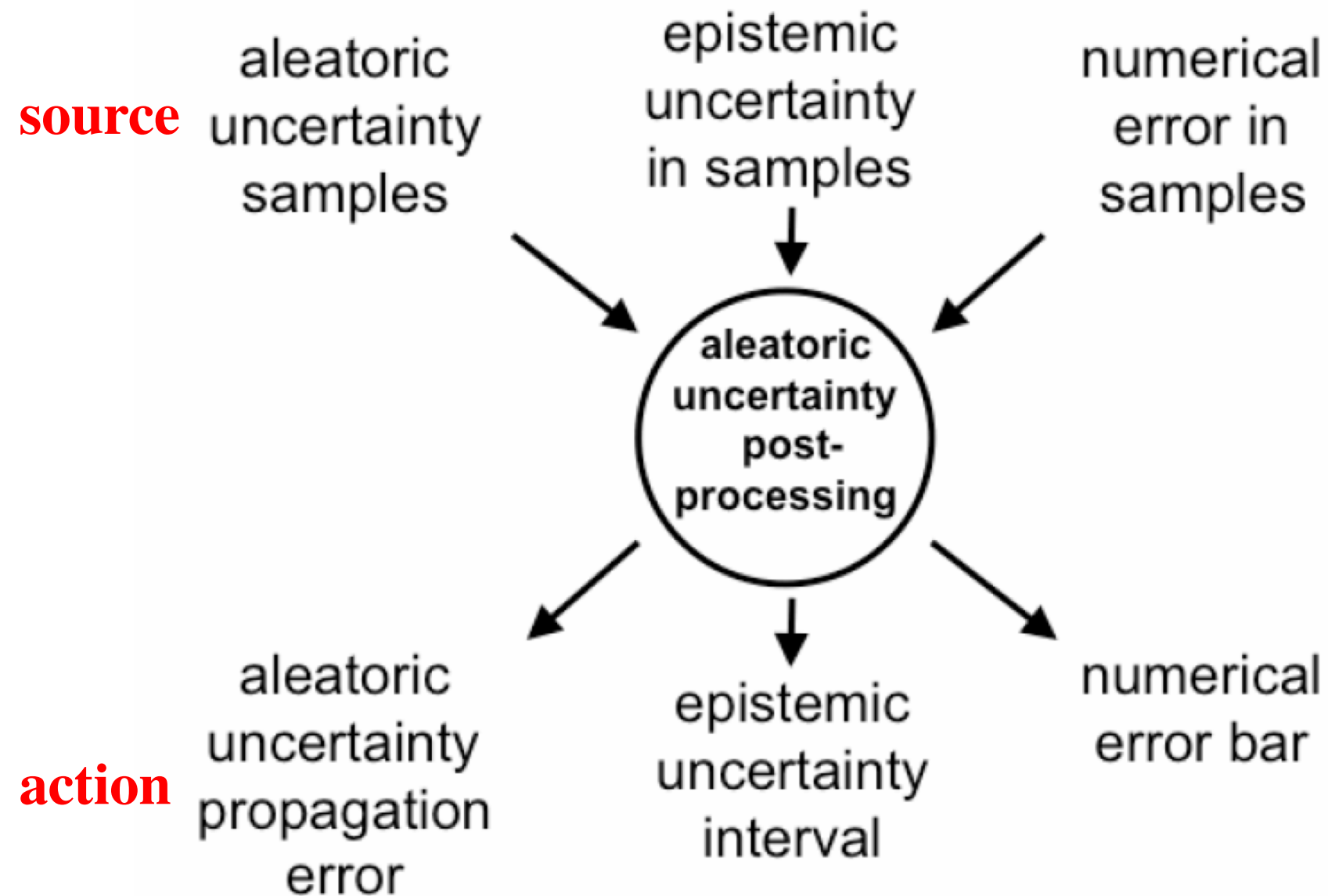


Fig. 5. The principle of the model validation.

How to increase fidelity ?



How to describe errors/uncertainties?



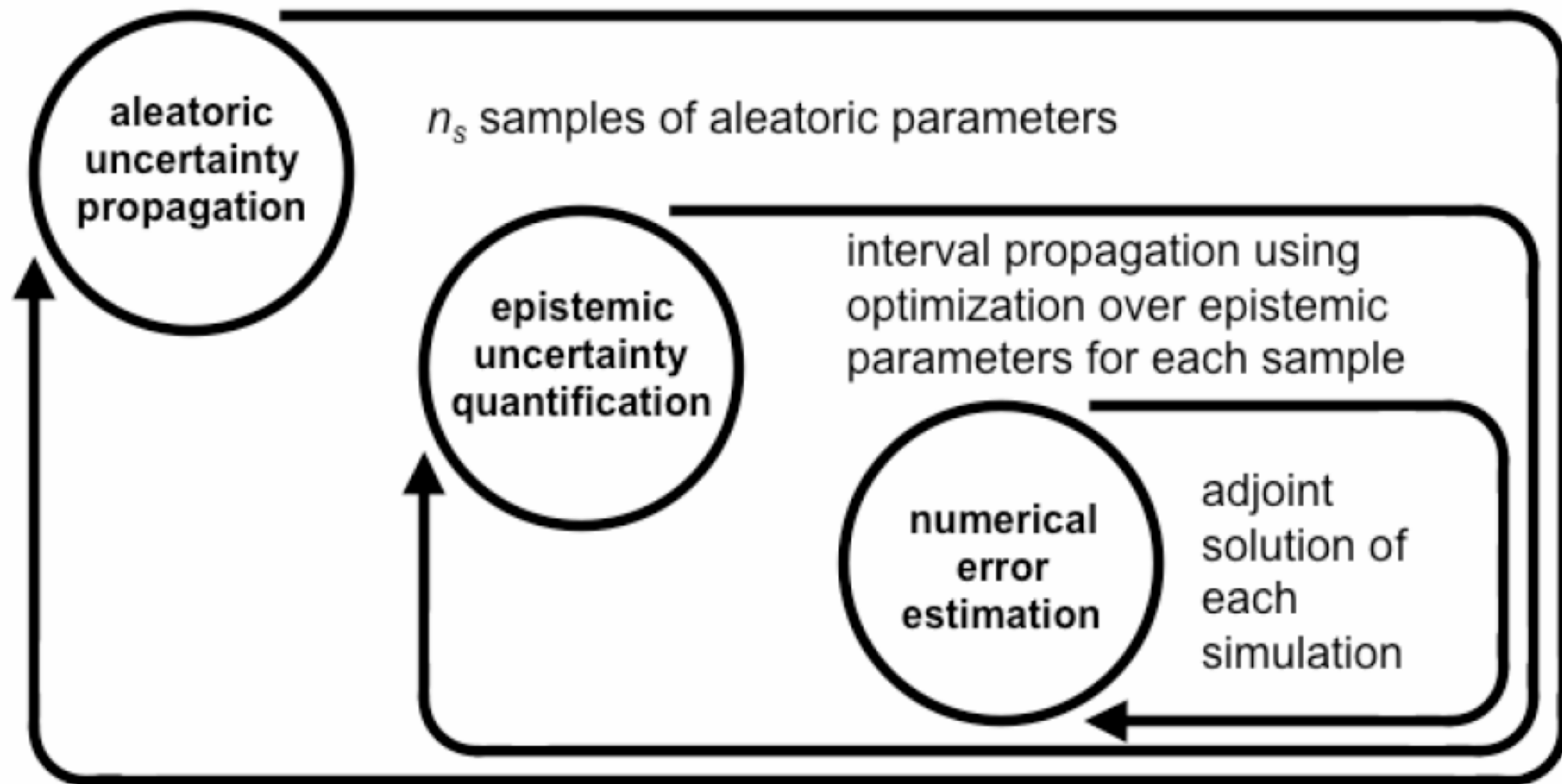
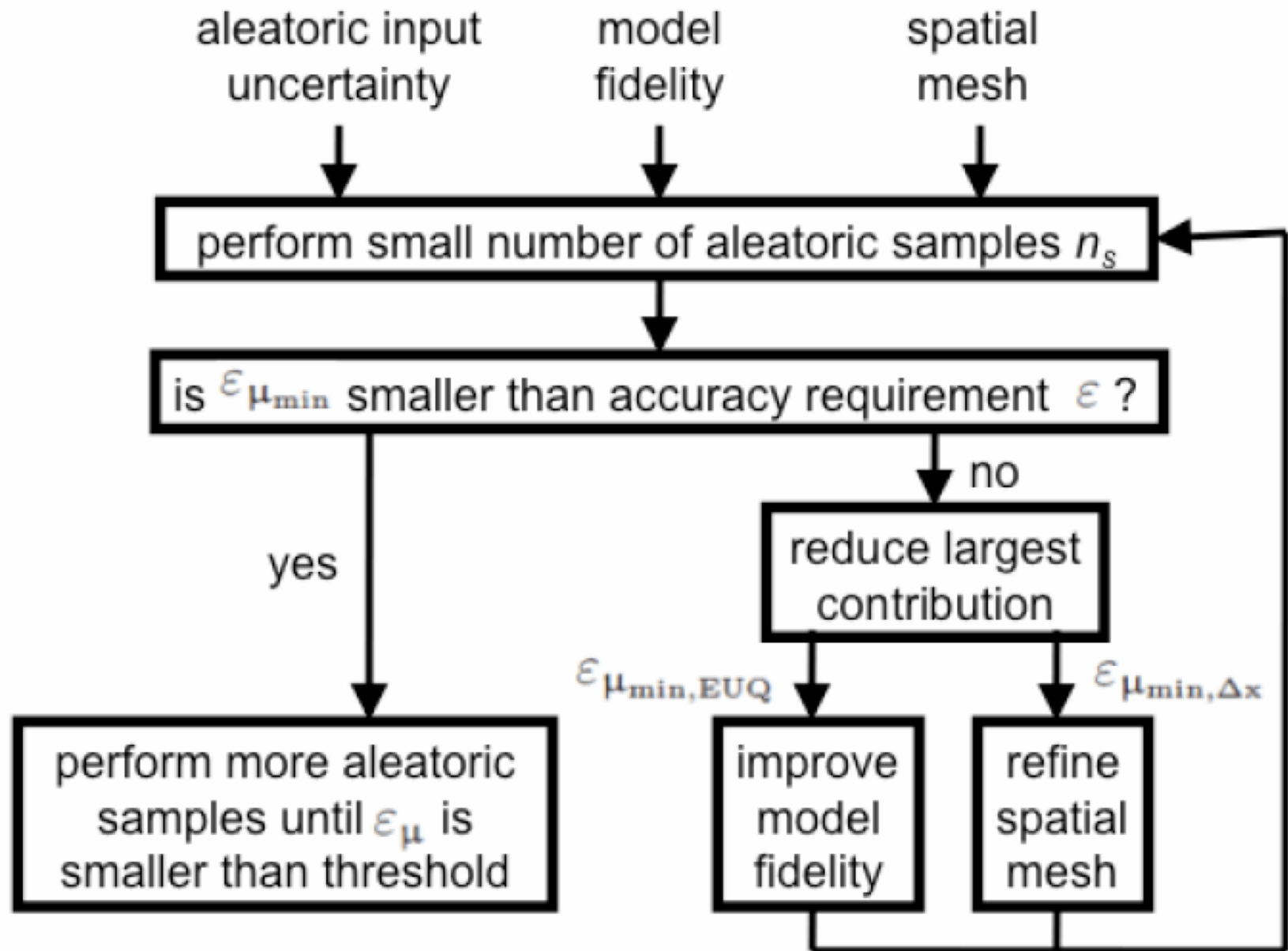


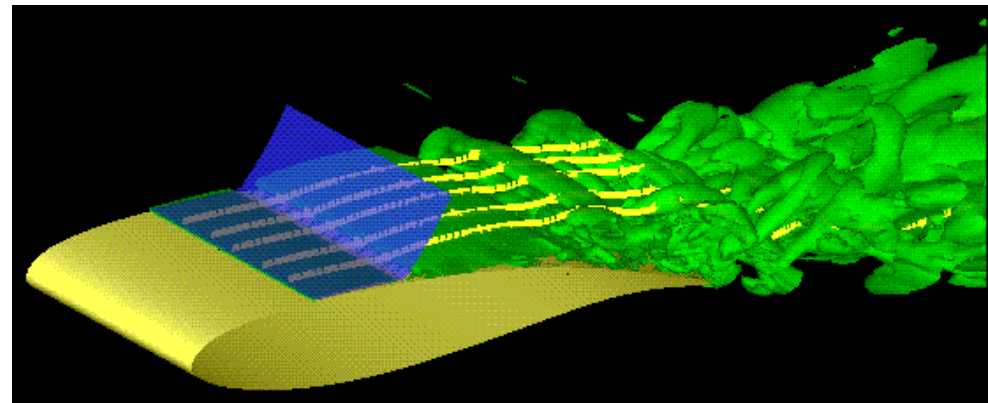
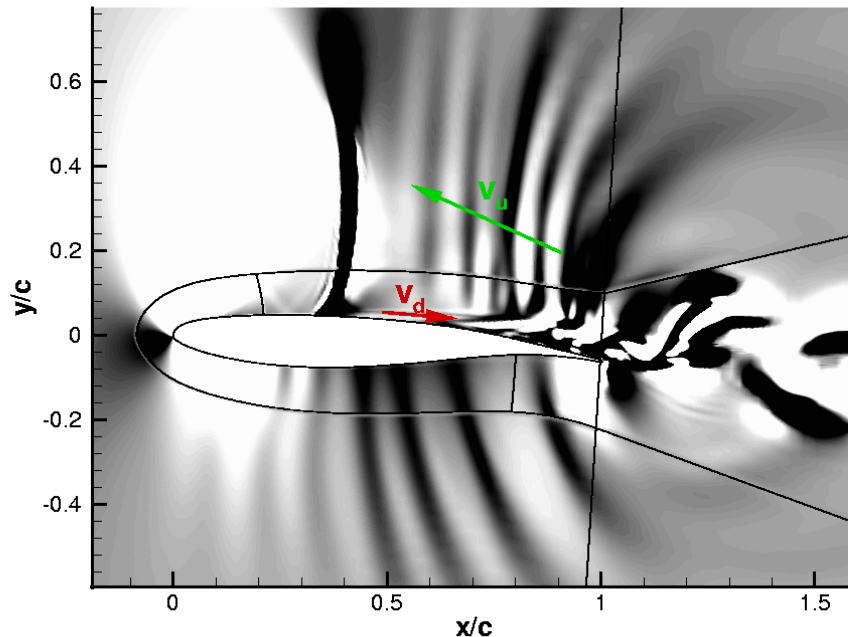
Figure 3. Nested uncertainty quantification and error estimation loop.



- ① Validation using experimental data
- ② Couplings between numerics and SGS models
- ③ LES numerics: beyond order of accuracy
- ④ Uncertainty quantification tools

**First issue: validation using
experimental data ?**

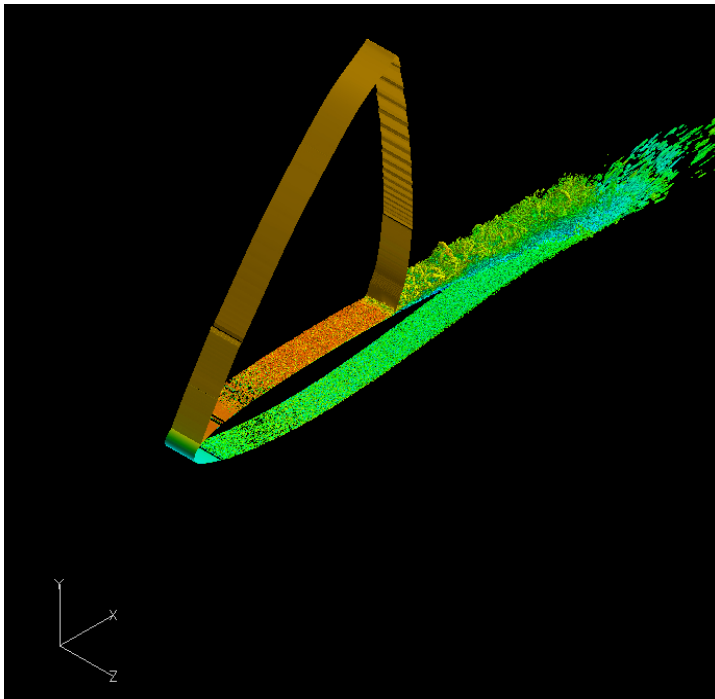
Buffeting on OAT15A airfoil (DES) (case with well-defined spectral peaks)



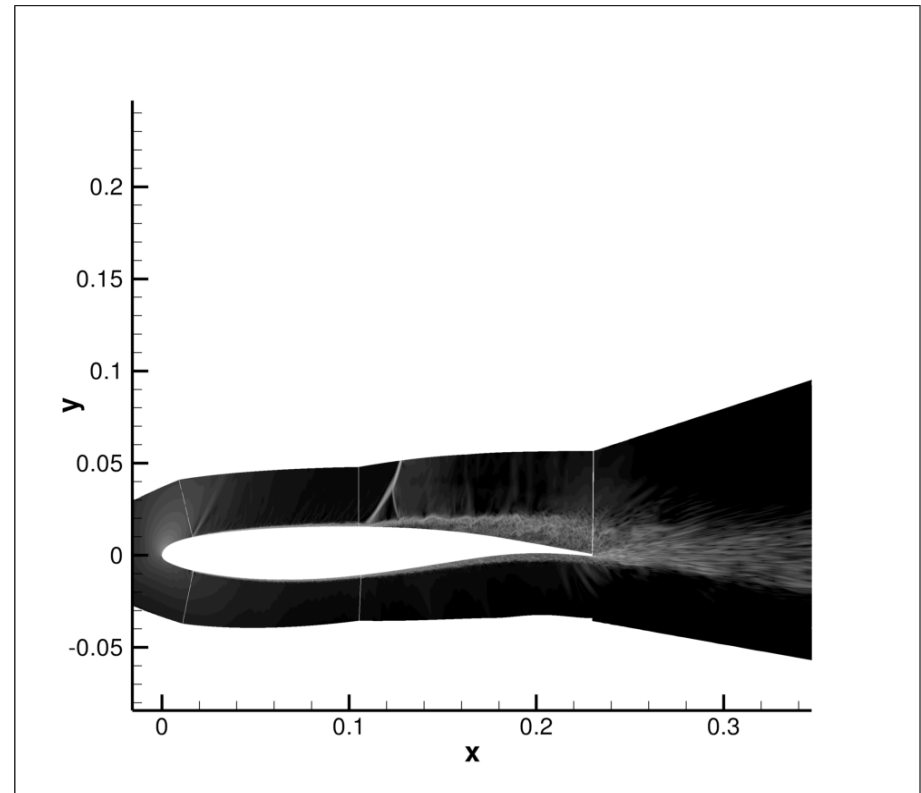
- Buffeting: large-amplitude pseudo-periodic shock oscillation
- Instability due to shock/boundary layer interaction
- High lift & Mach number

Instantaneous flow (grid C)

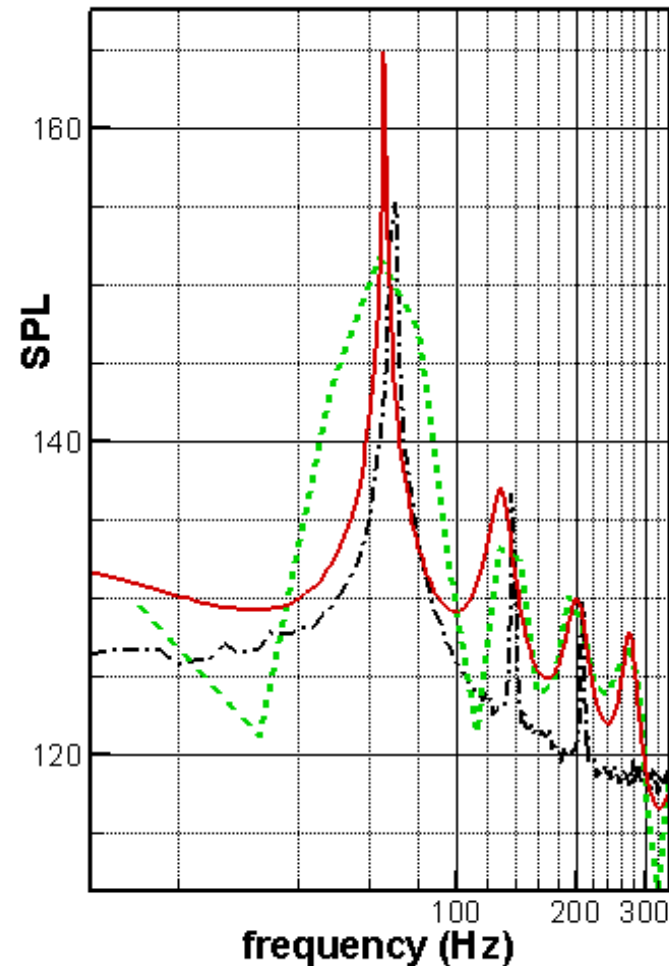
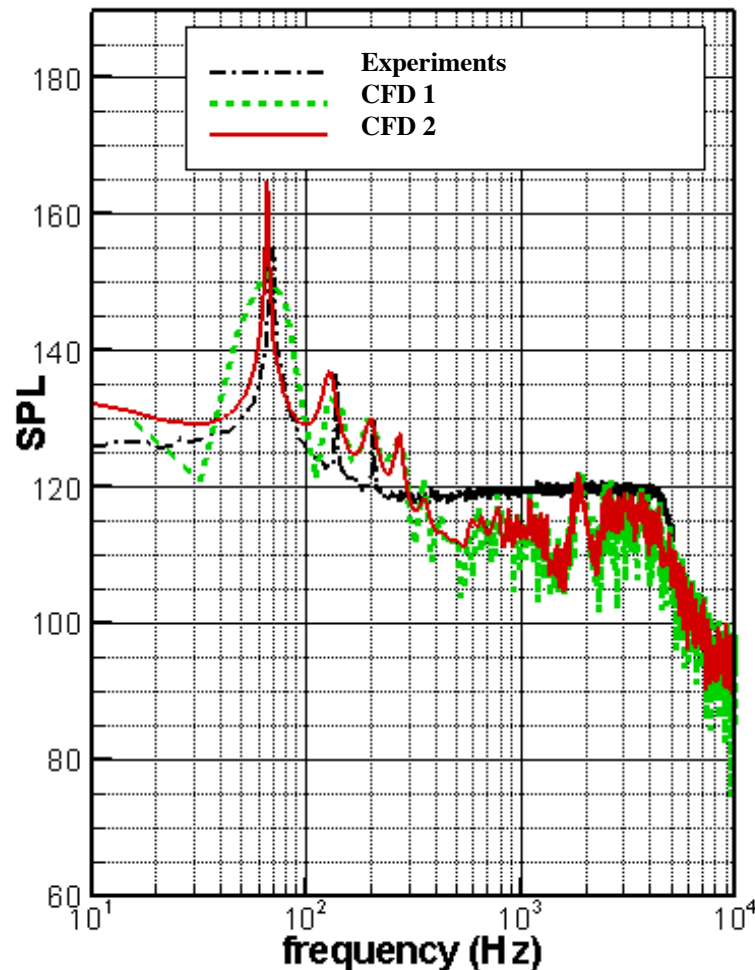
- 2D/3D interface (reflective) => numerically forced transition on the pressure side
- Periodic movement of the shock



Q criterion colored by the longitudinal velocity (grid C)



QUIZZ: which is the best CFD solution ?



Experimental vs. Numerical data

Buffeting on a 2D OAT15A airfoil

	Sampling time	Sampling frequency	Time step
Wind tunnel experiment	50 s	10 kHz	-
CFD (ZDES)	0.08s	100 kHz	$5 \cdot 10^{-7}$ s

Exact definition of **Power Spectral Density**

$$S_{xx}(f) = \lim_{T \rightarrow +\infty} \frac{1}{T} E [|X(f, T)|^2]$$

Fourier transform

$$X(f, T) = \int_0^T x(t) e^{-2i\pi ft} dt$$

A first approximation: **periodogram**

$$S_{xx}^{per}(f) = \frac{1}{T} |X(f, T)|^2$$

- Periodogram is easy to compute, but not consistent: variance of the coefficient does not decrease when sampling time is increased
⇒ an averaged variant must be used

Initial sample set split into L non-overlapping blocks with length M

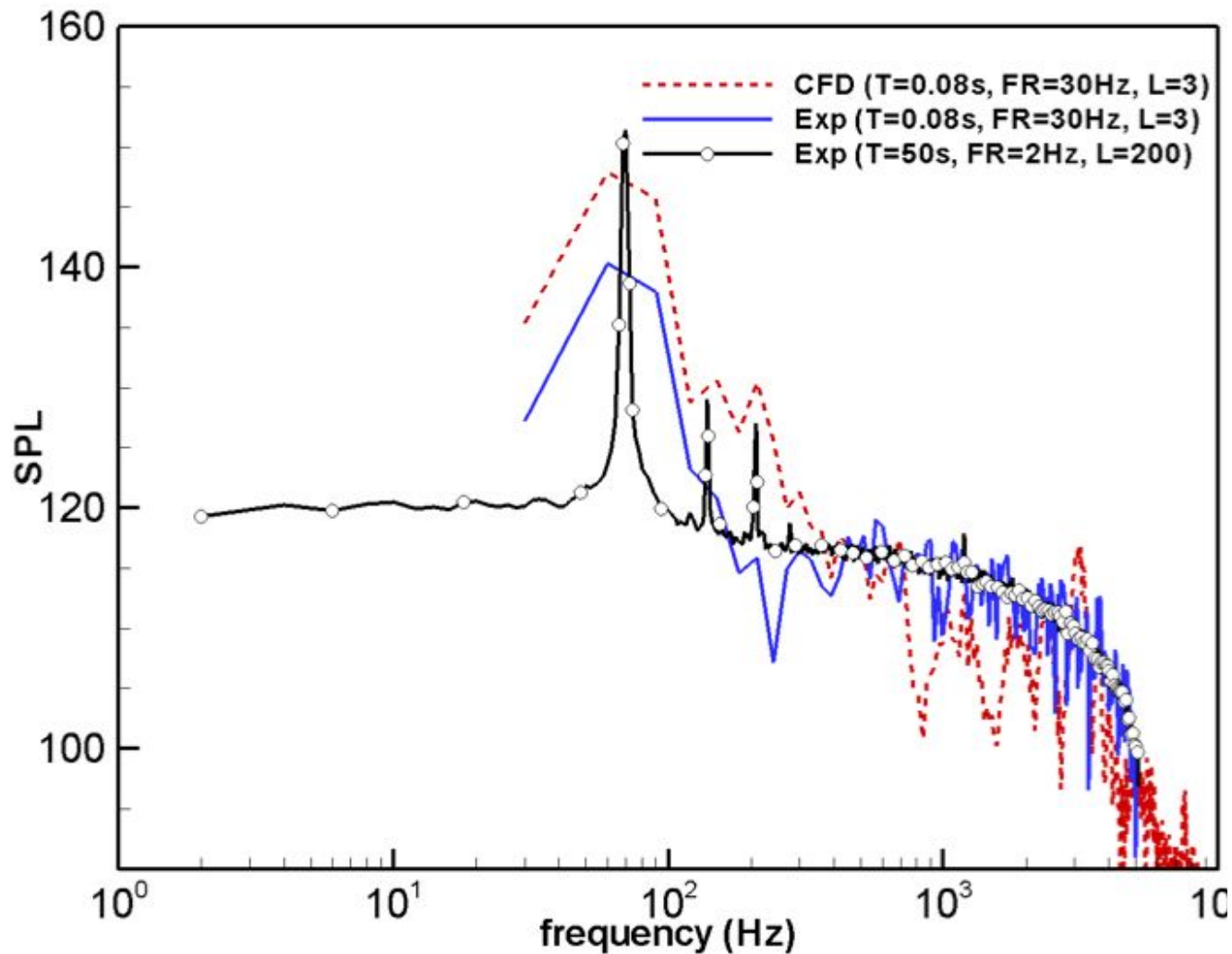
$$x_M^{[i]}[k] = x[iM + k] \quad 0 \leq i \leq L-1 \quad 0 \leq k \leq M-1$$

Averaged PSD estimate

Period of x(t)


$$S_{xx}^{aper}(f) = \frac{T_e}{L} \sum_{i=0}^{L-1} \frac{|X_M^{(i)}(f)|^2}{M}$$

Pressure PSD



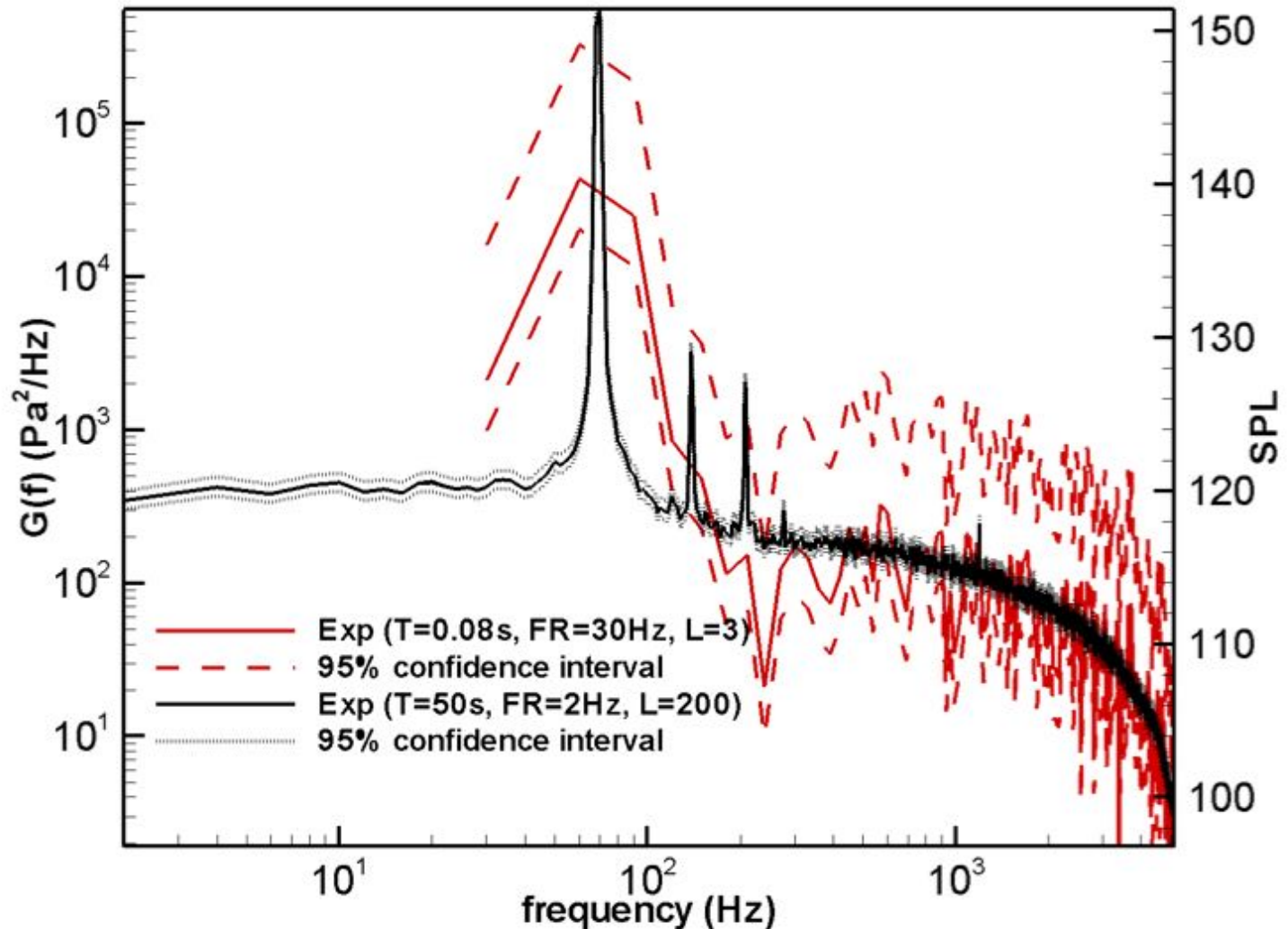
Is that a fair comparison ?

Kay's formula for confidence interval on PSD

- confidence interval $(1-\alpha) \times 100\%$
- L non-overlapping sample sets
- $\chi_n^2 = \chi^2$ distribution with n degrees of freedom
- $\alpha =$ confidence level such that $\text{Prob} [\chi_n^2 < \chi_{n;\alpha}^2] < \alpha$

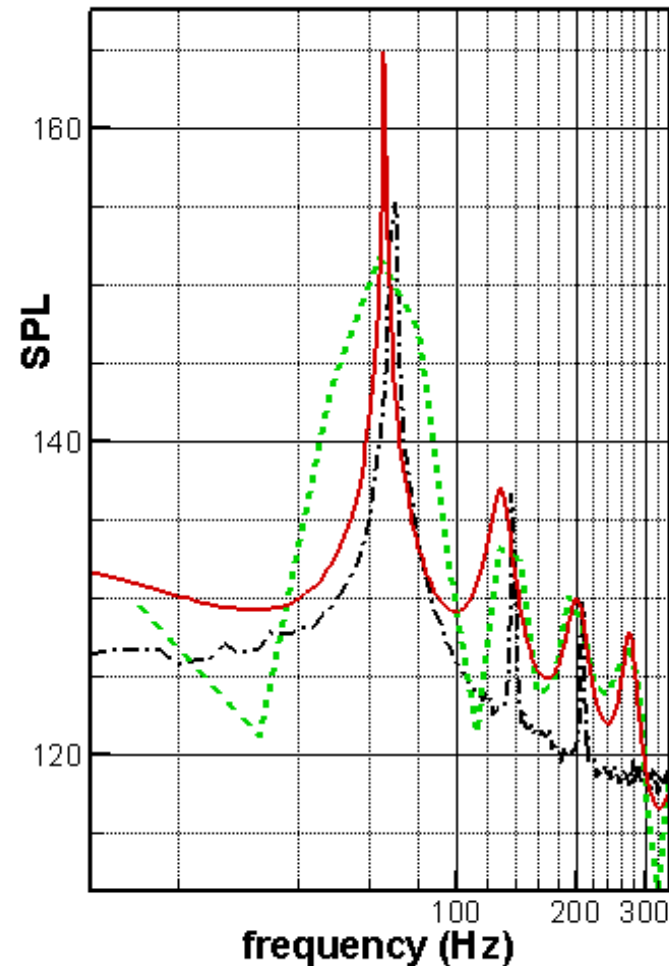
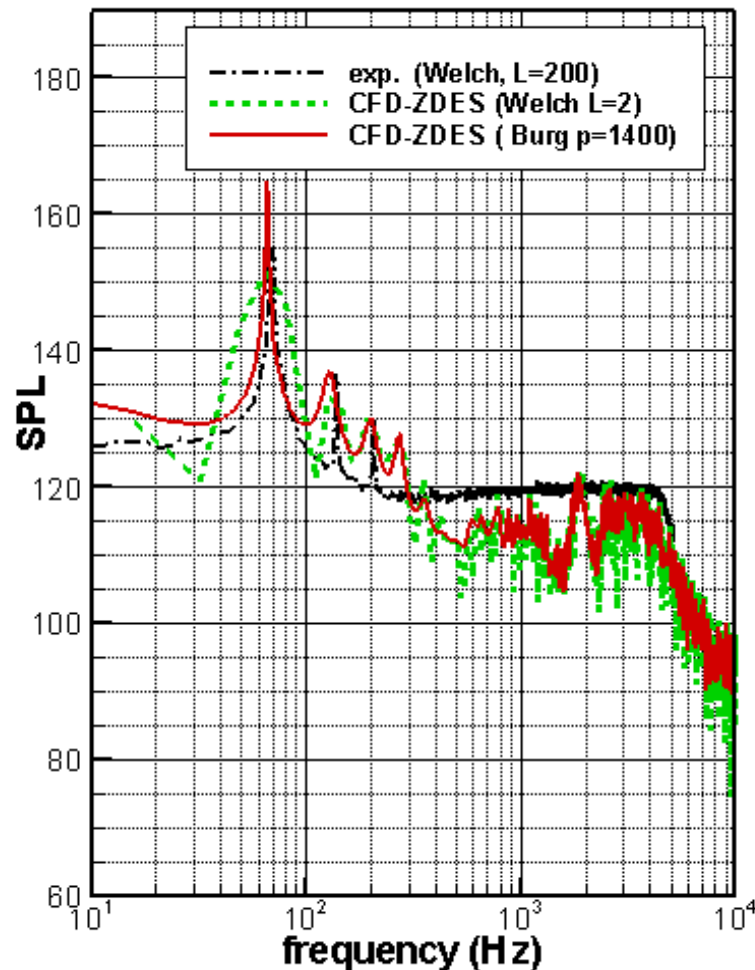
$$\left[\frac{2L.S_{xx}^{aper}(f)}{\chi_{2L;1-\alpha/2}^2}, \frac{2L.S_{xx}^{aper}(f)}{\chi_{2L;\alpha/2}^2} \right]$$

Pressure PSD - 95% confidence interval



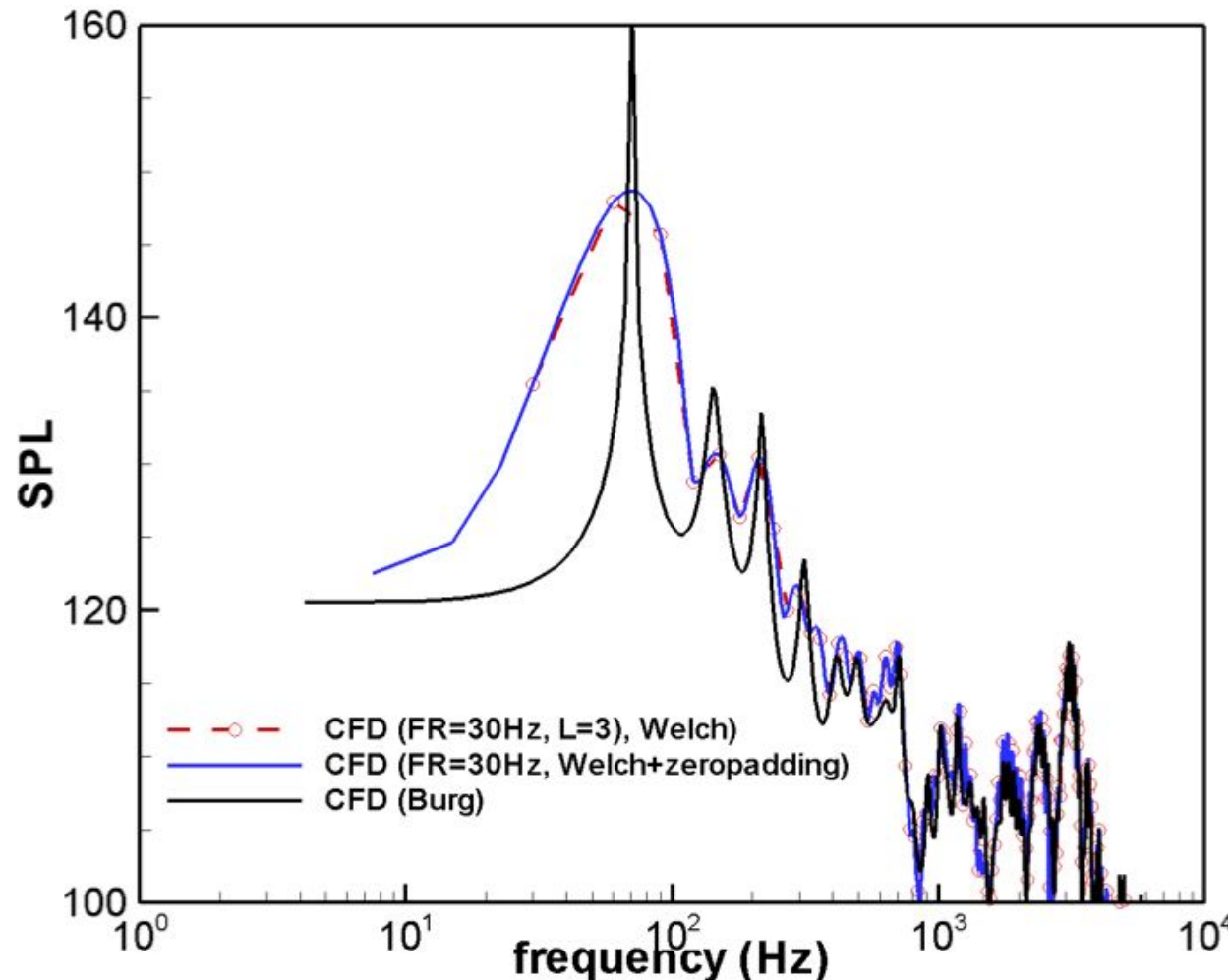
⇒ *confidence intervals should be always plotted on PSD !*

Pressure PSD

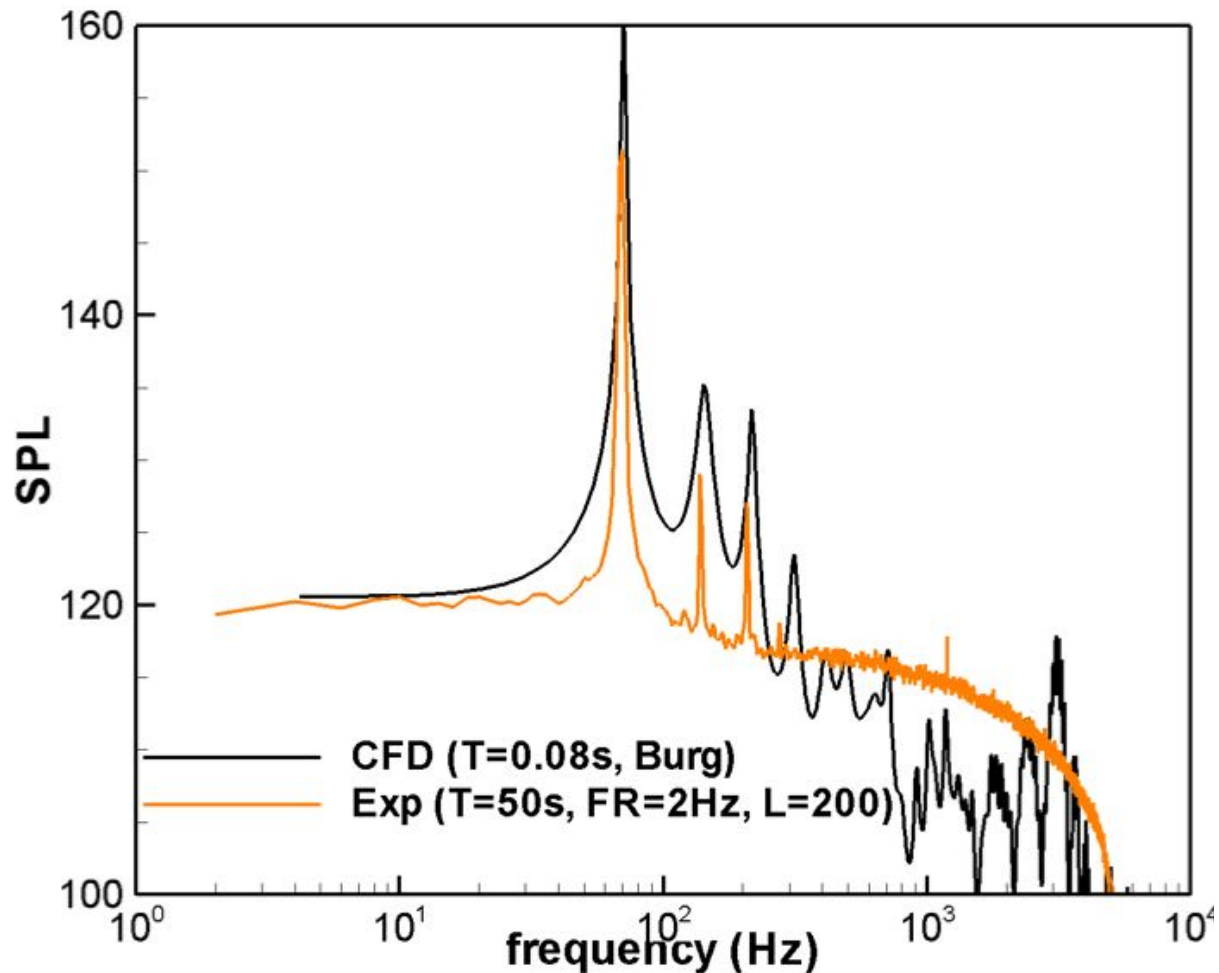


QUIZZ solution: there was only 1 CFD data set !

Pressure PSD - AR-Burg vs. Welch Sensitivity to details of Welch's method

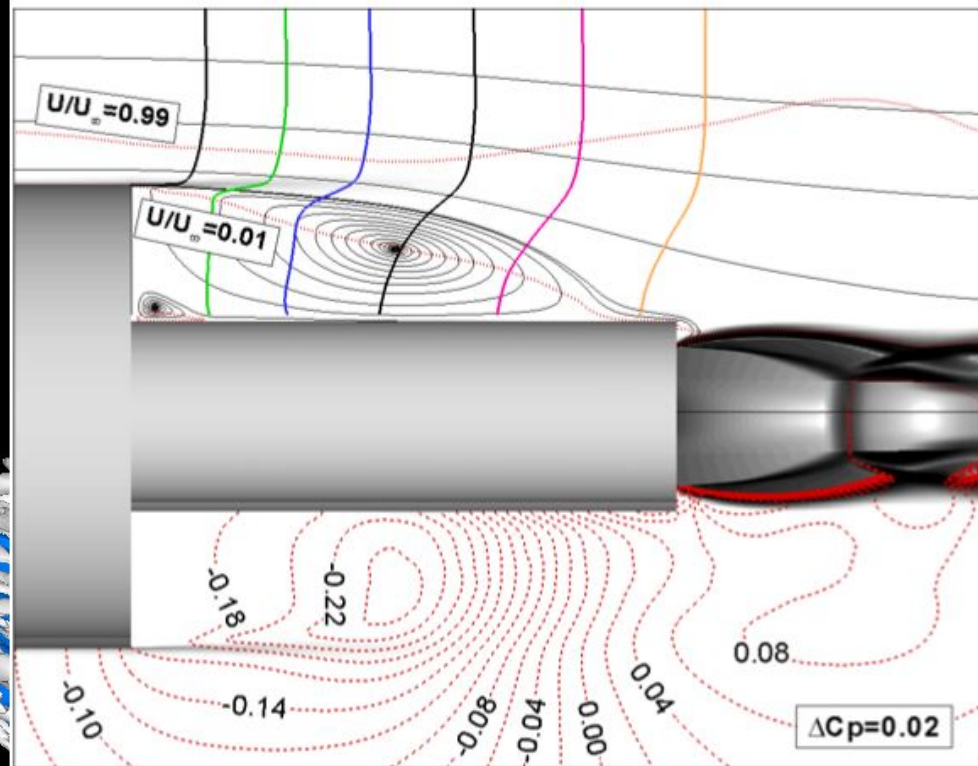
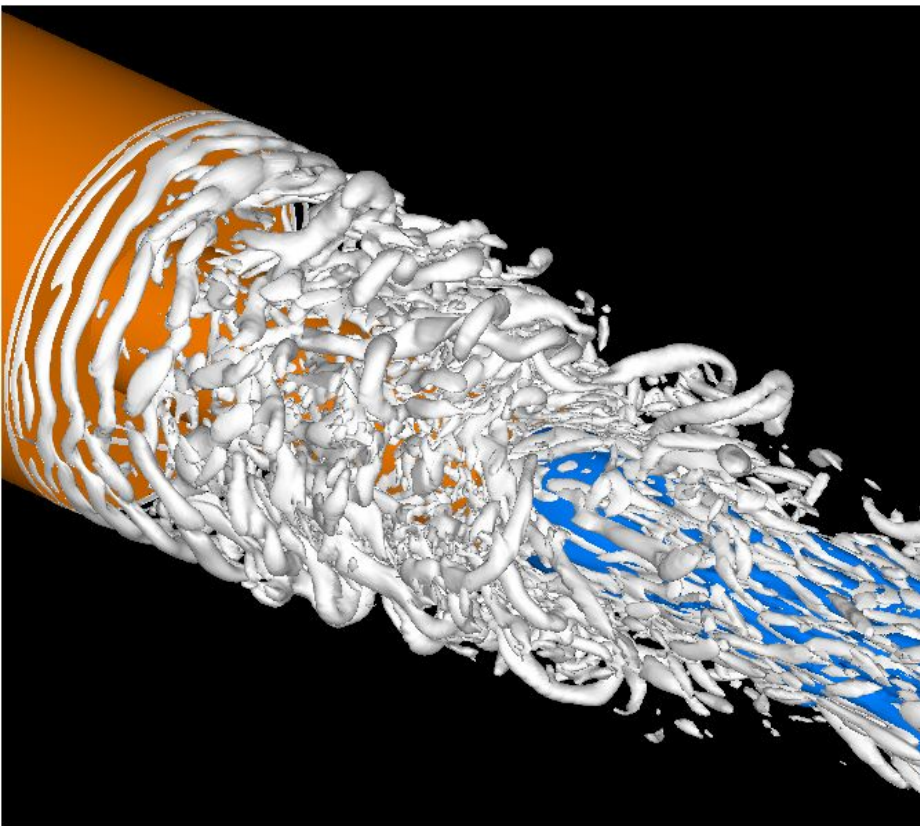


Pressure PSD - AR-Burg



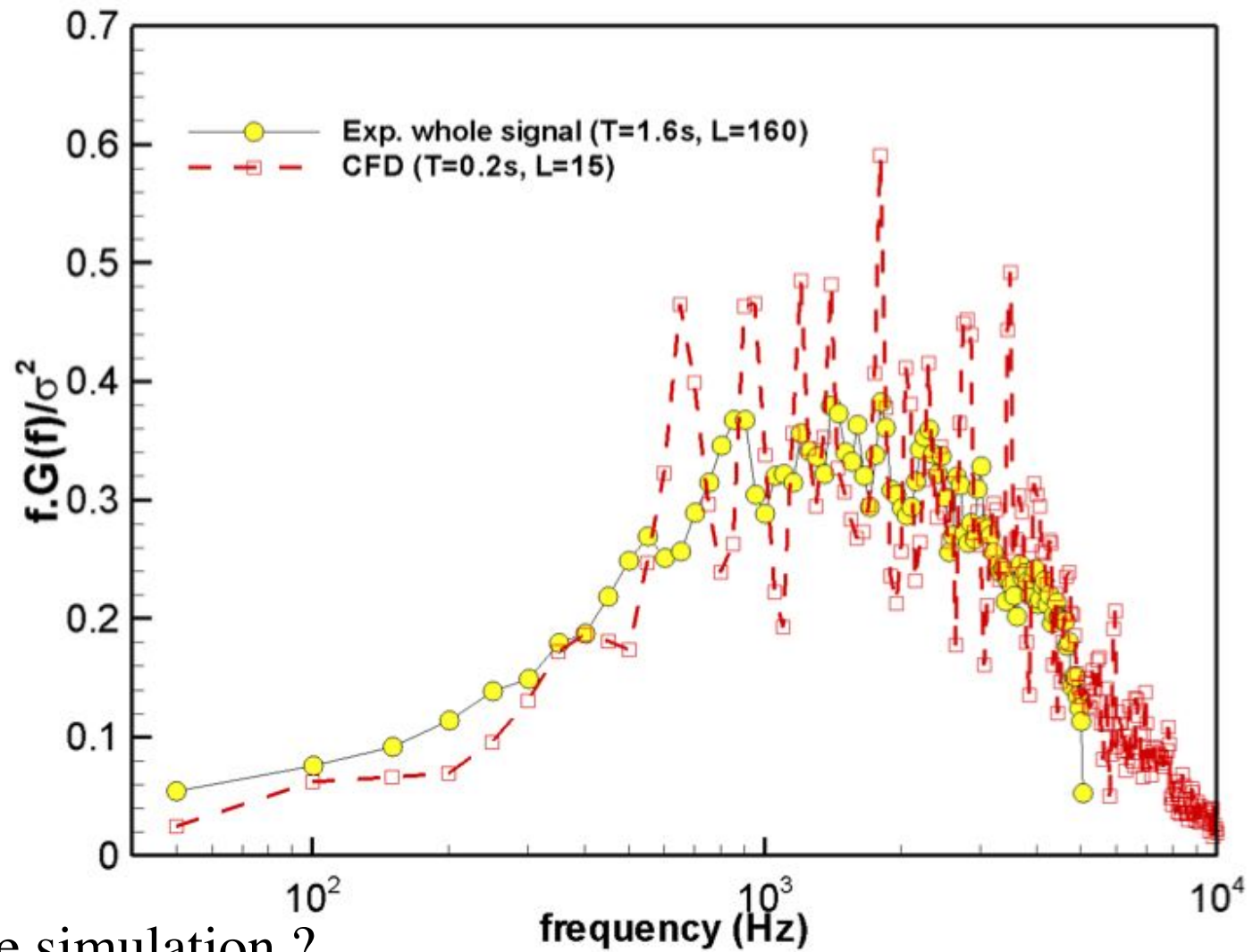
- Existence of peaks is accurately predicted
- not very efficient for spectral plateau capturing

Flow past a supersonic truncated contoured nozzle (case with broadband spectrum)



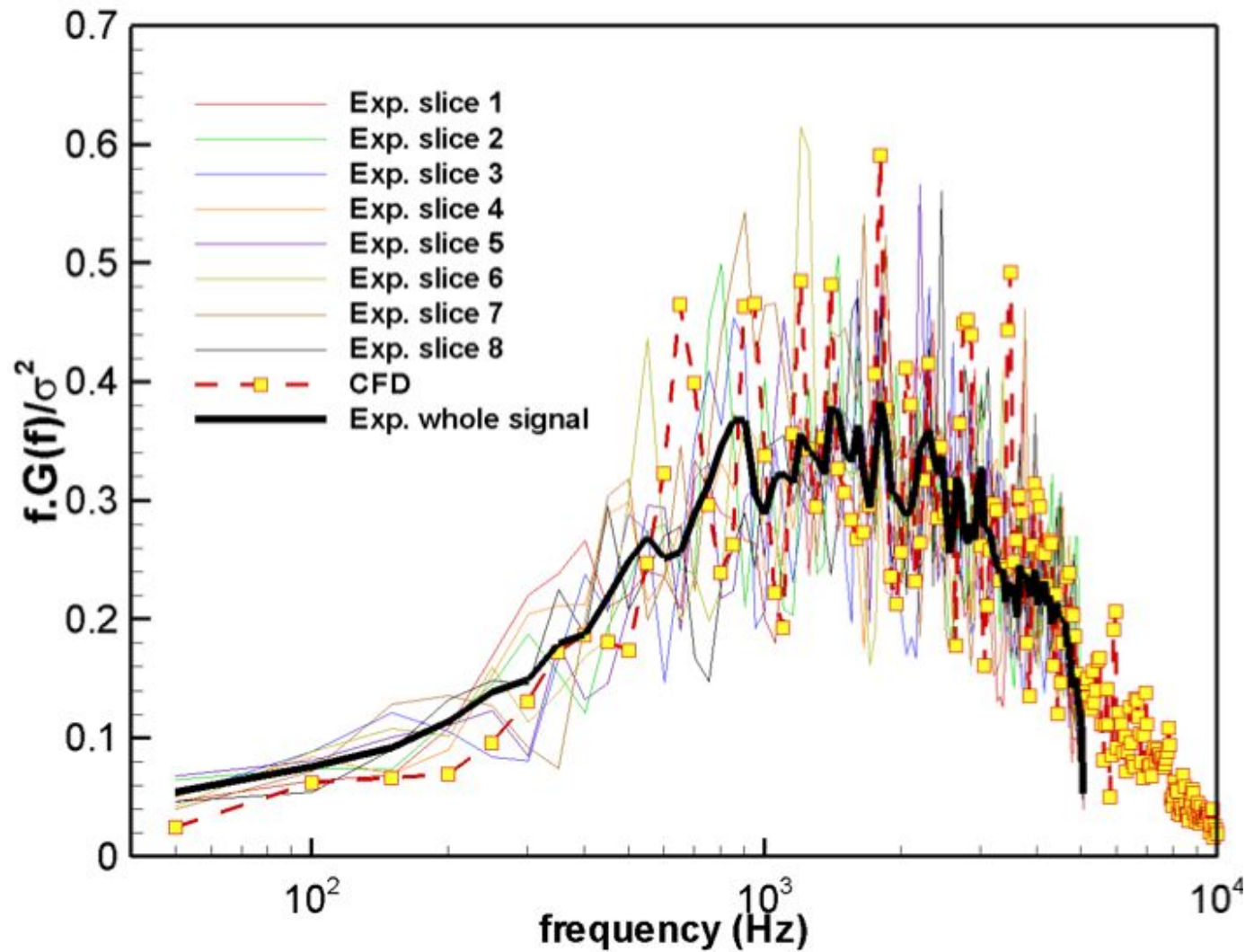
- experiments: sampling time = 1.6 s, sampling frequency = 10 kHz
- CFD (DES): sampling time = 0.2 s, sampling frequency = 100 kHz

Pressure PSD - Welch



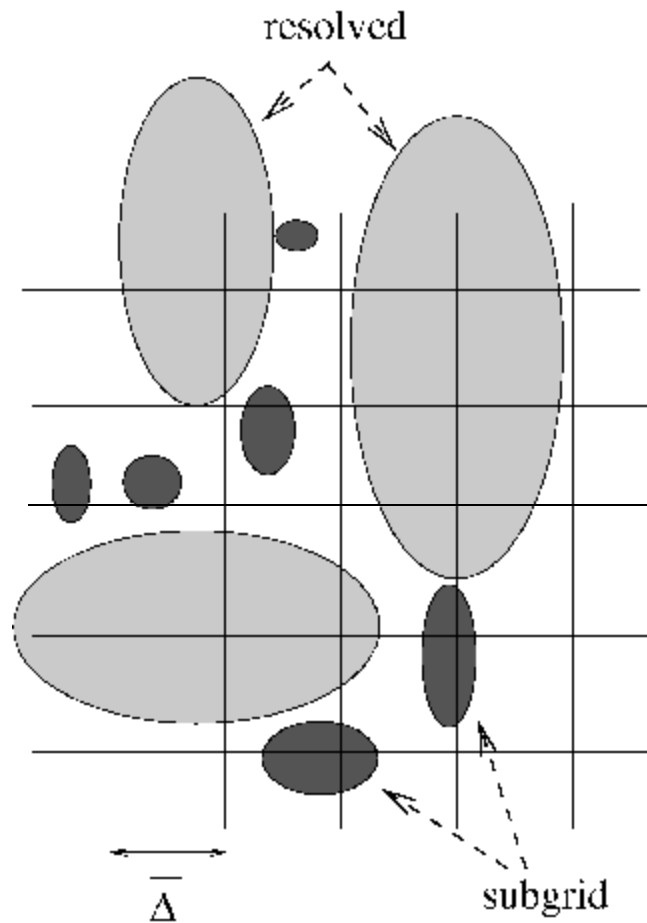
Inaccurate simulation ?

Pressure PSD - Welch

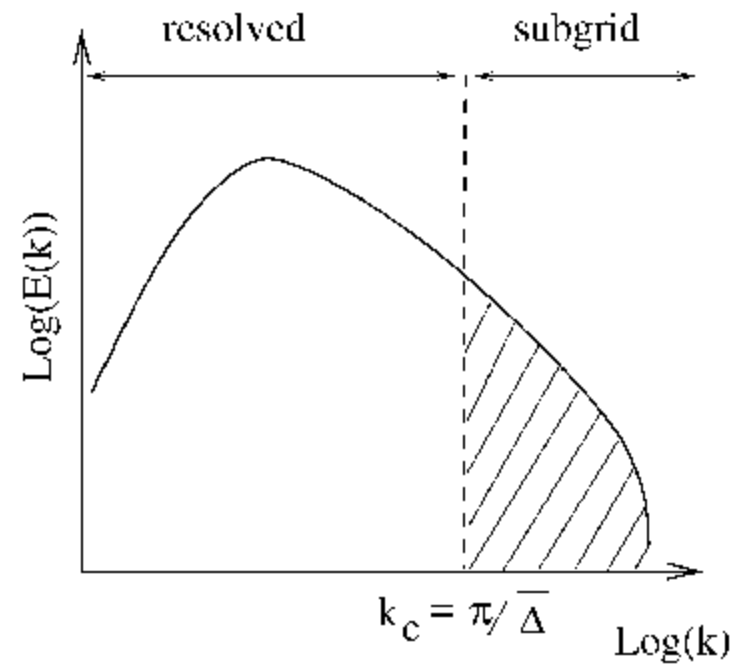


Introduction to numerical errors

LES: the key observation



PHYSICAL SPACE

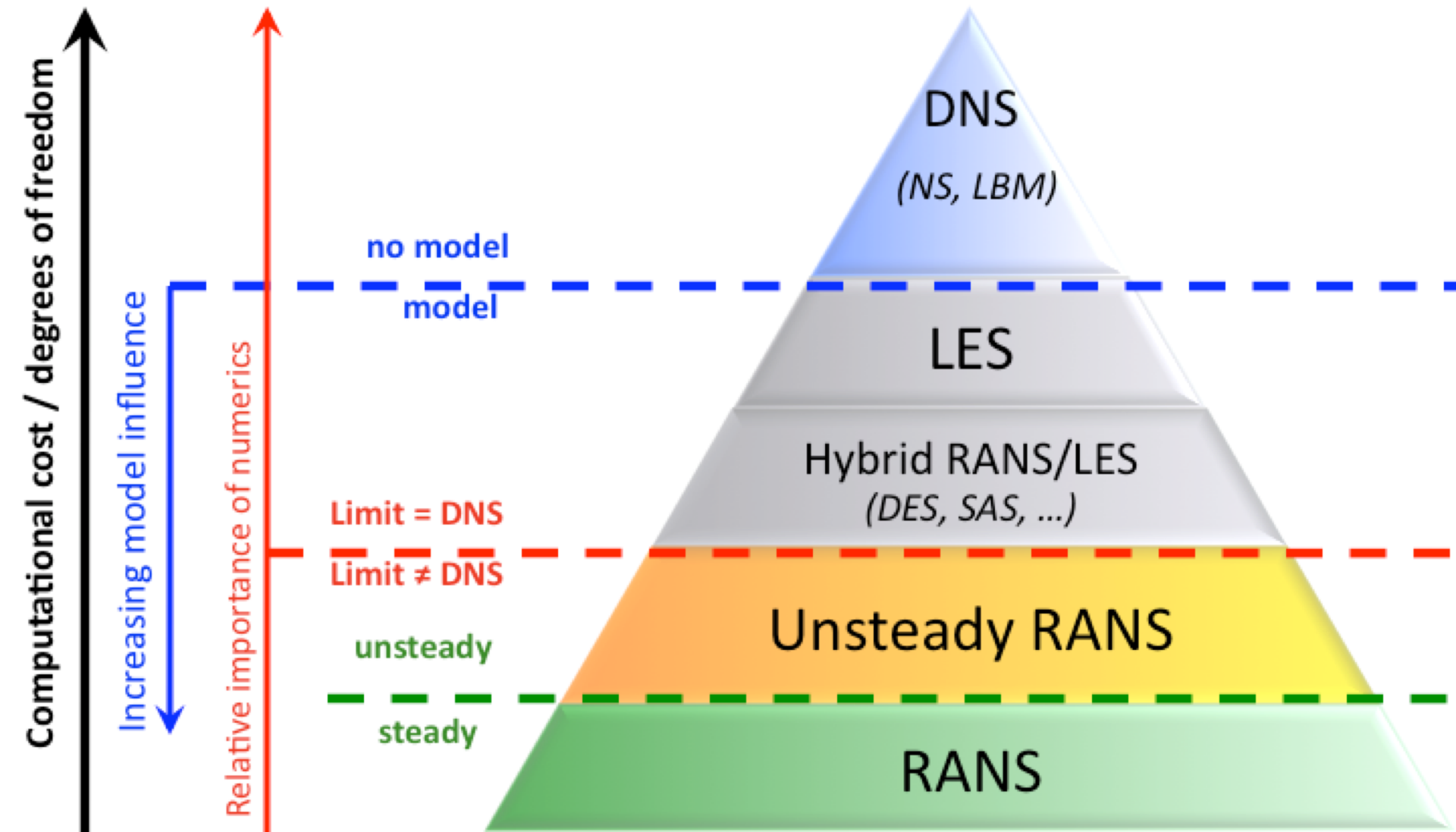


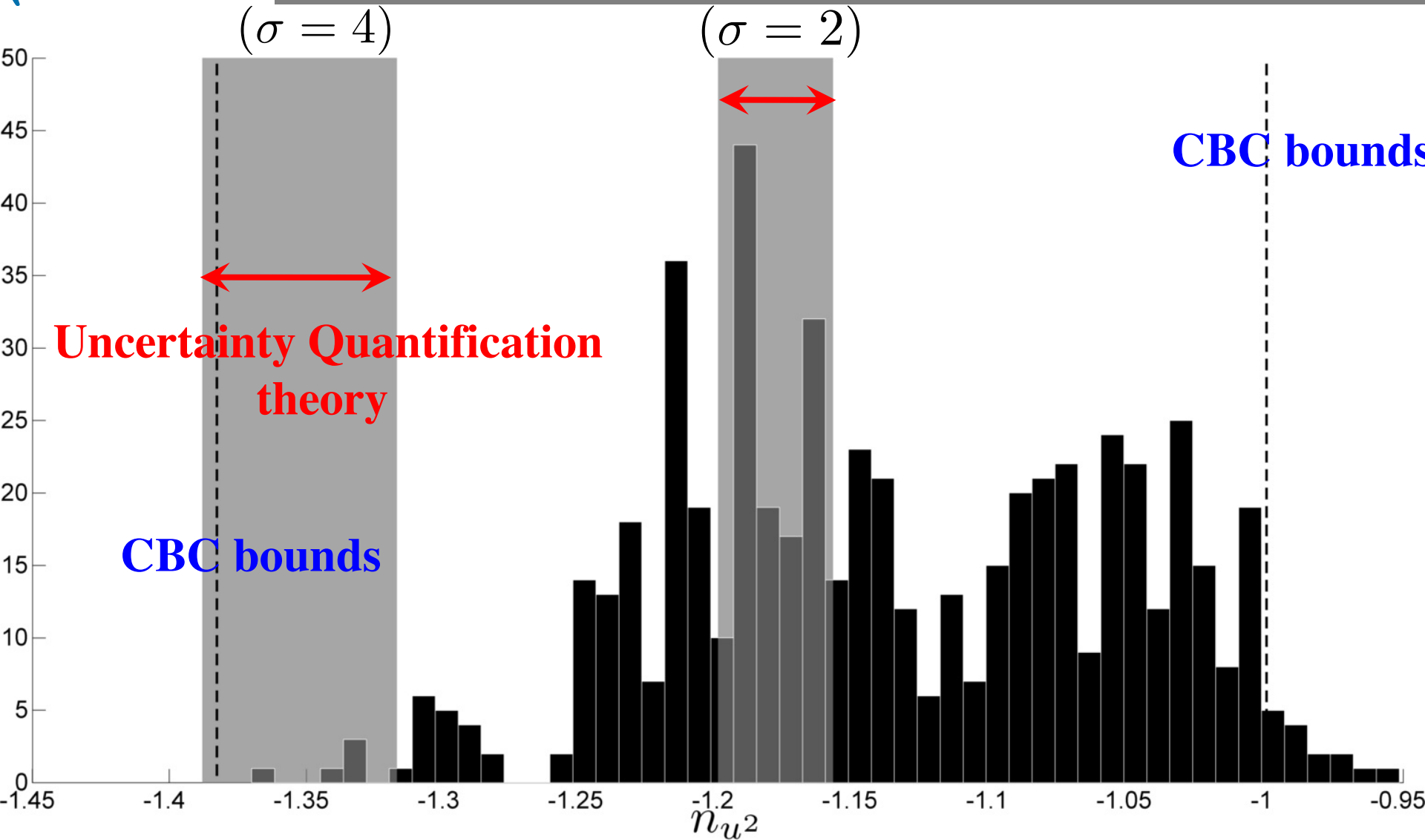
FOURIER SPACE

Hierarchy of CFD methods

Multiscale & Multiresolution approaches for turbulence, 2nd edition

Sagaut, Deck & Terracol, Imperial College Press, 2013





[Meldi et al. , J. Fluid Mech. 668, 2011]

[Meldi & Sagaut, J. Fluid Mech. 711, 2012]

- ***Fundamental theorem of Blind CFD***
 1. Solving equations of crude physical model
 2. Using poorly accurate numerical methods
 3. Implemented in a bugged code

→ It is possible to fit experimental data in many cases
- ***Proof***: daily practice !
- But when it doesn't work, we need to understand what is going on!

$$\frac{\partial u}{\partial t} + F(u, u) = 0$$

u : exact turbulent solution

$$\frac{\delta u_h}{\delta t} + F_h(u_h, u_h) = 0$$

u_h : discrete solution

LES:

- discrete solution $\Rightarrow$ **projection error**
- approximate integro-differential operators $\Rightarrow$ **discretization error**
- unresolved scales $\Rightarrow$ **resolution error**

Discrete numerical model

$$\frac{\delta u_h}{\delta t} + F_h(u_h, u_h) = 0$$

Equivalent continuous PDE

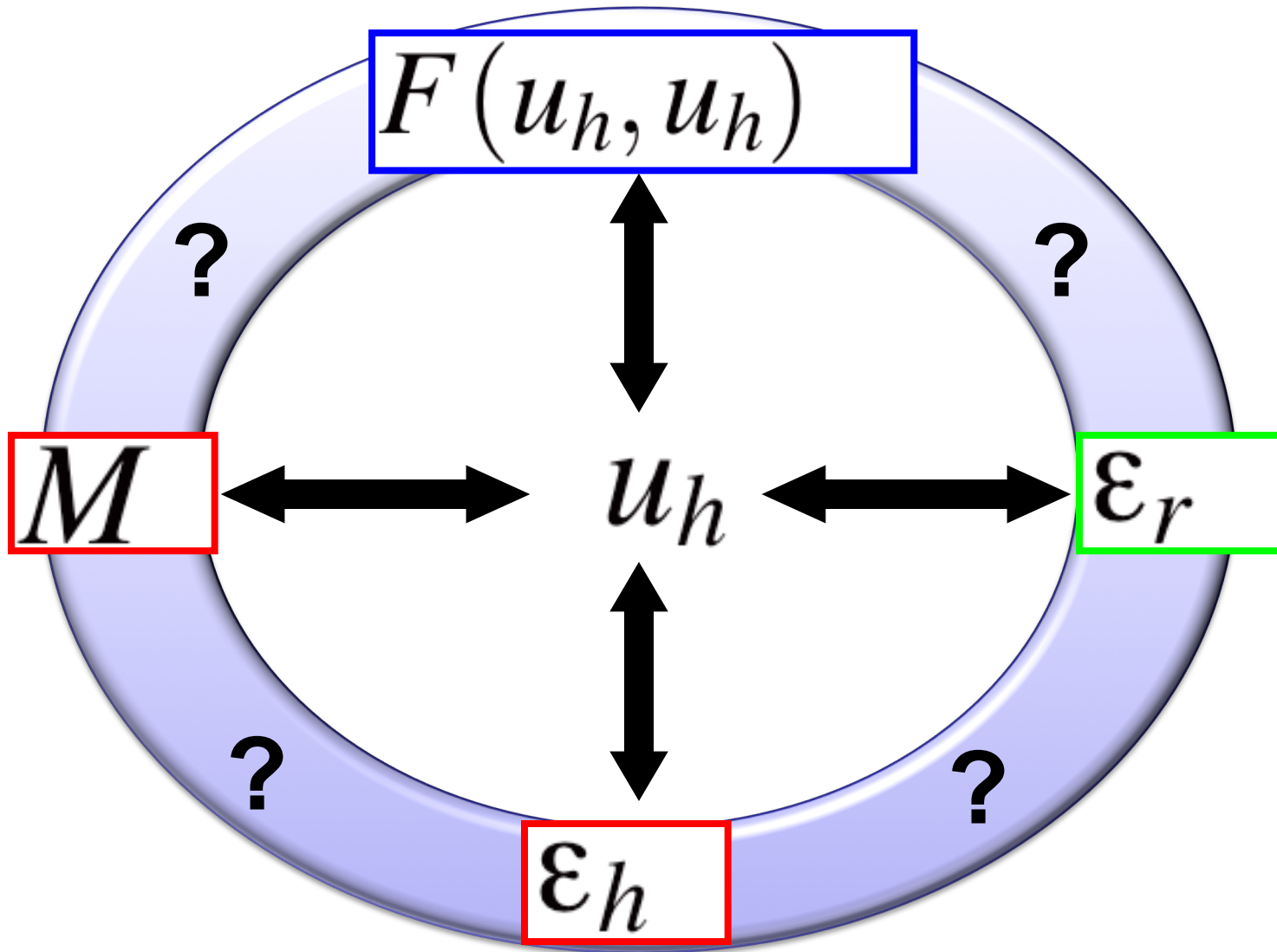


$$\frac{\partial u_h}{\partial t} + F(u_h, u_h) = M + \varepsilon_r + \varepsilon_h$$

Exact subgrid model

numerical error

Subgrid modelling error



$$e(h_r, \eta) = e_{\Pi}(h_r, \eta) + e_h(h_r, \eta) + e_r(h_r, \eta)$$

projection

discretization

resolution

Discretization/numerical error

$$\nabla_h(u) - \nabla(u) \neq 0 \implies e_h(h_r, \eta)$$

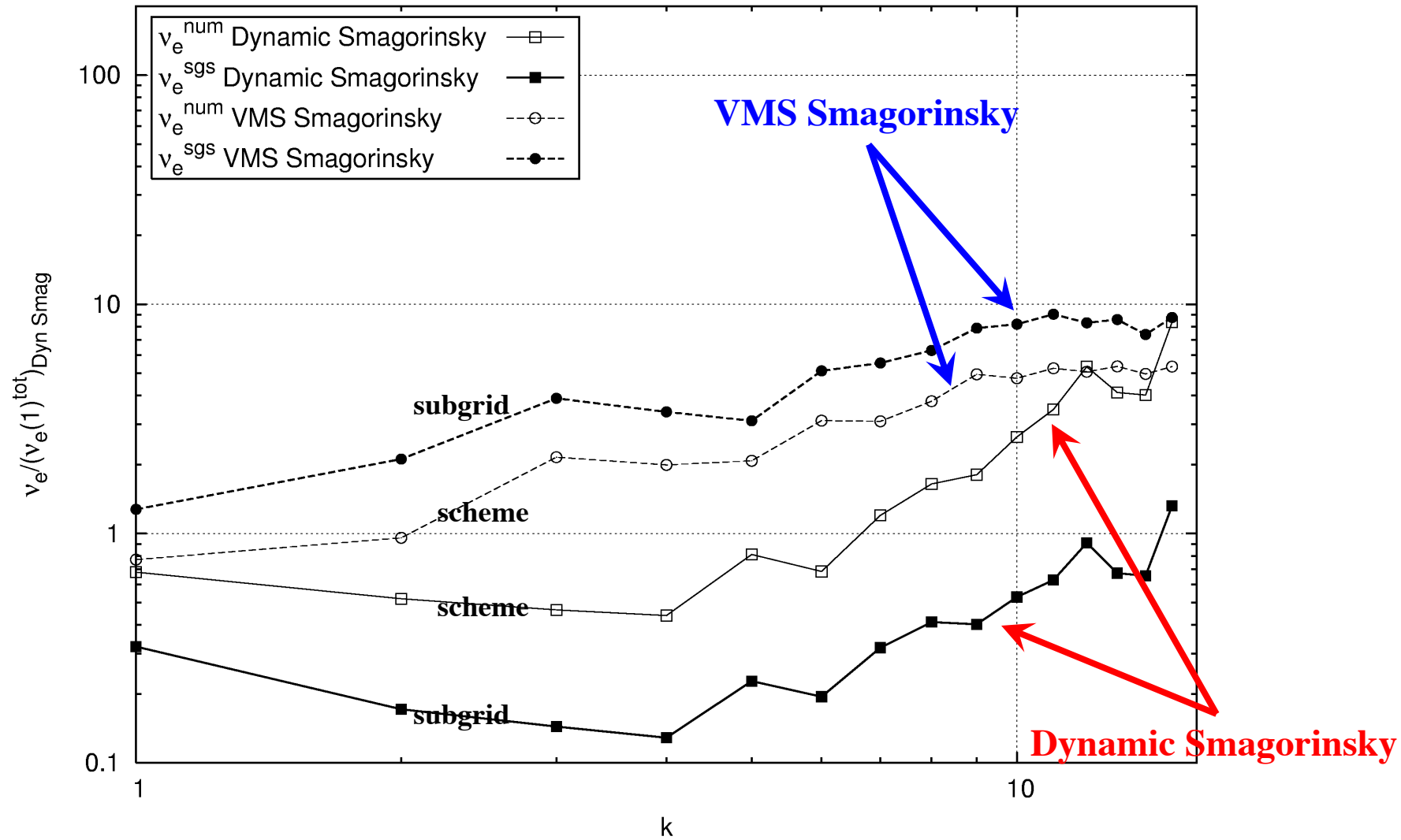
Resolution/modelling error

$$F_h(u_h, u_h) \neq F_h(u, u) = F(u, u) \implies e_r(h_r, \eta)$$

DNS as a limit (consistency)

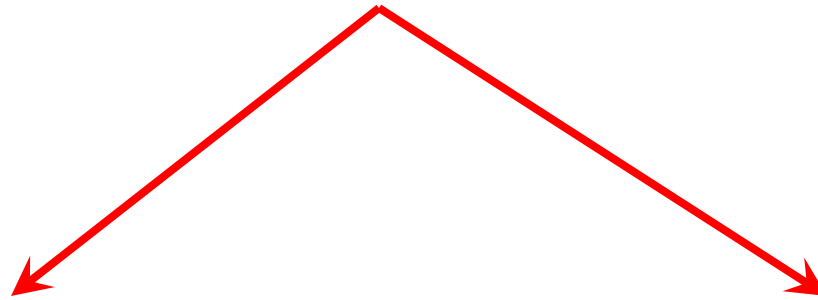
$$\lim_{h_r \rightarrow \eta} e(h_r, \eta) = 0$$

(Ciardi & al., JOT, 2006)



Definition of an optimal LES solution

$$\min(e_\pi + e_h + e_r) = \min(e_\pi)$$



$$e_h + e_r = 0$$

$$e_h = e_r = 0$$

Implicit LES:

- No physical subgrid model
- Ad hoc scheme

Classical LES:

- Explicit subgrid model
- Neutral scheme

2. Couplings between numerics and SGS models

- **Linearized analysis:**
 - reminiscent of von Neuman analysis for numerical schemes
 - linearized terms
 - monochromatic plane wave solution
 - Drawbacks: linearization, turbulent spectrum ignored, no feedback
- **Static analysis**
 - nonlinear terms
 - turbulent spectrum prescribed
 - Drawback: no feedback
- **Dynamic analysis**
 - nonlinear terms
 - turbulent solution computed
 - feedback & couplings taken into account
 - Issue: simplified model

Lin's equation for turbulent spectrum (1939)

$$\frac{\partial}{\partial t} E(k, t) + 2\nu k^2 E(k, t) = T(k, t)$$

EDQNM model (Orszag, 1971)

$$\frac{\partial}{\partial t} E(k, t) + 2\nu k^2 E(k, t) = T_{EDQNM}(k, t)$$

Exact/ideal EDQNM-LES

$$\frac{\partial}{\partial t} \bar{E}(k, t) + 2\nu k^2 \bar{E}(k, t) = T_{EDQNM}^{<}(k, t) + T_{EDQNM}^{>}(k, t)$$

EDQNM-LES with subgrid model

$$\frac{\partial}{\partial t} \bar{E}(k, t) + 2\nu k'^2 \bar{E}(k, t) = T'_{EDQNM}(k, t) + T'_{sg}(k, t)$$

$$\frac{\partial}{\partial t} \bar{E}(k, t) + 2\nu k'^2 \bar{E}(k, t) = T'^{<}_{EDQNM}(k, t) + T'_{sg}(k, t)$$

- Discretization errors in viscous terms

$$2\nu k^2 \bar{E}(k, t) \longrightarrow 2\nu k'^2 \bar{E}(k, t)$$

- Discretization errors in convection & pressure terms

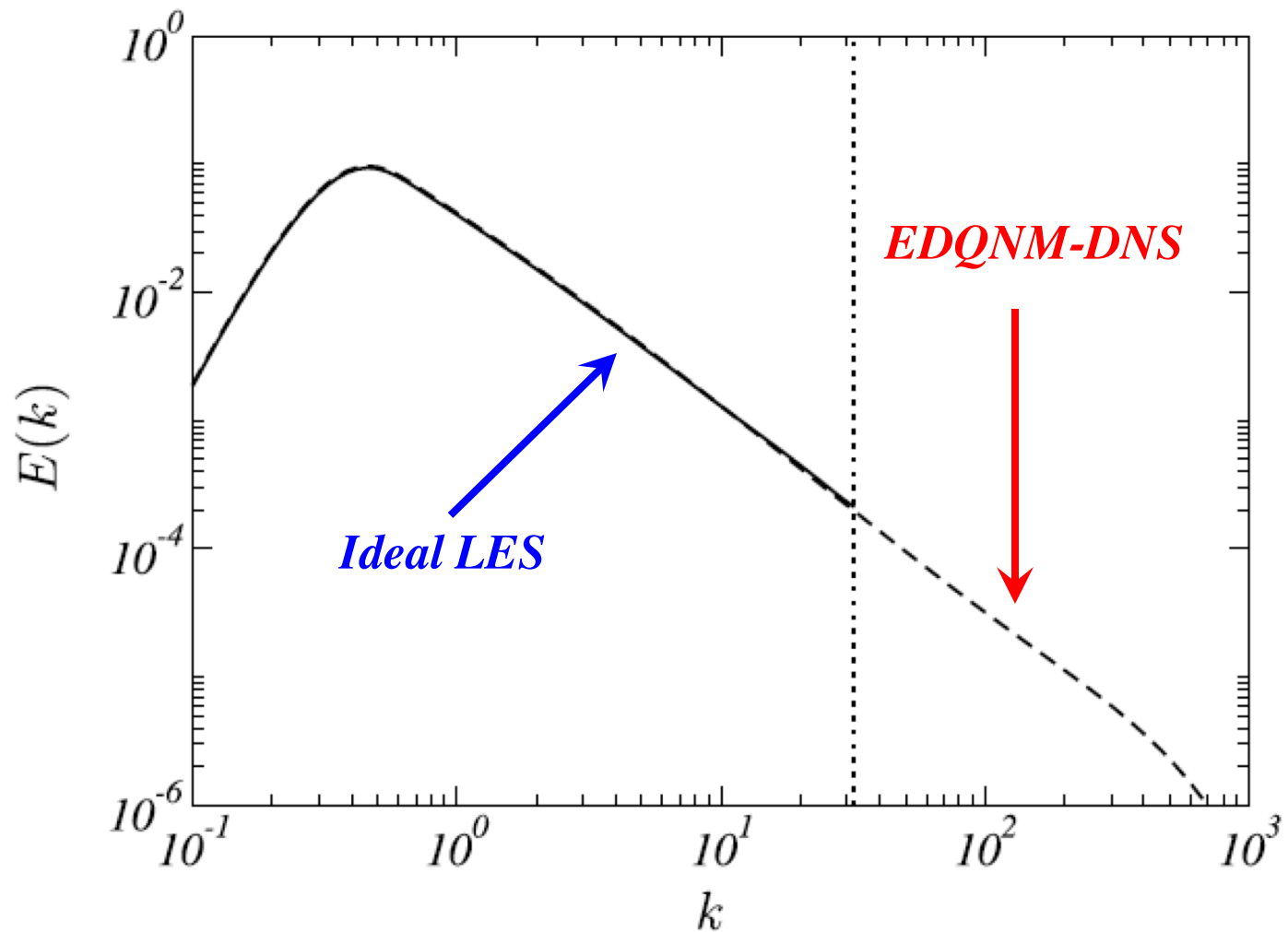
$$T^{<}_{EDQNM}(k, t) \longrightarrow T'^{<}_{EDQNM}(k, t)$$

- Subgrid modelling errors

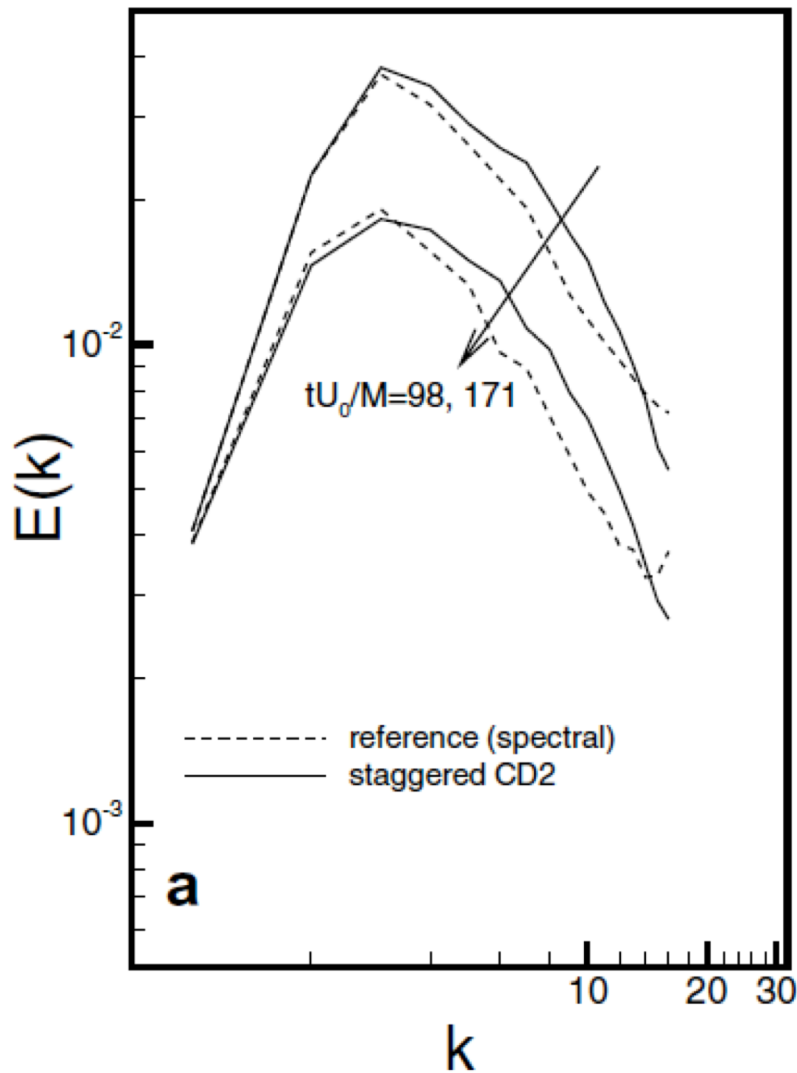
$$T^{>}_{EDQNM}(k, t) \longrightarrow T_{sg}(k, t)$$

- Discretization errors in subgrid model

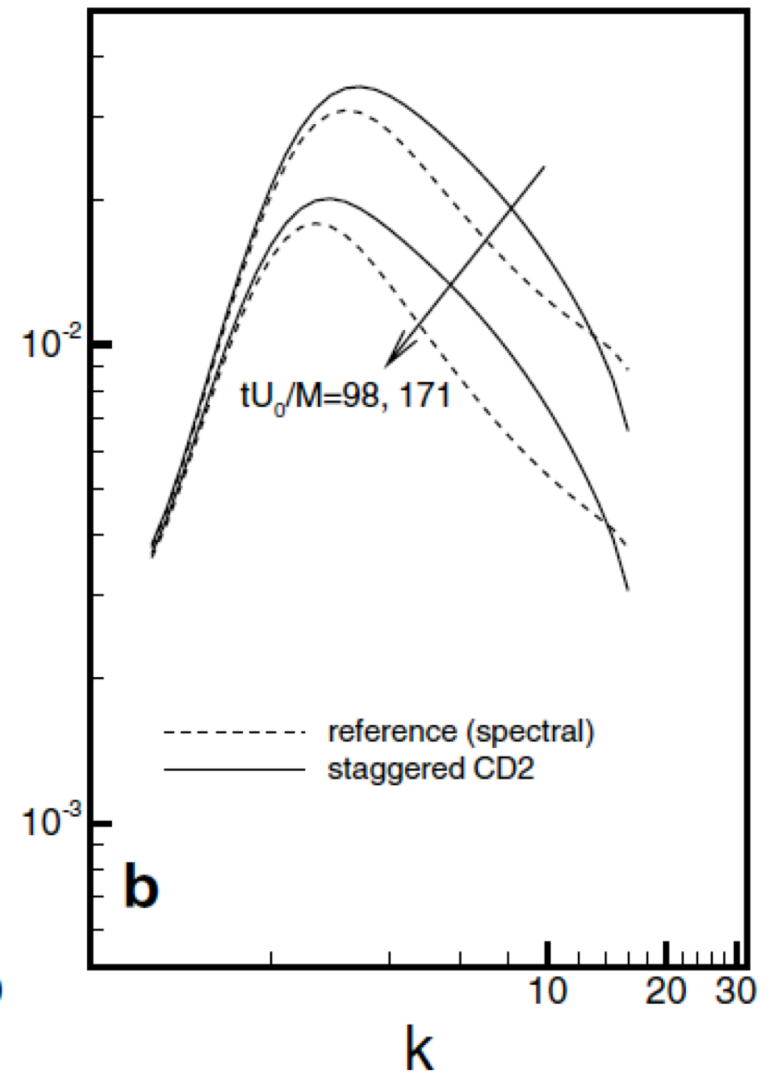
$$T_{sg}(k, t) \longrightarrow T'_{sg}(k, t)$$

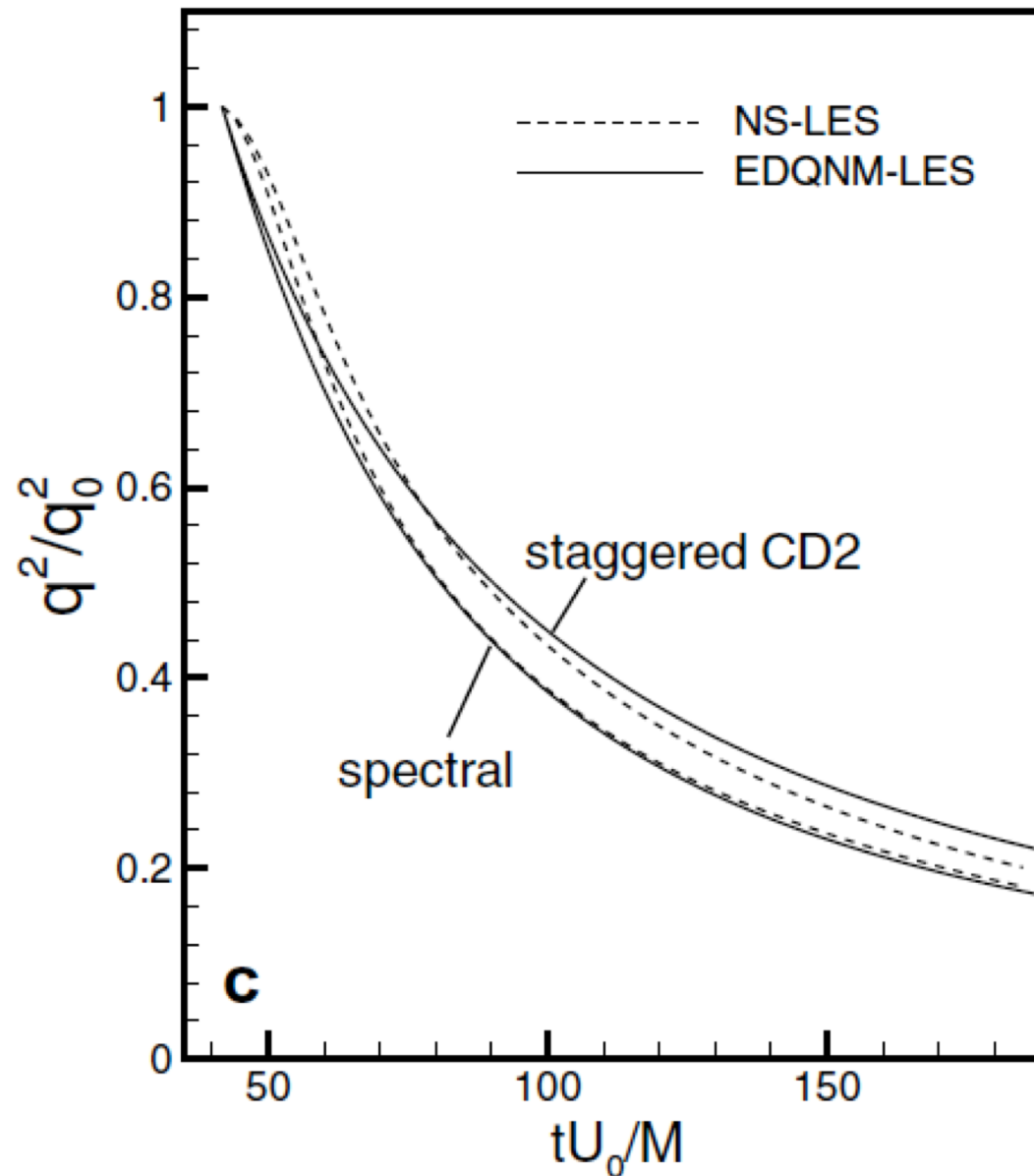


Real physical-space LES



EDQNM-LES





3-pt 2nd order scheme

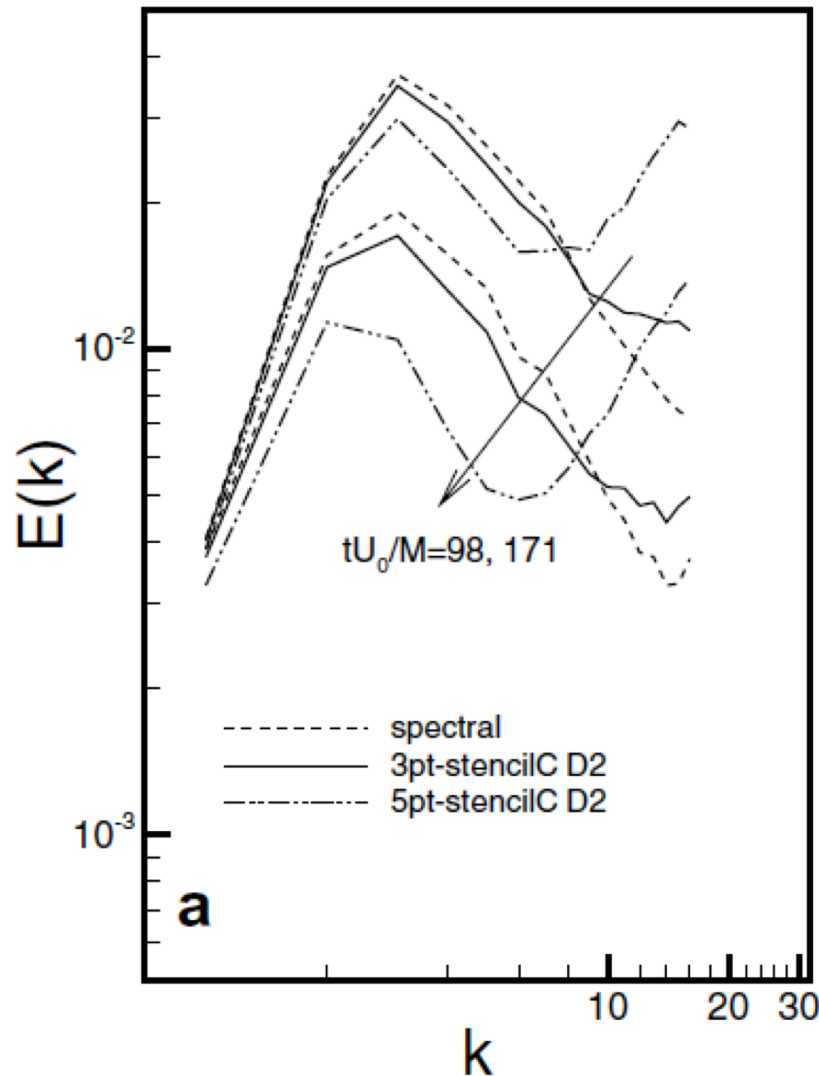
$$\left. \frac{\partial^2}{\partial x^2} u \right|_i \sim \frac{u_{i-1} - 2u_i + u_{i+1}}{h^2}$$

5-pt 2nd order scheme

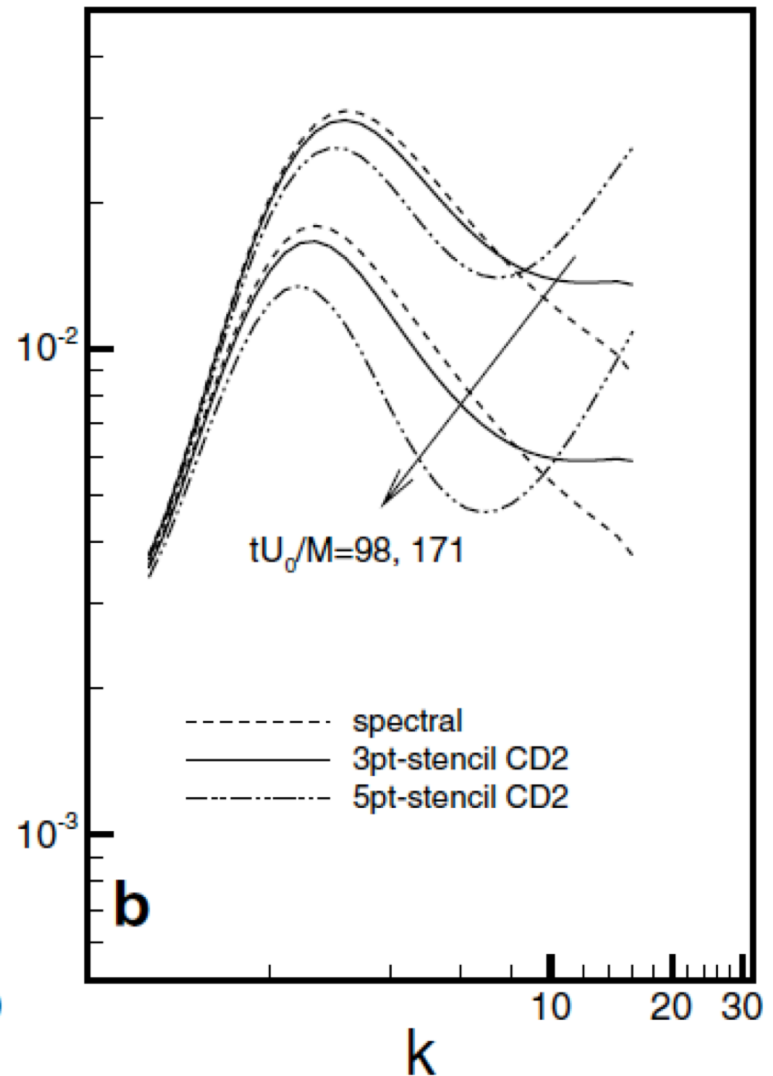
$$\left. \frac{\partial^2}{\partial x^2} u \right|_i \sim \frac{u_{i-2} - 2u_i + u_{i+2}}{4h^2}$$

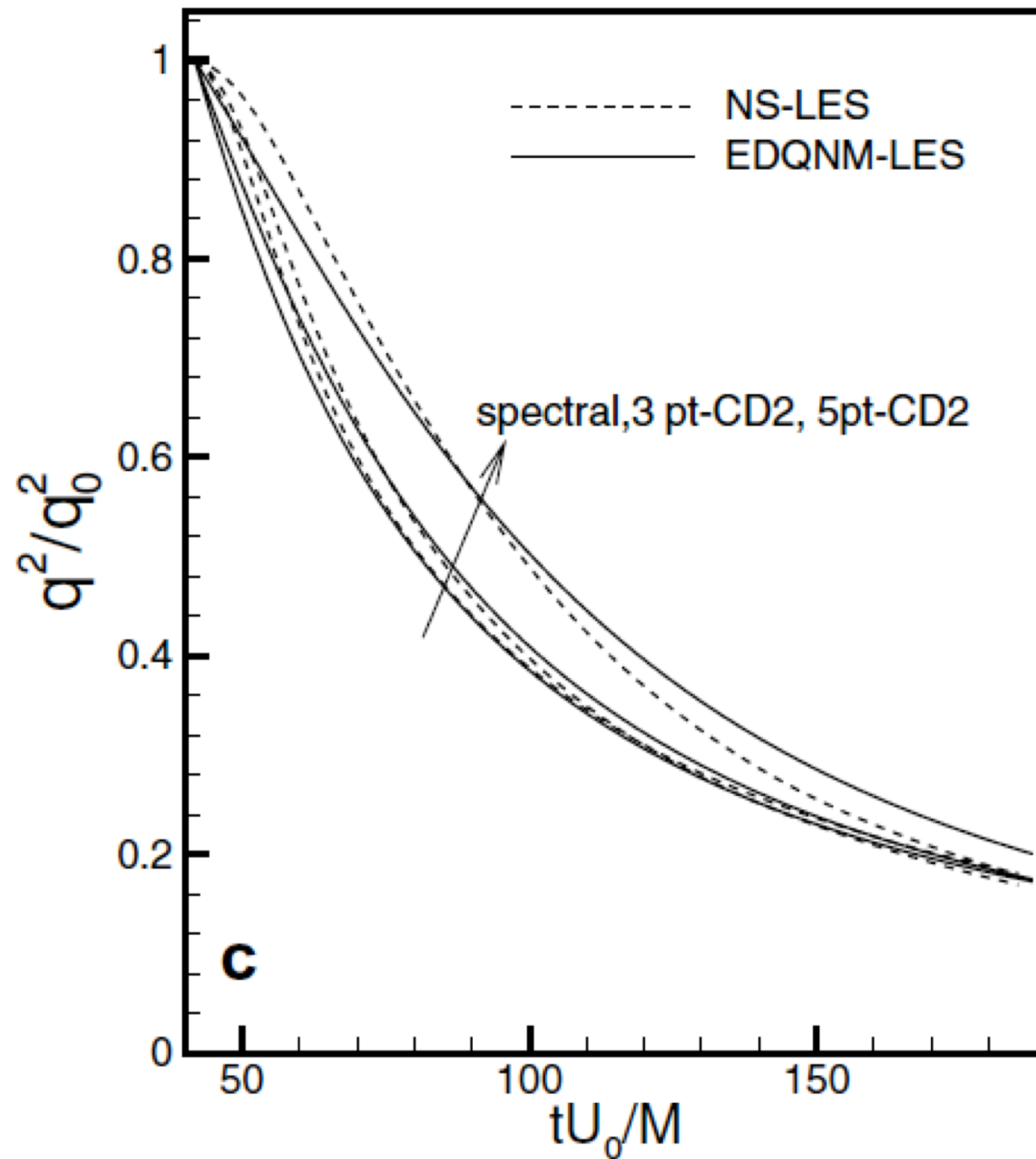
Ghost mode $u = (\dots, 1, -1, 1, -1, \dots)$

Real physical-space LES

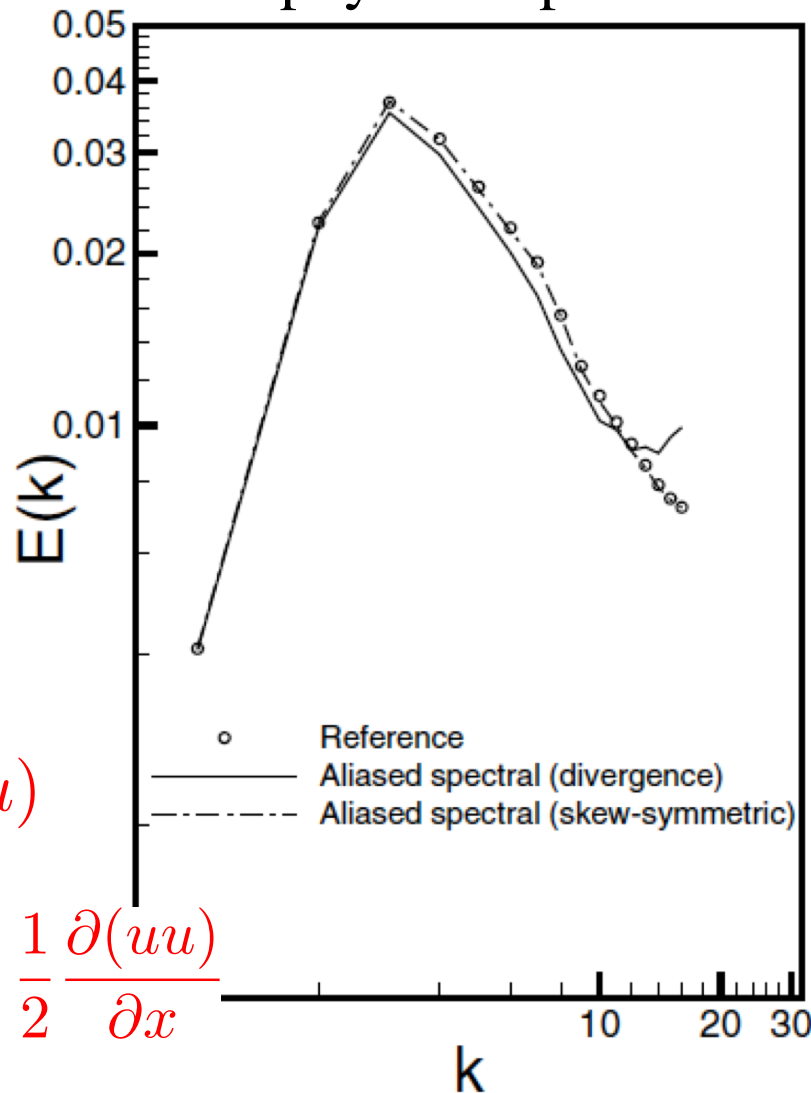


EDQNM-LES

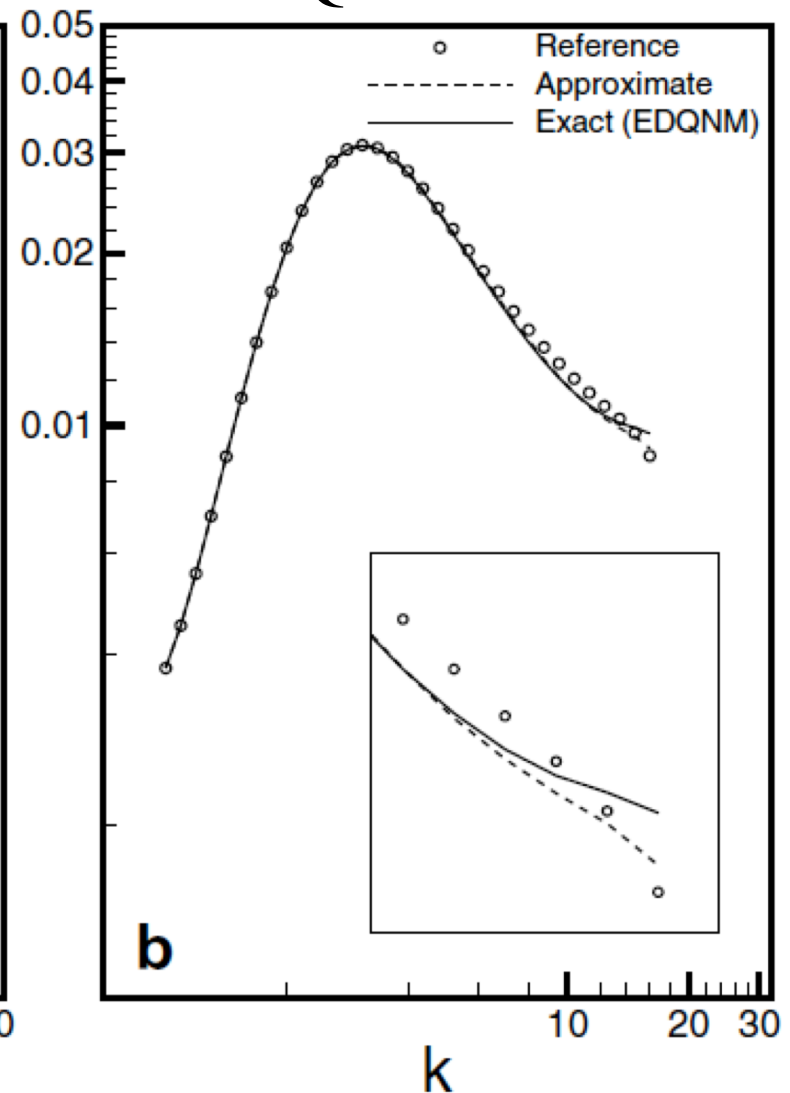




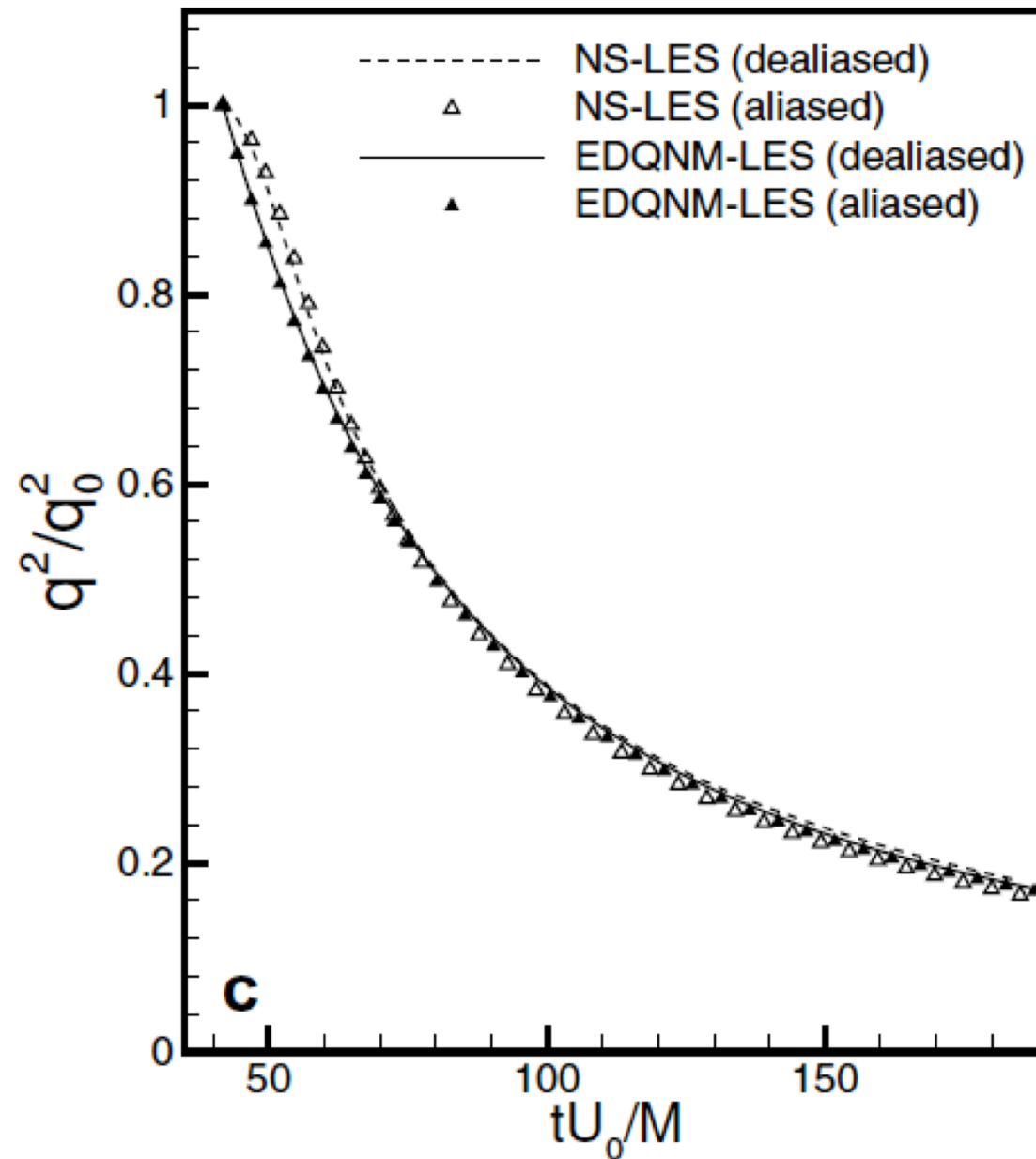
Real physical-space LES



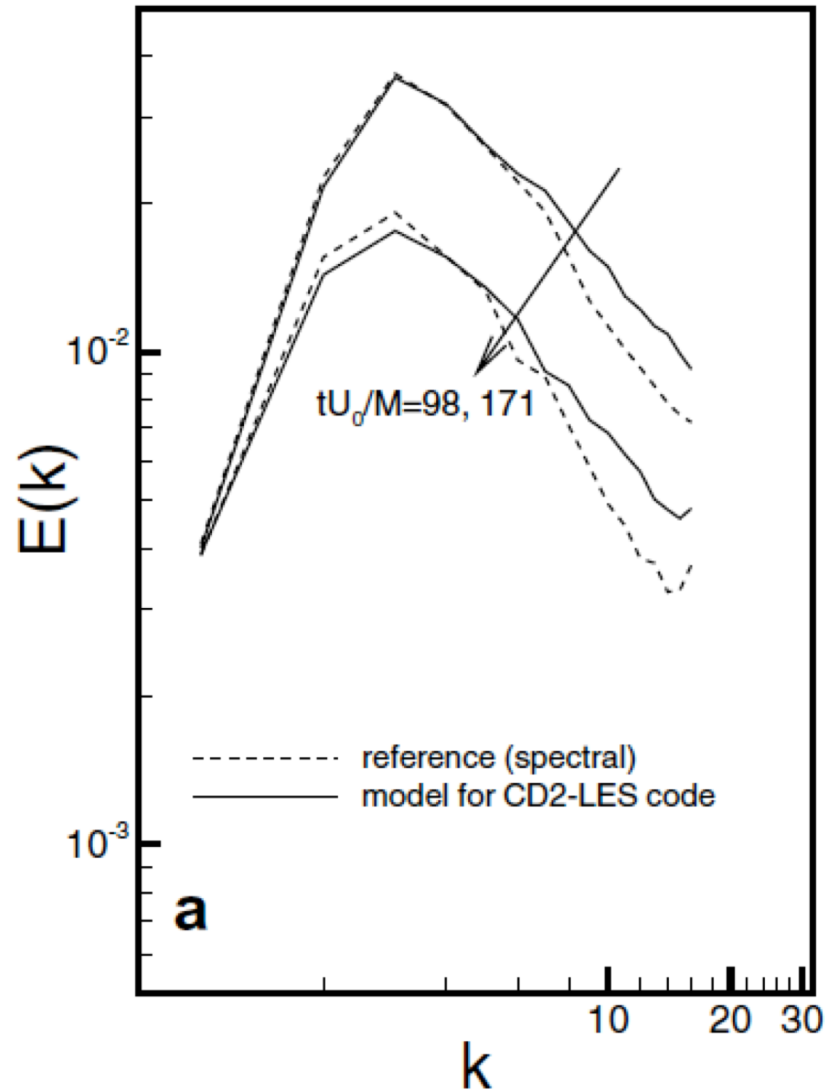
EDQNM-LES



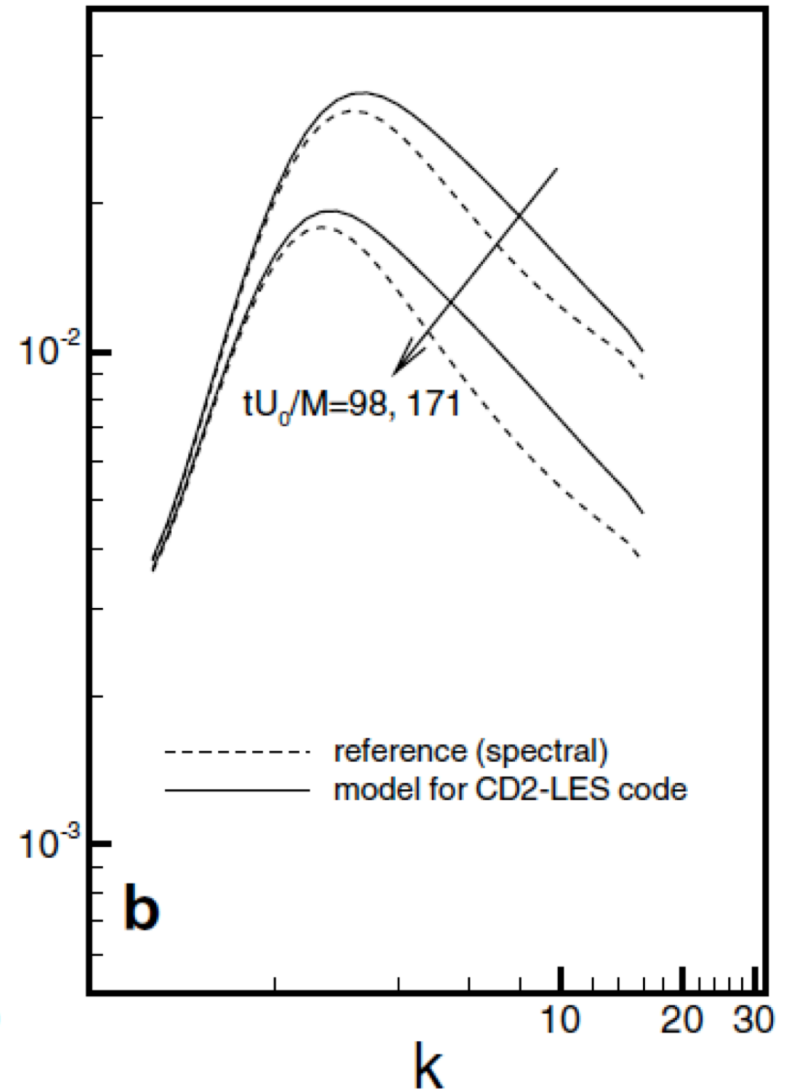
Sensitivity to aliasing errors

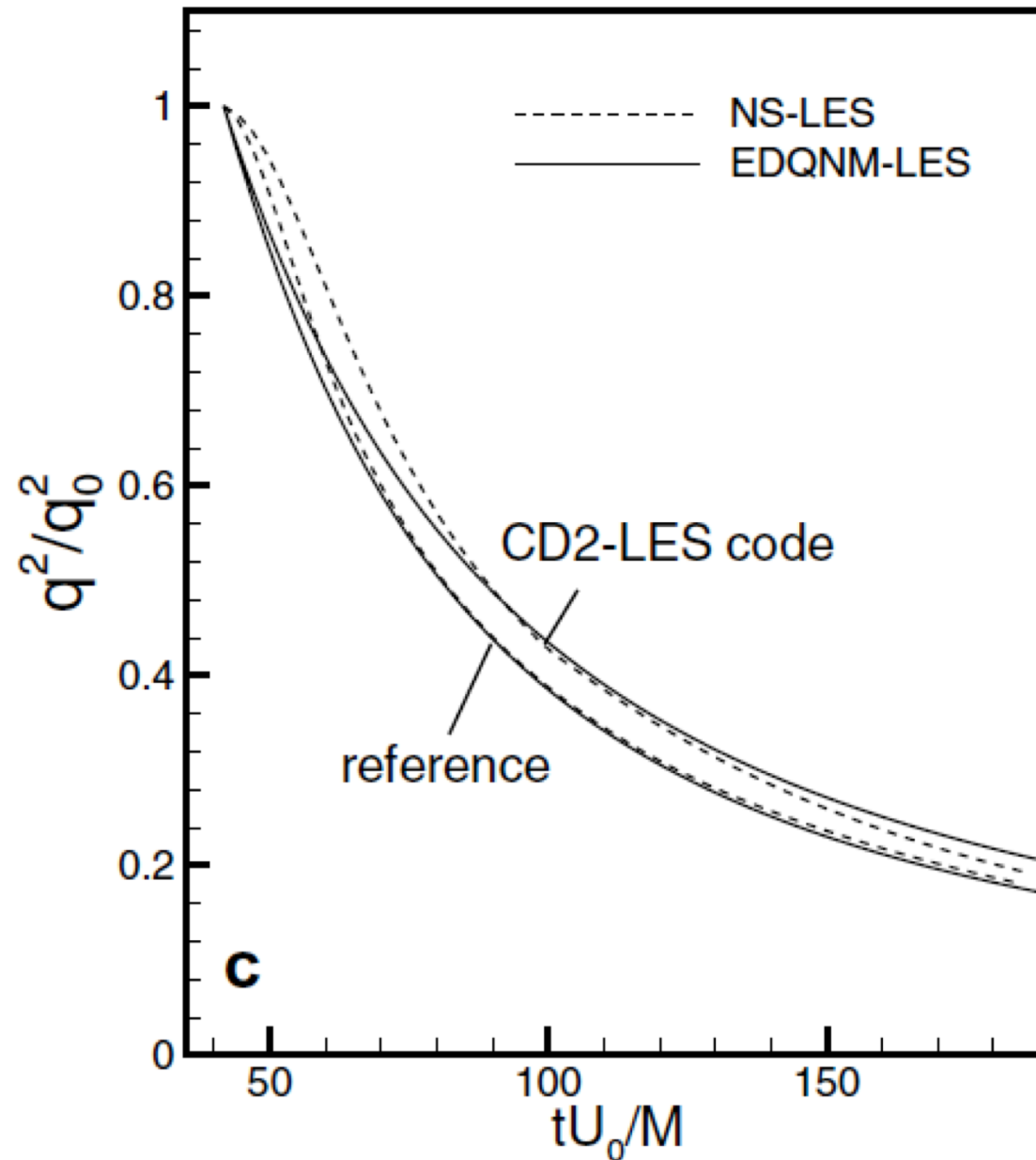


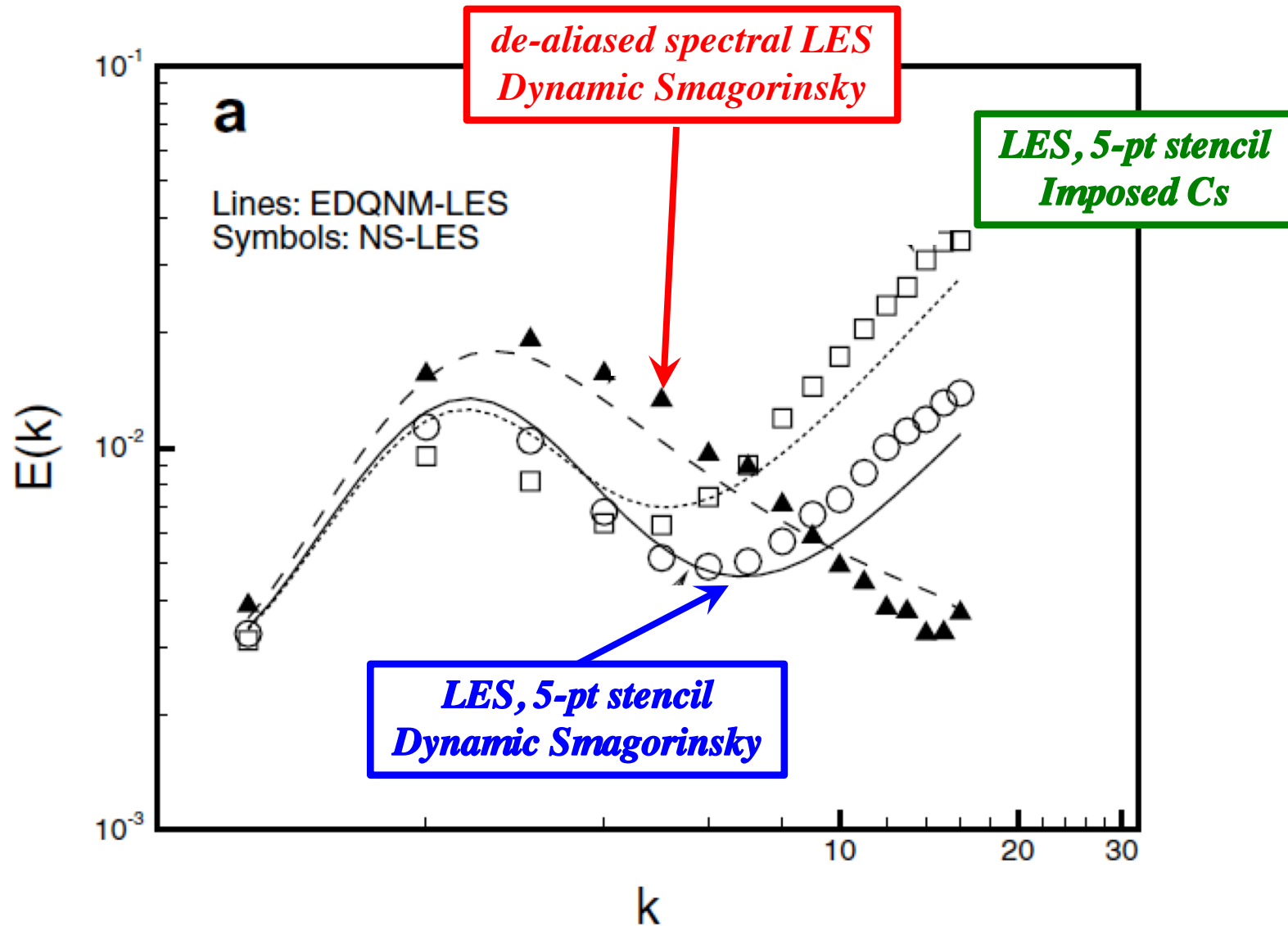
Real physical-space LES

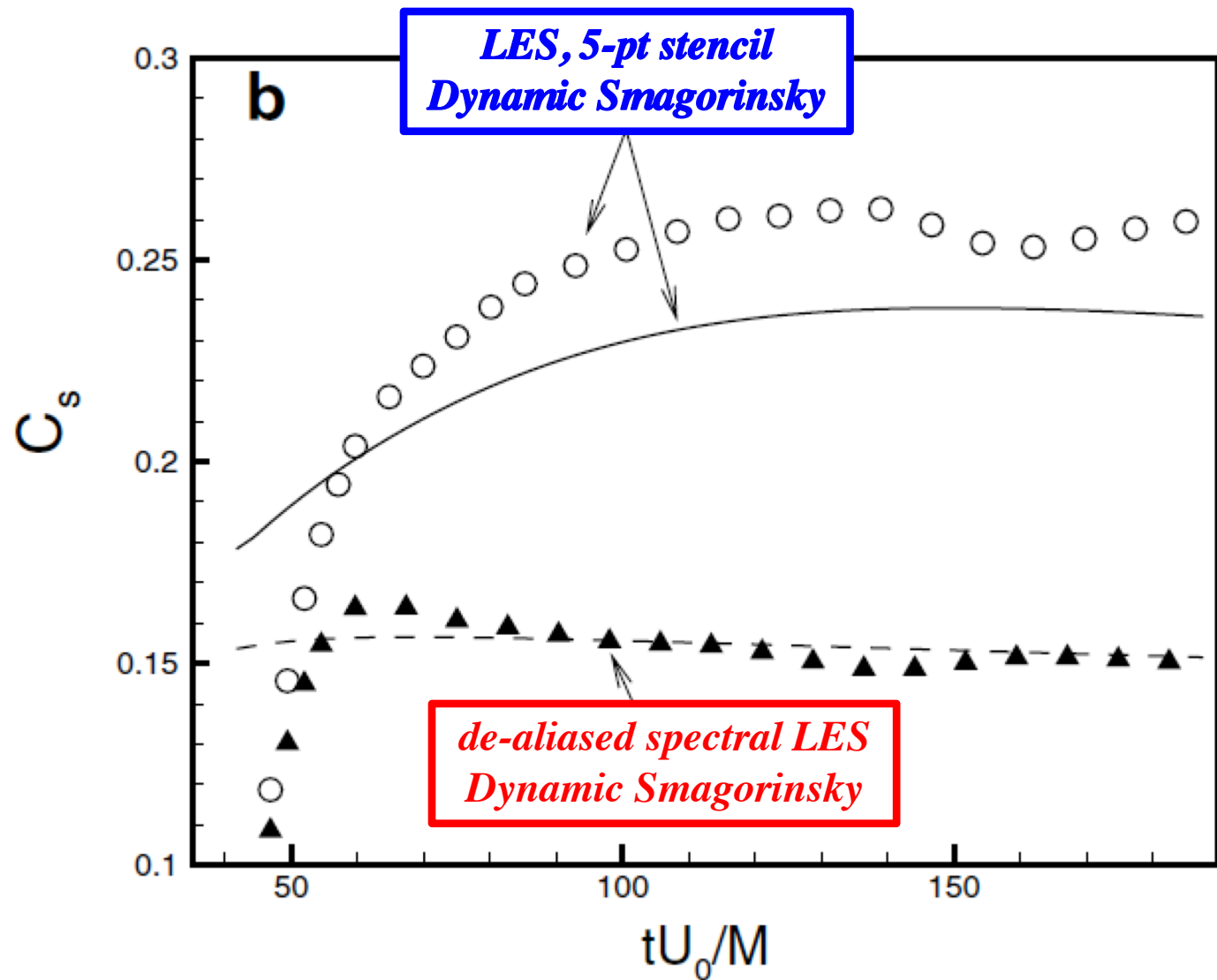


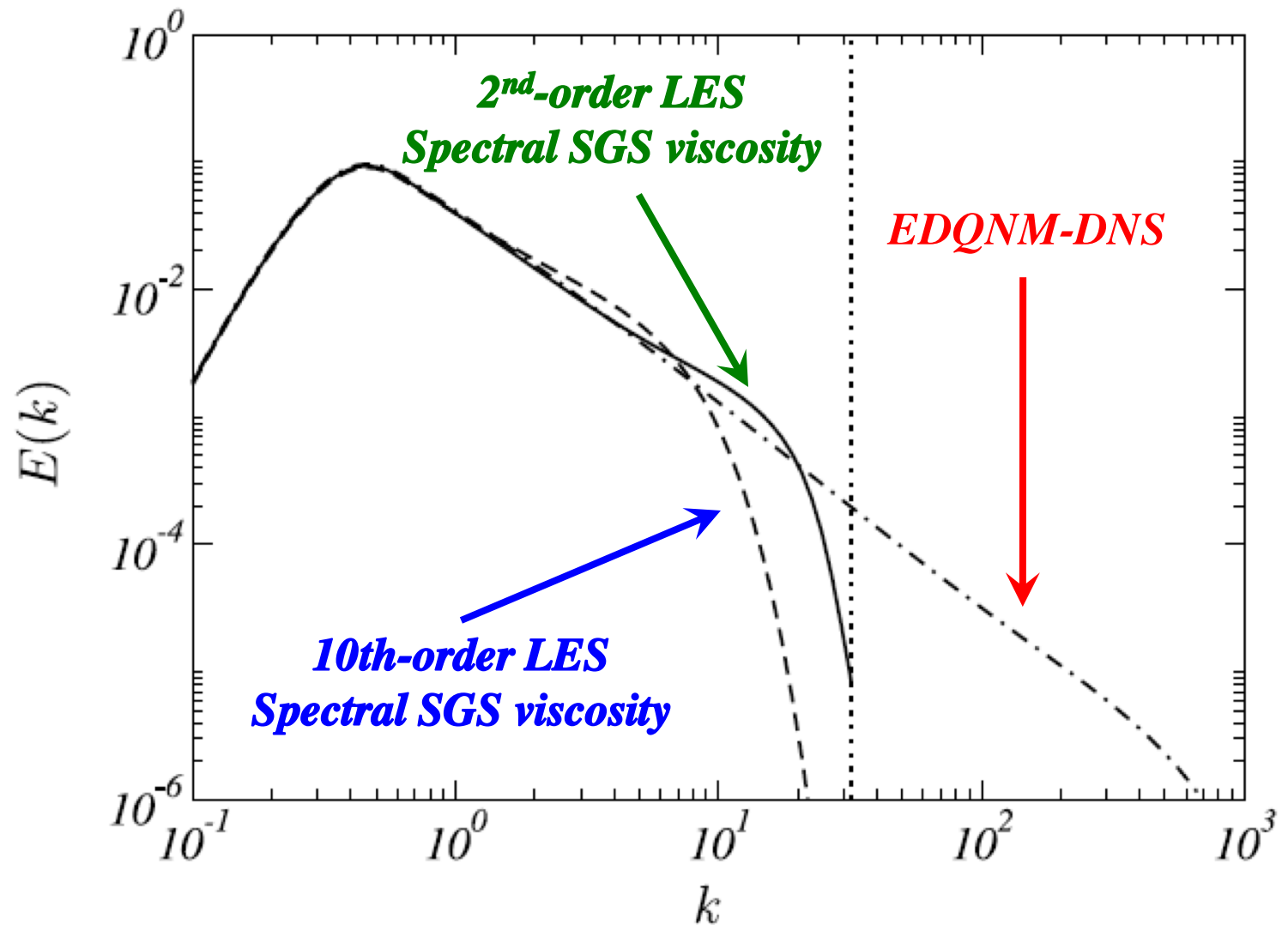
EDQNM-LES





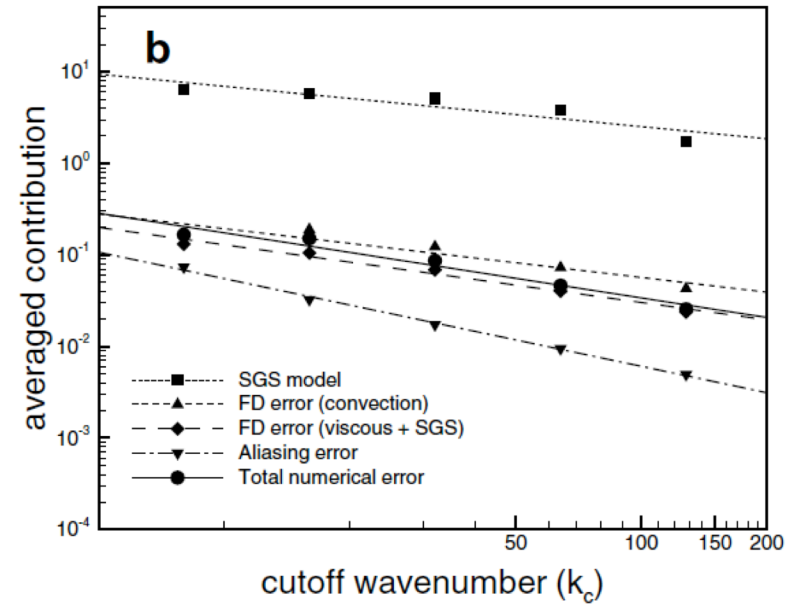
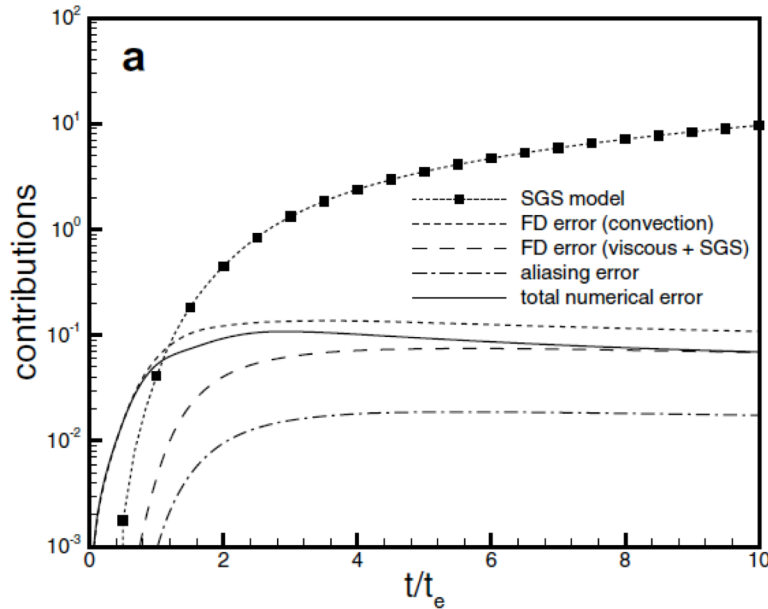




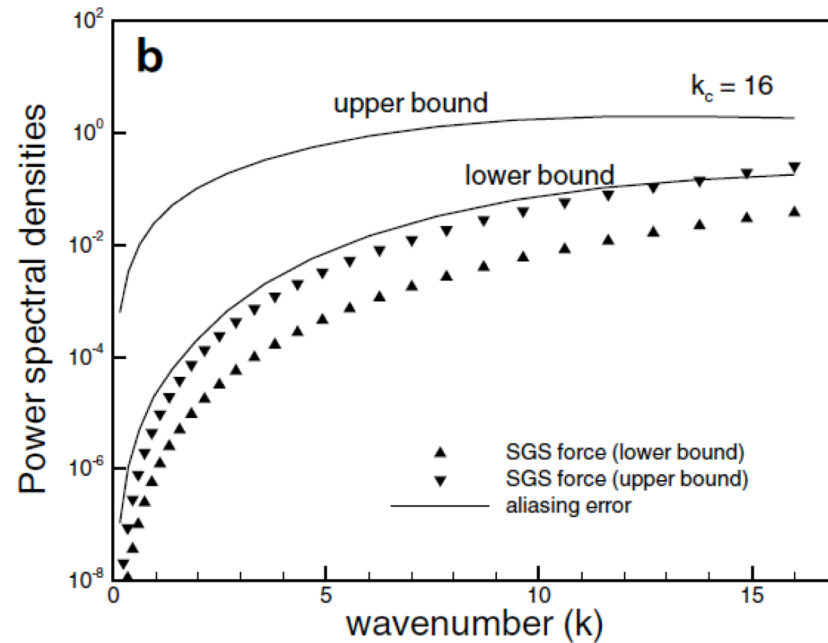
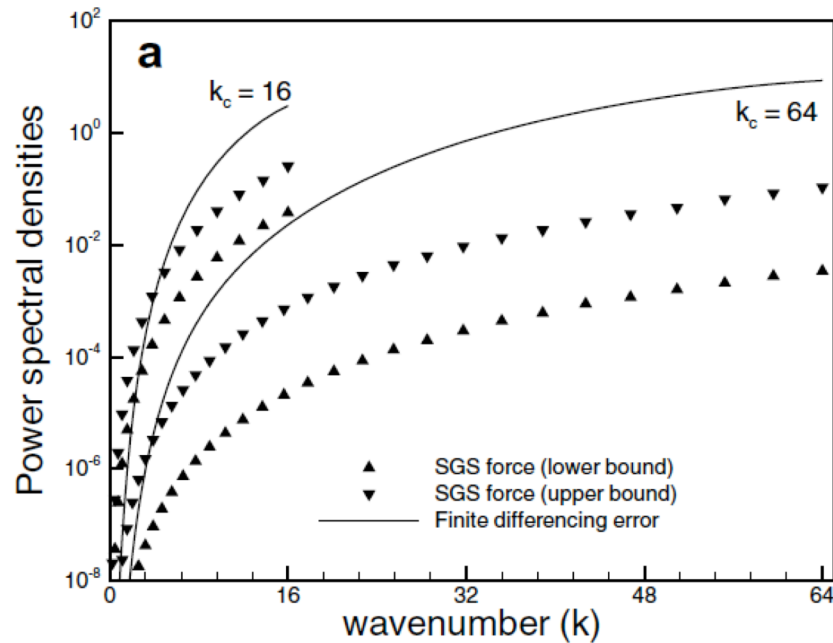


Static vs. Dynamic analysis

dynamic

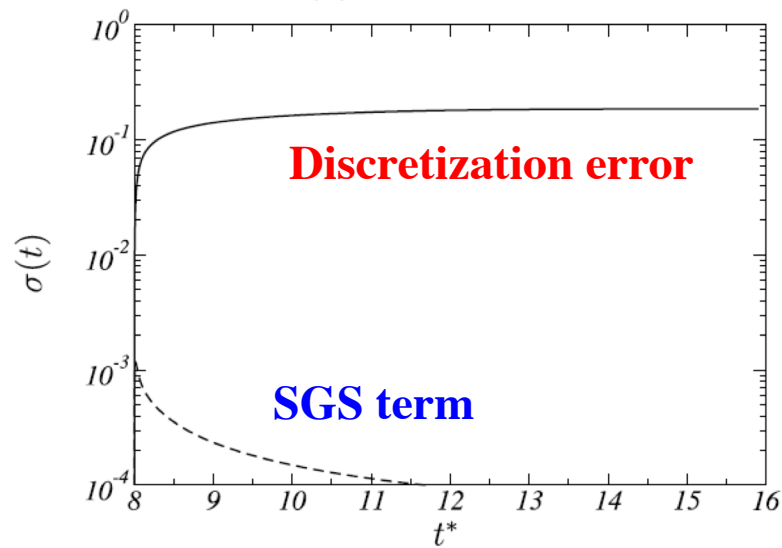


static

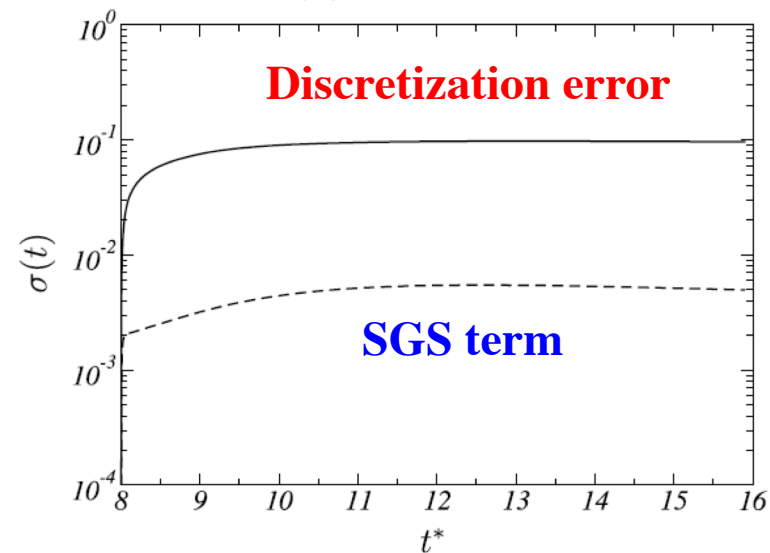


Effect of order of accuracy (*dynamic*)

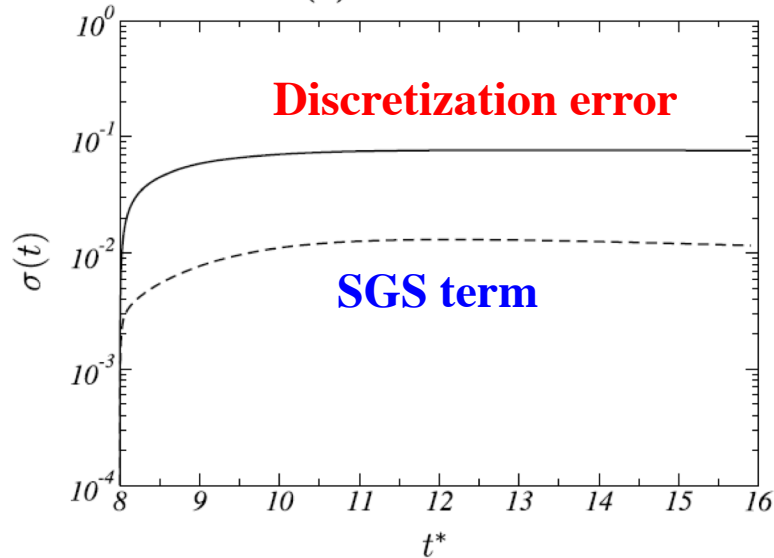
(a) - 2nd order



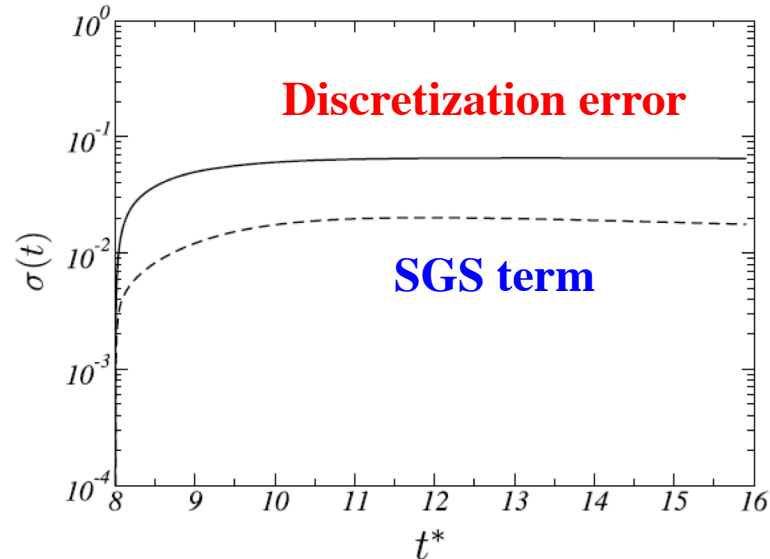
(b) - 6th order



(c) - 10th order



(d) - 14th order

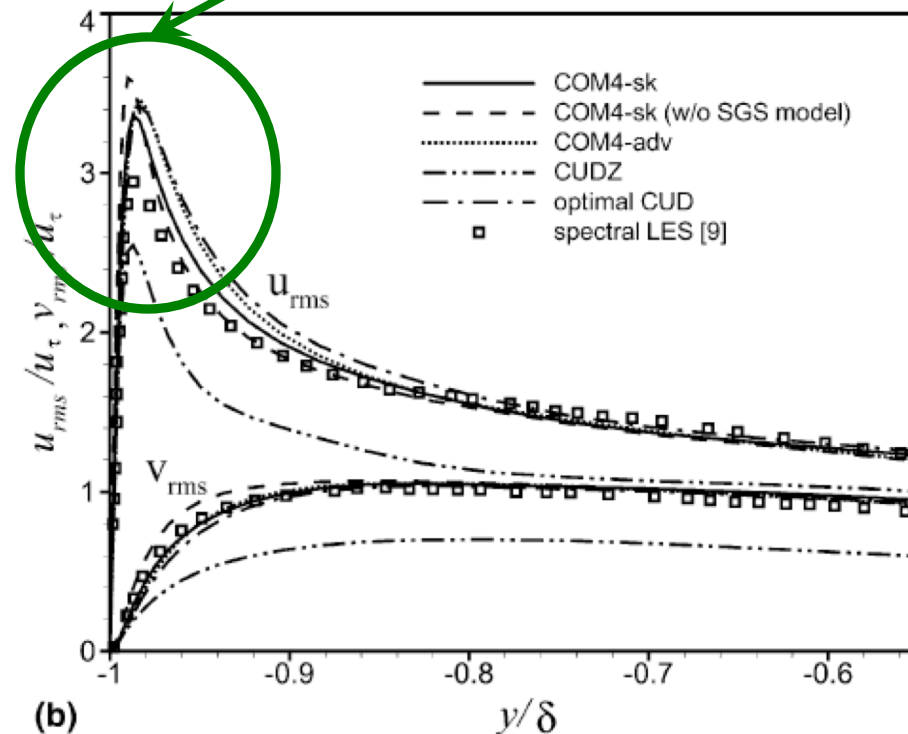
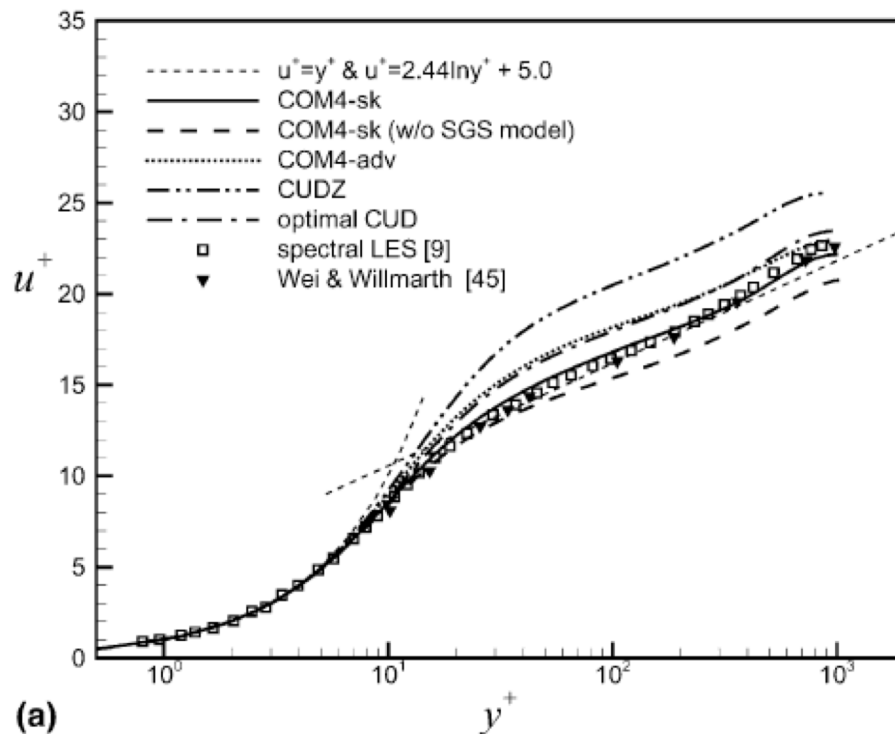


Simulations of turbulent channel flow with various discretization schemes

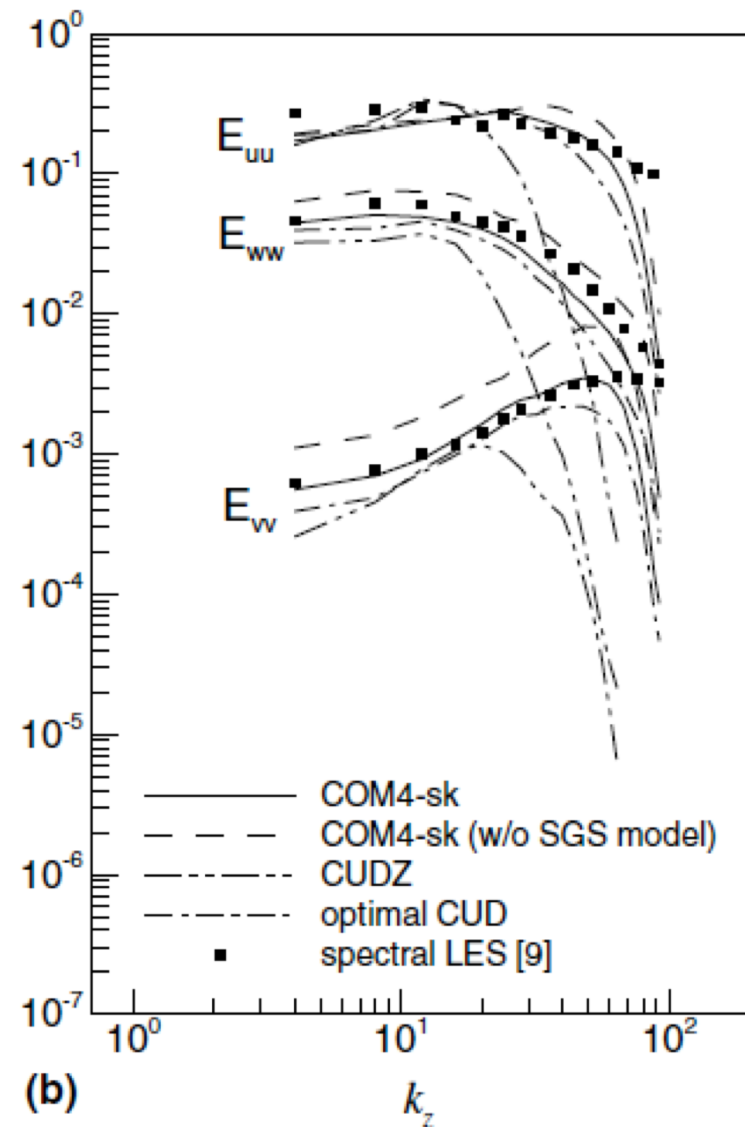
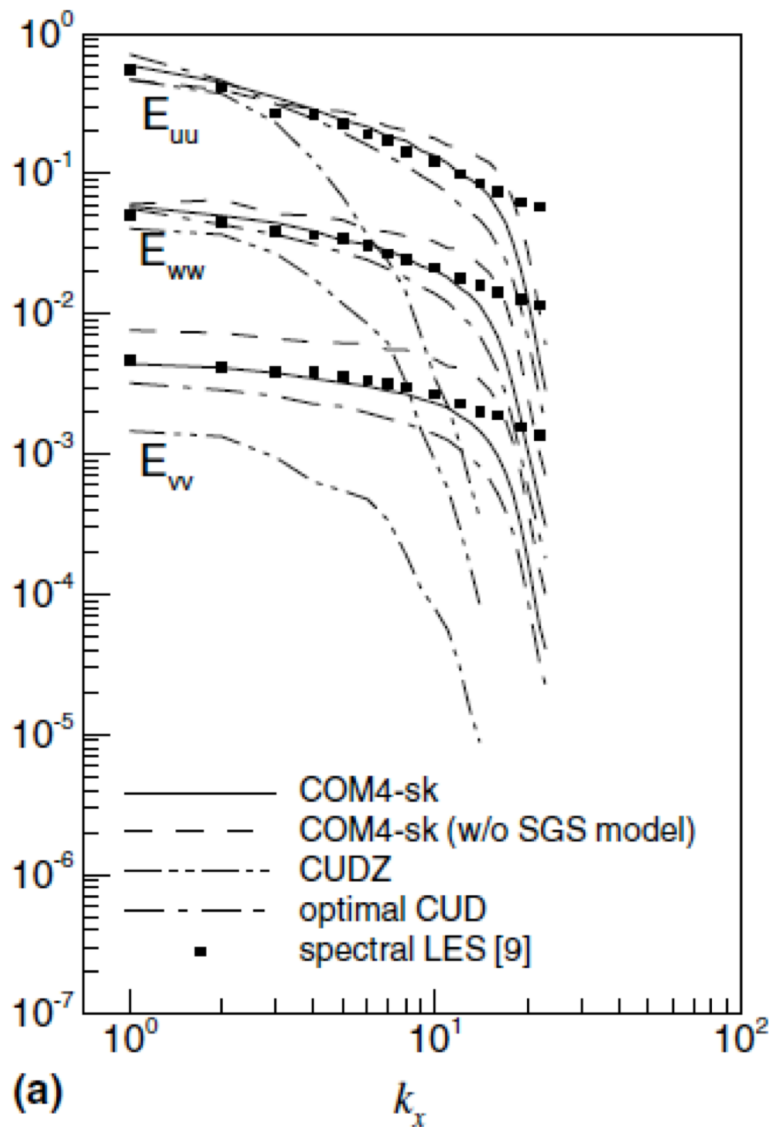
Scheme	Nonlinear terms	SGS model	Stability
COM4	Divergence	○	↑
COM4	Skew-symmetric	○, ×	●
COM4	Advective	○ (×)	● (†)
CUDZ	Divergence	○, ×	↓
CUD3	Divergence	○, ×	↓
CUD ($\alpha = 0.008$)	Divergence	○, ×	●
WENO3	Divergence	○, ×	↓

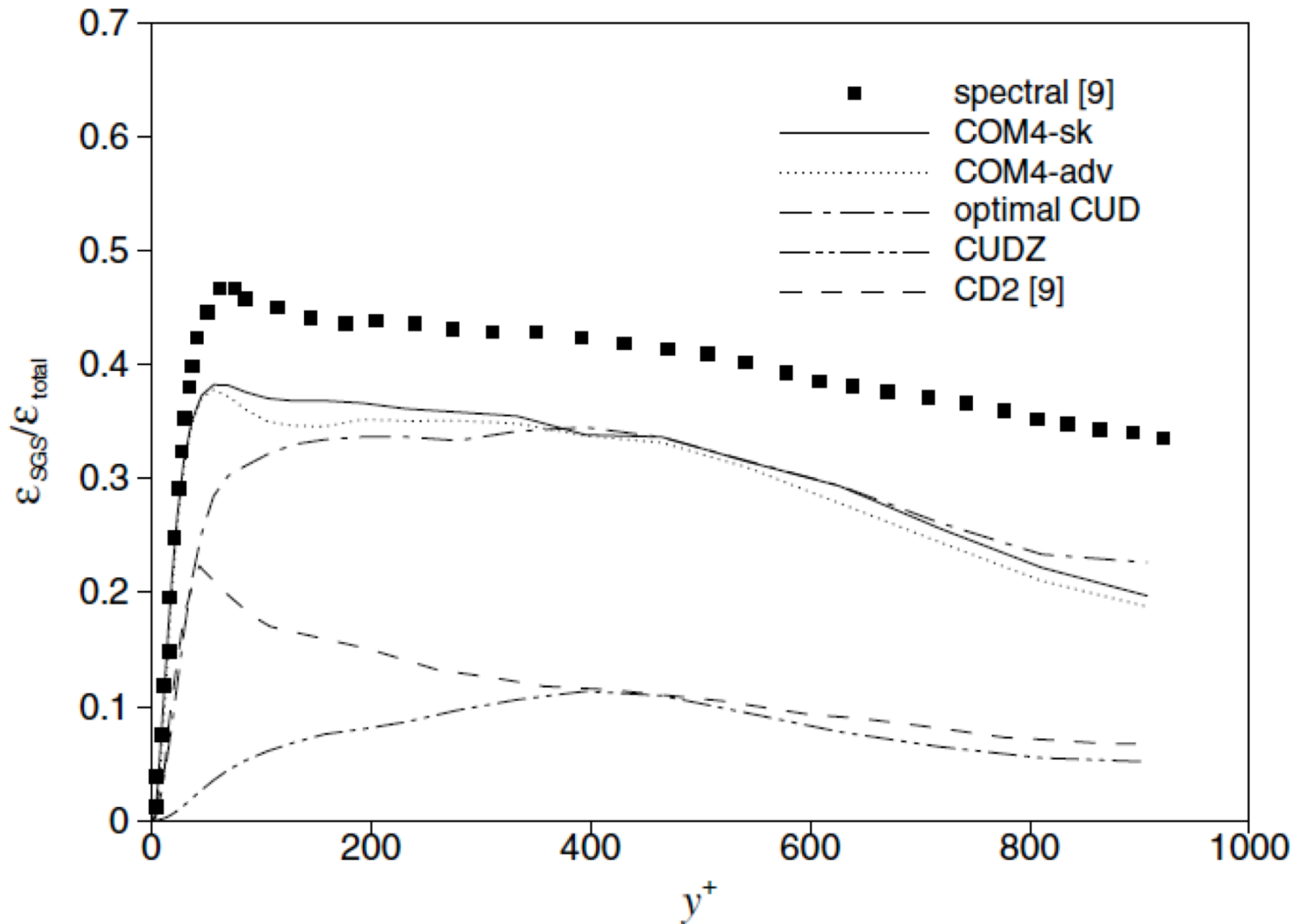
●, stable; ↑, numerically unstable; ↓, flow laminarizes.

Overshoot → no filter !

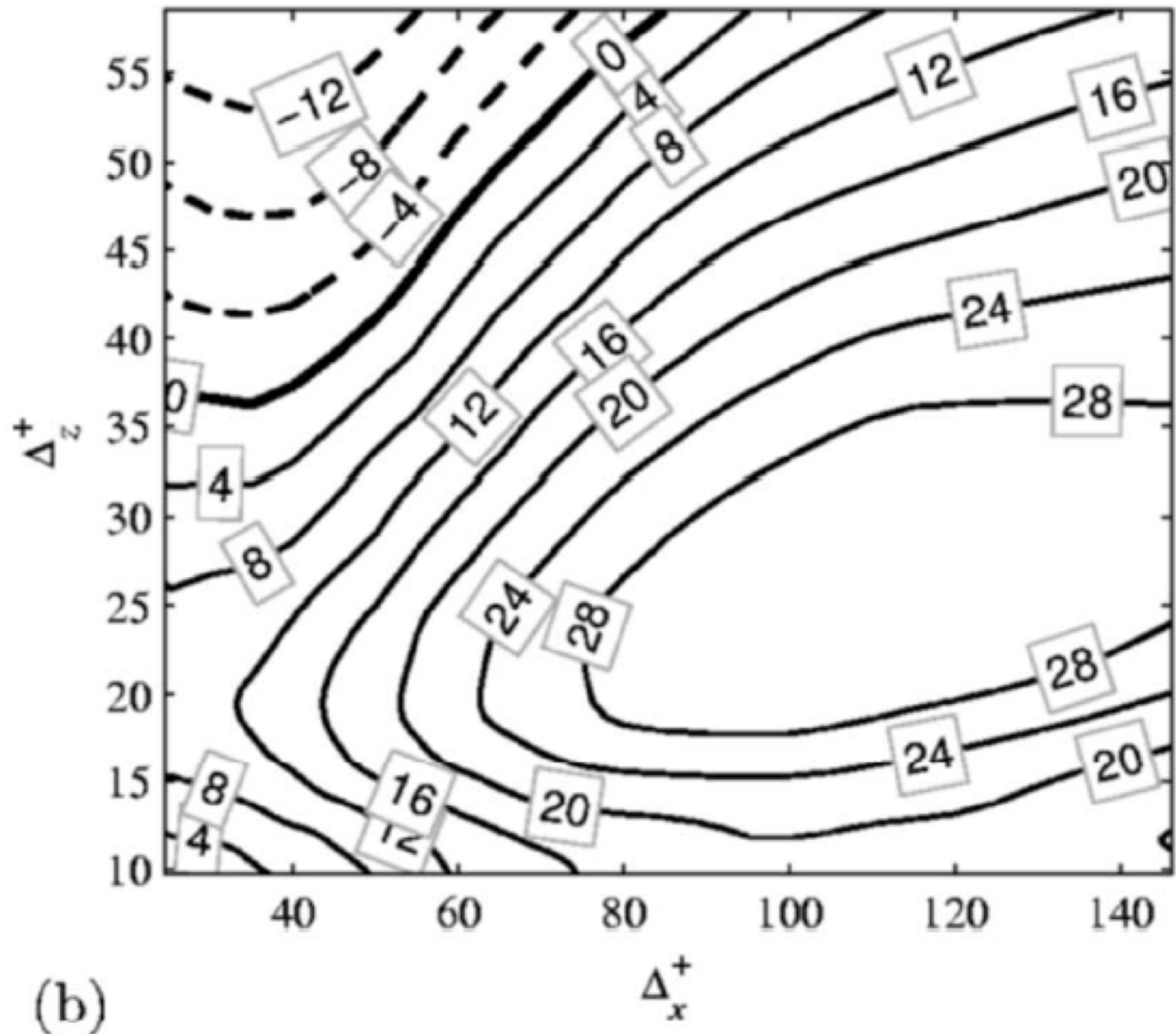


Channel flow case

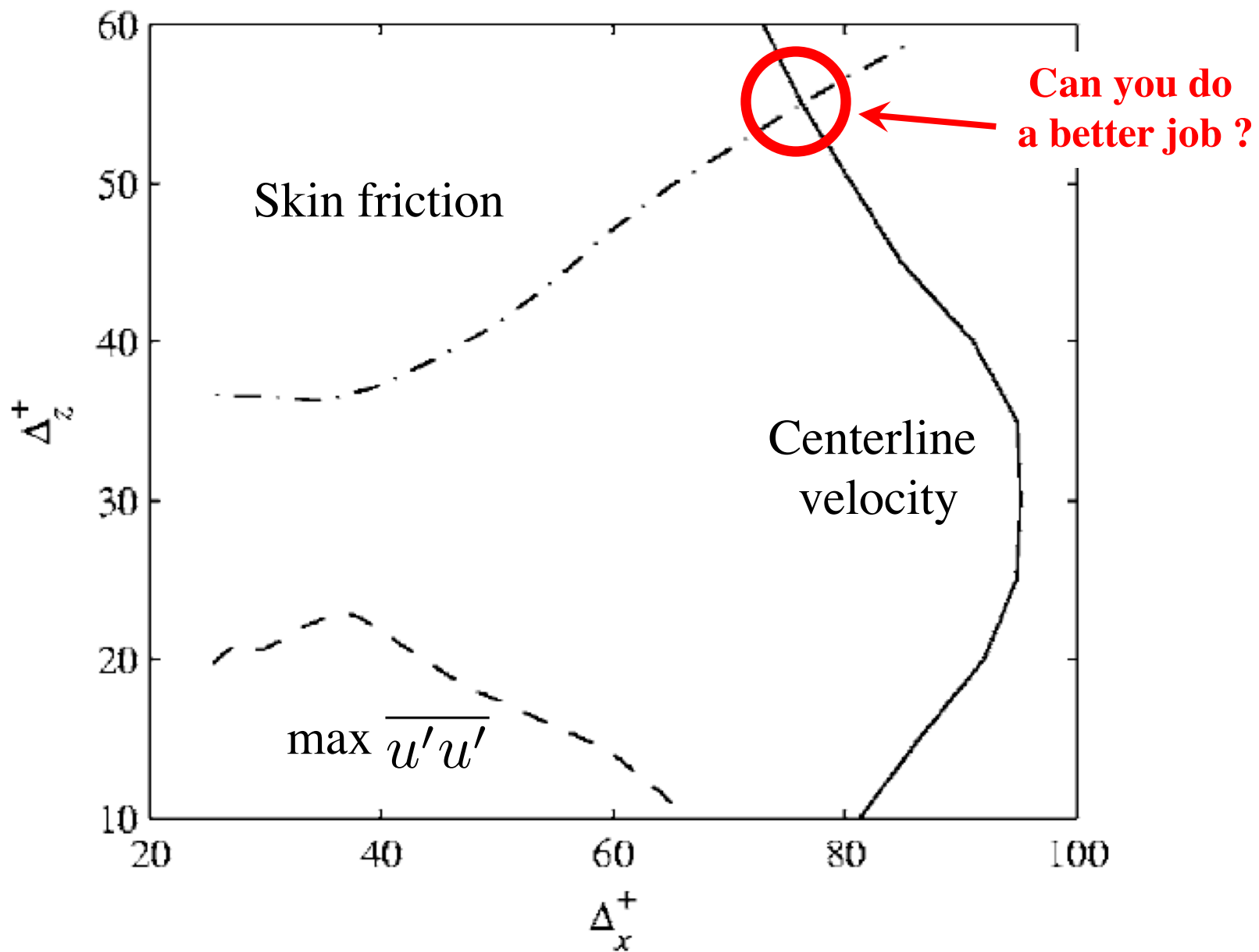




Skin friction error map (model-free)



Zero error trajectories



Circular cylinder ($Re = 3900$)

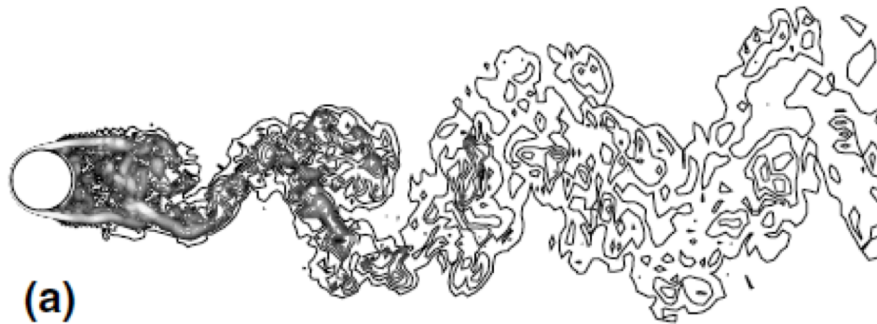
Flow parameters for the flow over a circular cylinder at $Re = 3900$

Data from	C_D	$-C_{p_b}$	St	L_{rec}/D	U_{min}
Run1	1.02	0.89	0.209	1.37	-0.33
Run2	1.04	0.93	0.209	1.26	-0.34
Run3	0.95	0.80	0.210	1.79	-0.27
Experiment [49–51]	0.99 ± 0.05	0.88 ± 0.05	0.215 ± 0.05	1.4 ± 0.1	-0.24 ± 0.1
UD7 (DSM) [19]	1.00	0.95	0.203	1.36	-0.32
UD7 (SM) [19]	0.96	0.81	0.209	1.74	-0.33
CD2 (DSM) [21]	1.00	0.93	0.207	1.40	-0.35
QUICK (SM) [22]	0.97	0.87	–	1.69	-0.24
B-spline (DSM) [47]	1.04	0.94	0.210	1.35	-0.37

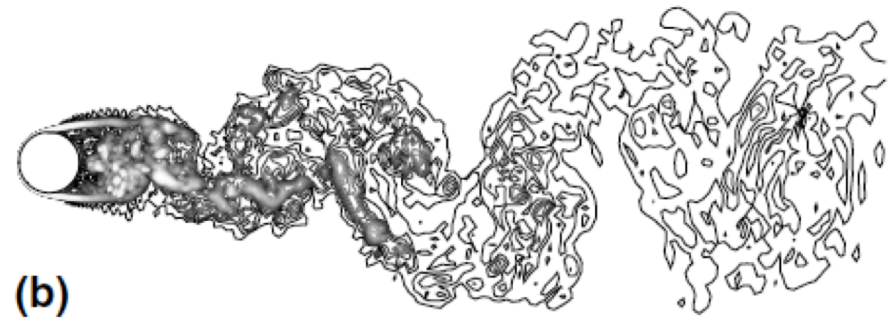
Run1, COM4-sk (DSM); Run2, COM4-sk (no SGS model); Run3, CUDZ (DSM).

Circular cylinder ($Re = 3900$)

Compact scheme with
dynamic Smagorinsky model

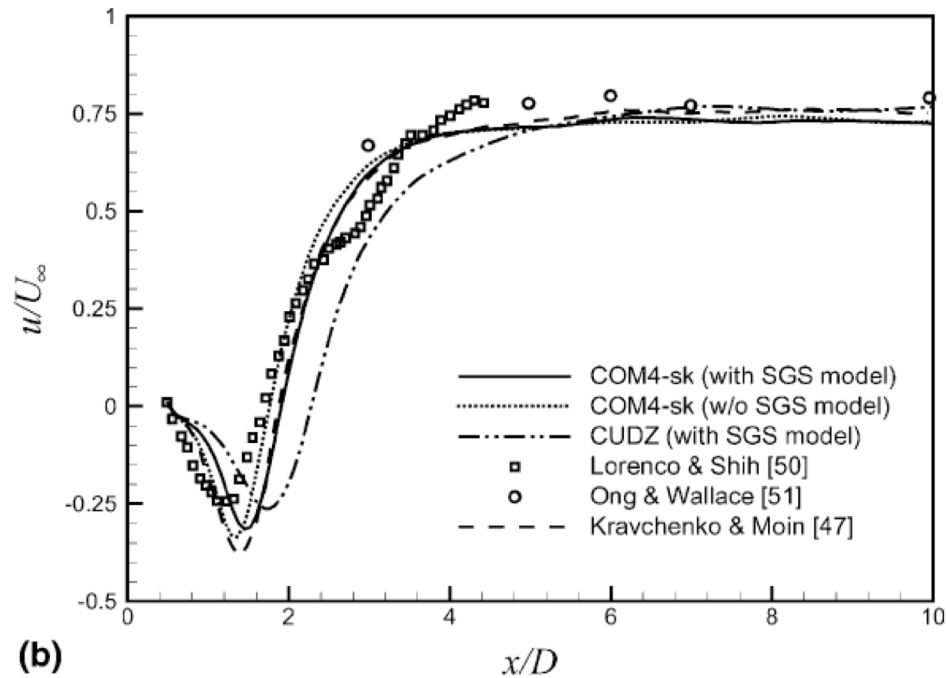
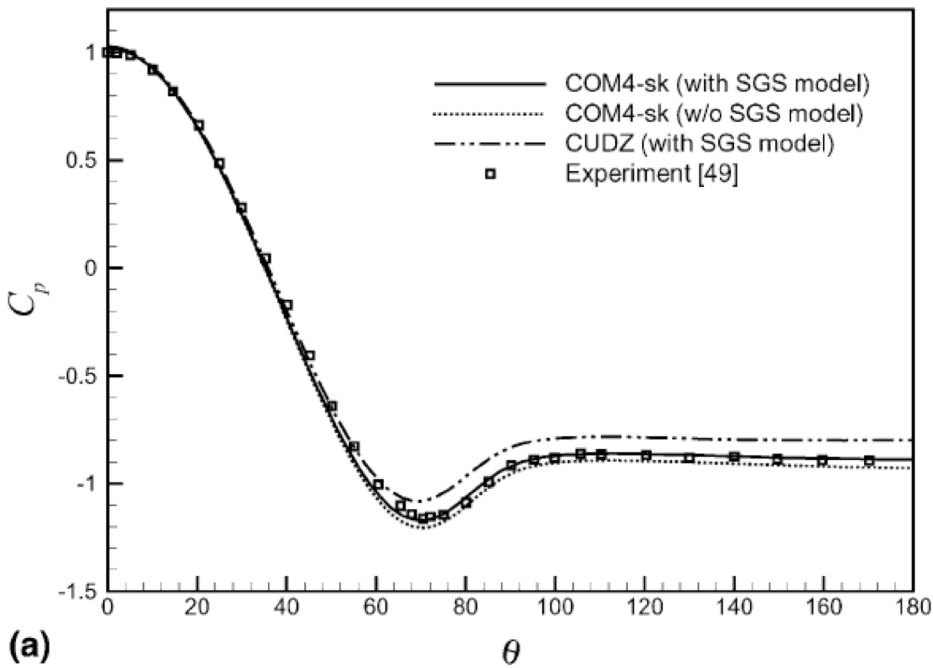


Compact scheme without
Subgrid model

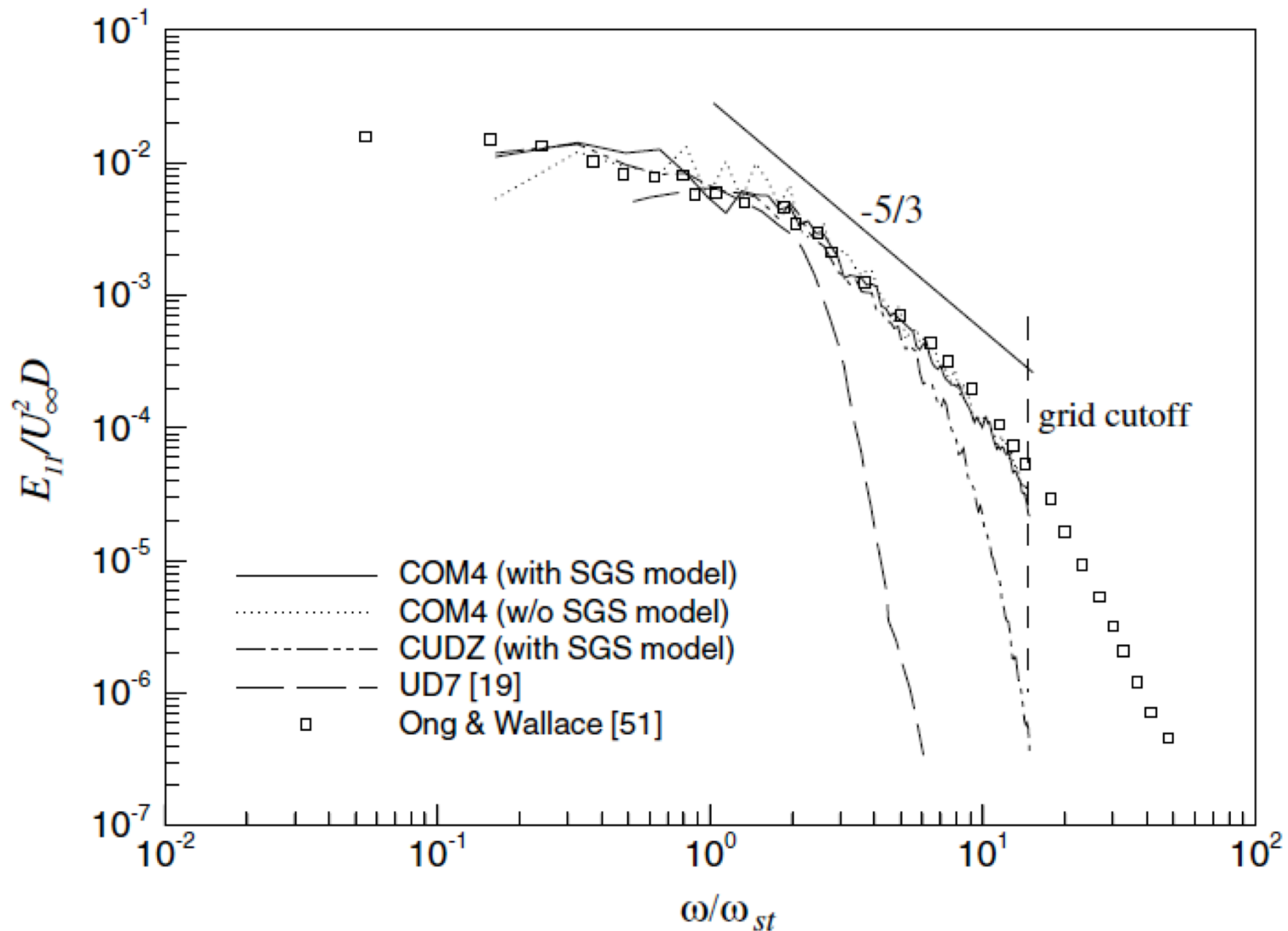


Compact upwind scheme
without subgrid model

Circular cylinder ($Re = 3900$)



Circular cylinder ($Re = 3900$)



3. LES numerics: beyond order of accuracy

Order of accuracy:

- asymptotic behavior of numerical error when $h \longrightarrow 0$
but: h kept as large as possible in practice
- LES solution has lot of kinetic energy at the cutoff
→ convergence of Taylor series expansions not obvious

e.g.

$$\overline{[u(x+r) - u(x)]^p} = O(r^p)?$$

No ! Intermittency effect:



$$\overline{[u(x+r) - u(x)]^p} \sim (\varepsilon r)^{p/3} \Phi(p, r/L, Re)$$

- **order of accuracy** is not sufficient to describe LES numerics
- other features are important, such as
 - conservation of some **invariant quantities** (inviscid flows)
 - ↔ preservation of **Kraichnan's detailed conservation** properties
 - ↔ preservation of **NS symmetries**
- Low-order symmetry-preserving schemes may perform better than higher-order usual schemes
- **Numerical scheme neutrality** is physical-quantity-dependent
- No full-symmetry-group-preserving method known for Navier-Stokes !
- **Partial error cancellation** between discretization error and SGS error may occur

(Lie-group structure)

Time shift

$$(t, x, u, p) \longrightarrow (t + a, x, u, p)$$

Pressure shift

$$(t, x, u, p) \longrightarrow (t, x, u, p + \xi(t))$$

Rotations

$$(t, x, u, p) \longrightarrow (t, \mathcal{R}x, \mathcal{R}u, p)$$

Generalized Galilean Transformations

$$(t, x, u, p) \longrightarrow (t, x + \alpha(t), u + \dot{\alpha}(t), p - \rho x \ddot{\alpha}(t))$$

Rescaling (I)

$$(t, x, u, p) \longrightarrow (a^2 t, ax, u/a, p/a^2)$$

Rescaling (II)

$$(t, x, u, p, \mathbf{v}) \longrightarrow (t, ax, au, a^2 p, a^2 \mathbf{v})$$

Mirror symmetry

$$(t, x, u, p) \longrightarrow (t, \Lambda x, \Lambda u, p)$$

2D Material frame indifference (2D flow only)

$$(t, x, u, p) \longrightarrow (t, \mathcal{R}_{2D}(\omega, t)x, \mathcal{R}_{2D}(\omega, t)u, p - 3\omega\psi + \omega^2|x|^2/2)$$

Time reversal (inviscid flow only)

$$(t, x, u, p) \longrightarrow (-t, x, -u, p)$$

Navier-Stokes momentum equation (incompressible flow)

$$\frac{\partial}{\partial t} u + \underbrace{(u \cdot \nabla) u}_{C(u,u)} + \nabla p - \frac{1}{Re} \underbrace{\nabla \cdot \nabla u}_{\Delta u} = 0$$

Skew-symmetry of pressure & convection terms

$$(u \cdot \nabla)^* = -(u \cdot \nabla), \quad \nabla^* = -\nabla$$

$$(\nabla p, u) = -(p, \nabla \cdot u), \quad \underbrace{((u \cdot \nabla) v, w)}_{(C(u,v), w)} = - \underbrace{(v, (u \cdot \nabla) w)}_{(v, C(u, w))}$$

Global kinetic energy equation

$$\frac{d}{dt}(u, u) = -\frac{2}{Re}(\nabla u, \nabla u) \leq 0$$

Global vorticity equation

$$\frac{d}{dt} \frac{1}{2}(\omega, \omega) = -\frac{1}{Re}|\Delta u|^2 - (C(u, u), \Delta u)$$

Invariant quantities (inviscid flows/Kraichnan's detailed conservation)

- **2D: kinetic energy, vorticity**
- **3D: kinetic energy, helicity**

Idea: numerical schemes which preserve invariants/skew-symmetry

Momentum equation with filtered convection term

$$\frac{\partial}{\partial t} \bar{u} + C_n(\bar{u}, \bar{u}) + \nabla \bar{p} - \frac{1}{Re} \Delta \bar{u} = 0$$

Increasing-order filtered skew-symmetry-preserving nonlinear terms

$$\bar{\phi} = \tilde{\phi} + \phi'$$

Explicit symmetric (self-adjoint) filter

$$C_2(u, v) = \widetilde{C(\tilde{u}, \tilde{v})}$$

$$C_4(u, v) = C(\tilde{u}, \tilde{v}) + \widetilde{C(\tilde{u}, v')} + \widetilde{C(u', \tilde{v})}$$

$$C_6(u, v) = C(\tilde{u}, \tilde{v}) + C(\tilde{u}, v') + C(u', \tilde{v}) + \widetilde{C(u', v')}$$

Links with Adams' Approximate Deconvolution Method

$$u^* = Q \star \bar{u} = u + O(h^n) \simeq u$$

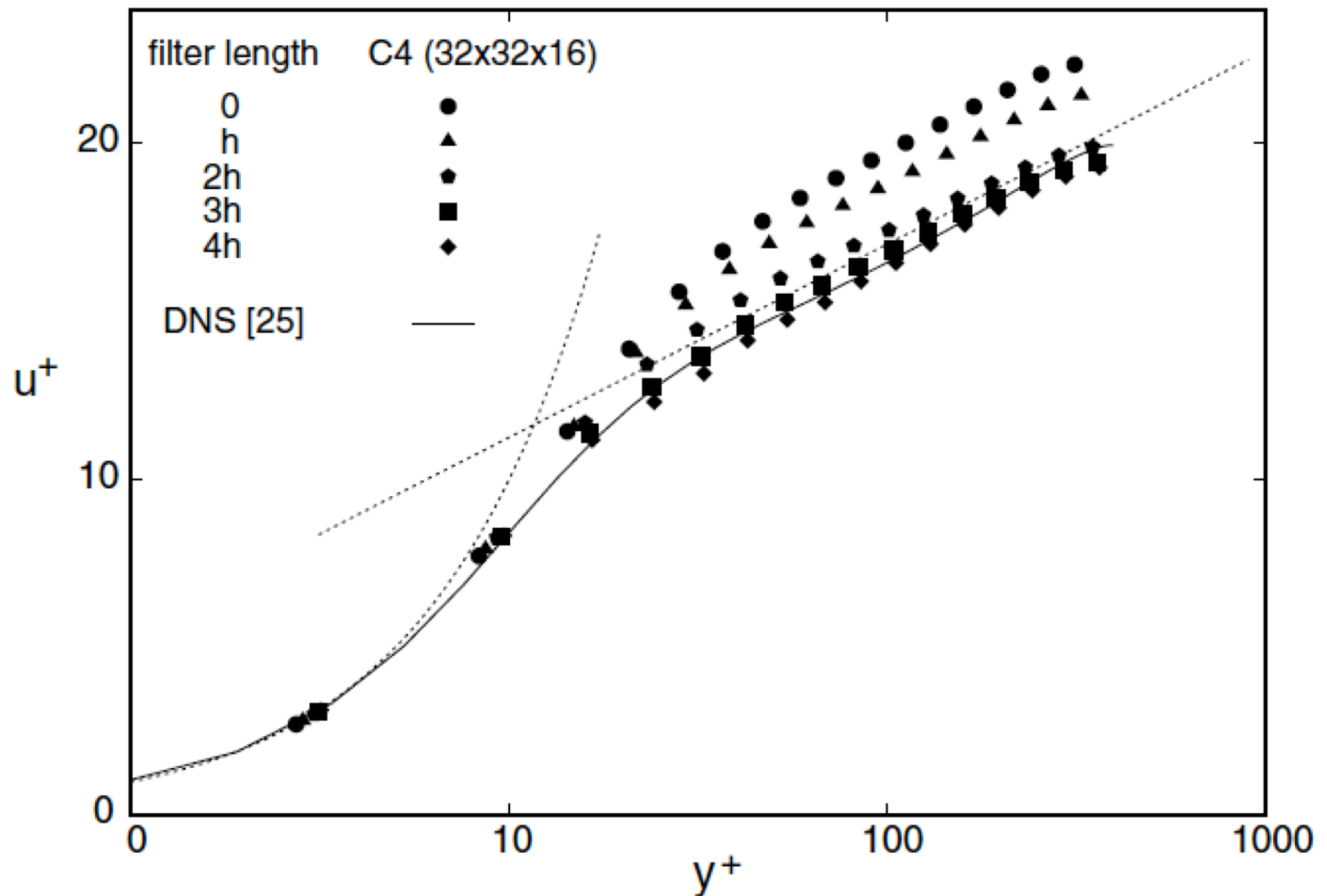
$$\frac{\partial}{\partial t} \bar{u} + \overline{C(u^*, u^*)} + \nabla \bar{p} - \frac{1}{Re} \Delta \bar{u} = 0$$

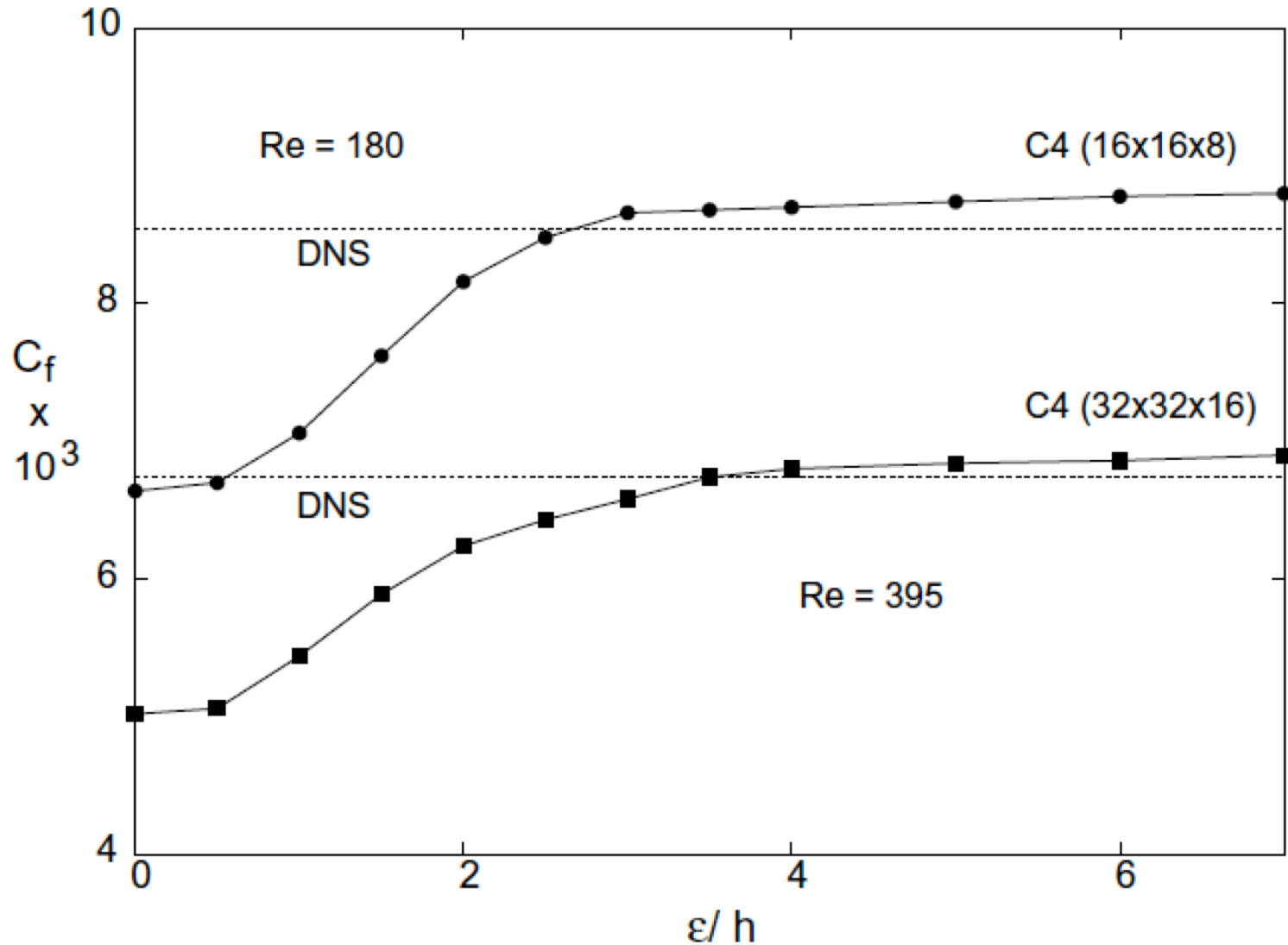
$$\overline{C(Q \star \bar{u}, Q \star \bar{u})} = \widetilde{C(\tilde{u}, \tilde{u})} + O(h^2)$$

Links with tensor diffusivity model

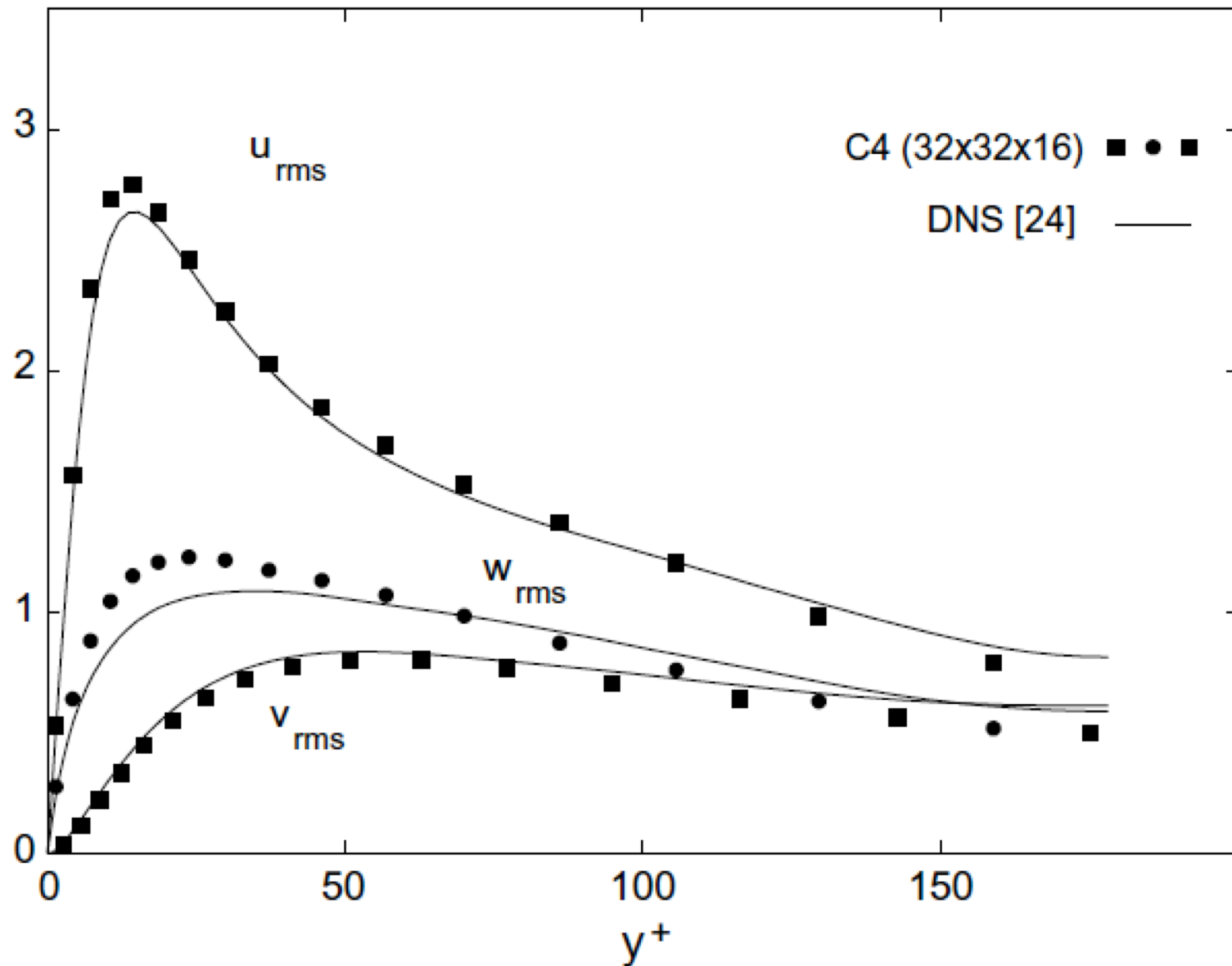
$$C_n(\bar{u}, \bar{u}) = C(\bar{u}, \bar{u}) - \frac{1}{12} \delta^2 \operatorname{div}(\nabla \bar{u} \cdot \nabla \bar{u}) + O(h^4, h^n)$$

Verstappen's method: plane channel flow





Verstappen's method: plane channel flow



Conservative variable formulation

$$\frac{\partial \rho}{\partial t} + \frac{\partial \rho u_i}{\partial x_i} = 0$$

$$\frac{\partial \rho u_i}{\partial t} + \frac{\partial \rho u_i u_j}{\partial x_j} + \frac{\partial p}{\partial x_i} = \frac{\partial \tau_{ij}}{\partial x_j}$$

$$\frac{\partial E}{\partial t} + \frac{\partial E u_j}{\partial x_j} + \frac{\partial p u_i}{\partial x_i} = \frac{\partial \tau_{ij} u_i}{\partial x_j} + \frac{\partial}{\partial x_i} \left(\kappa \frac{\partial T}{\partial x_i} \right)$$

$$\tau_{ij} = \mu \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) - \frac{2}{3} \mu \frac{\partial u_k}{\partial x_k} \delta_{ij}, \quad E = \rho \frac{u_i u_i}{2} + \rho C_v T$$

- **Problems:**
 - no scalar invariant directly appearing
 - splitting of fluctuations into different physical modes
 - **Kovasznay**: acoustic, vortical & entropy modes
 - **Helmholtz**: solenoidal & dilatational modes
 - transfers between modes
- **Possible solution**: looking at 2nd law of thermodynamics/entropy

$$\begin{aligned} \frac{\partial \rho s}{\partial t} + \frac{\partial \rho s u_j}{\partial x_j} &= \frac{1}{T} \left[\tau_{ij} \frac{\partial u_i}{\partial x_j} + \frac{\partial}{\partial x_i} \left(\kappa \frac{\partial T}{\partial x_i} \right) \right] \\ &= \frac{\partial}{\partial x_i} \left(\frac{\kappa}{T} \frac{\partial T}{\partial x_i} \right) + \frac{\tau_{ij}}{T} \frac{\partial u_i}{\partial x_j} + \frac{\kappa}{T^2} \frac{\partial T}{\partial x_i} \frac{\partial T}{\partial x_i} \end{aligned}$$

$$\frac{\partial \rho s^2}{\partial t} + \frac{\partial \rho s^2 u_j}{\partial x_j} = \frac{2s}{T} \left[\tau_{ij} \frac{\partial u_i}{\partial x_j} + \frac{\partial}{\partial x_i} \left(\kappa \frac{\partial T}{\partial x_i} \right) \right]$$

Global invariants in homogeneous flows:

$$\overline{\rho s}, \quad \overline{\rho s^2}$$

Skew-symmetric splittings

Quadratic nonlinear term splitting (also in incompressible flows)

$$\frac{\partial a u_j}{\partial x_j} \longrightarrow \frac{1}{2} \left(\frac{\partial a u_j}{\partial x_j} + a \frac{\partial u_j}{\partial x_j} + u_j \frac{\partial a}{\partial x_j} \right)$$

Entropy-consistent cubic nonlinear term splitting

$$\frac{\partial \rho a u_j}{\partial x_j} \longrightarrow \frac{1}{2} \left(\frac{\partial \rho a u_j}{\partial x_j} + a \frac{\partial \rho u_j}{\partial x_j} + \rho u_j \frac{\partial a}{\partial x_j} \right)$$

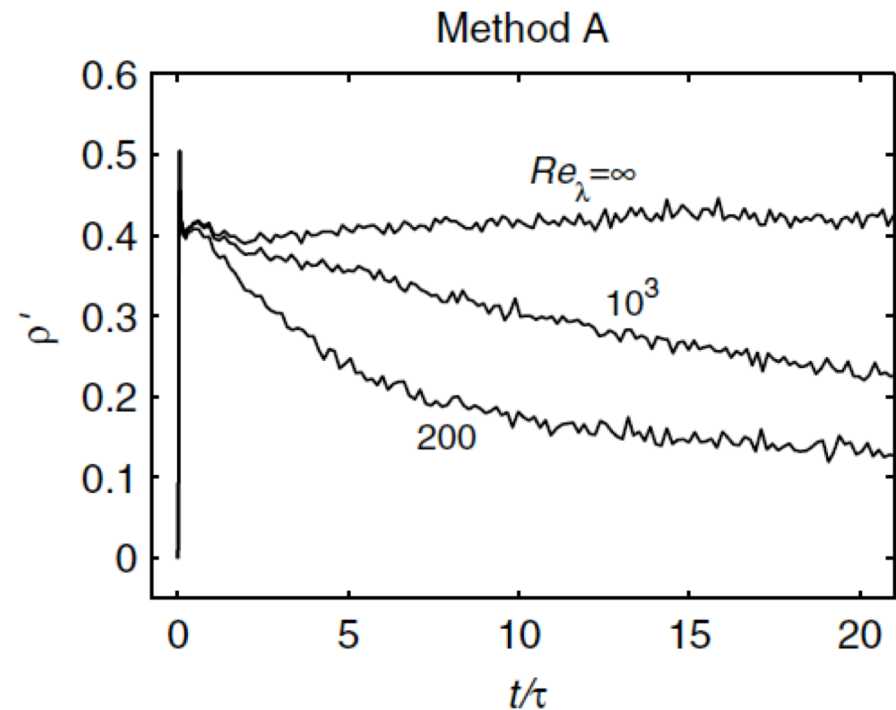
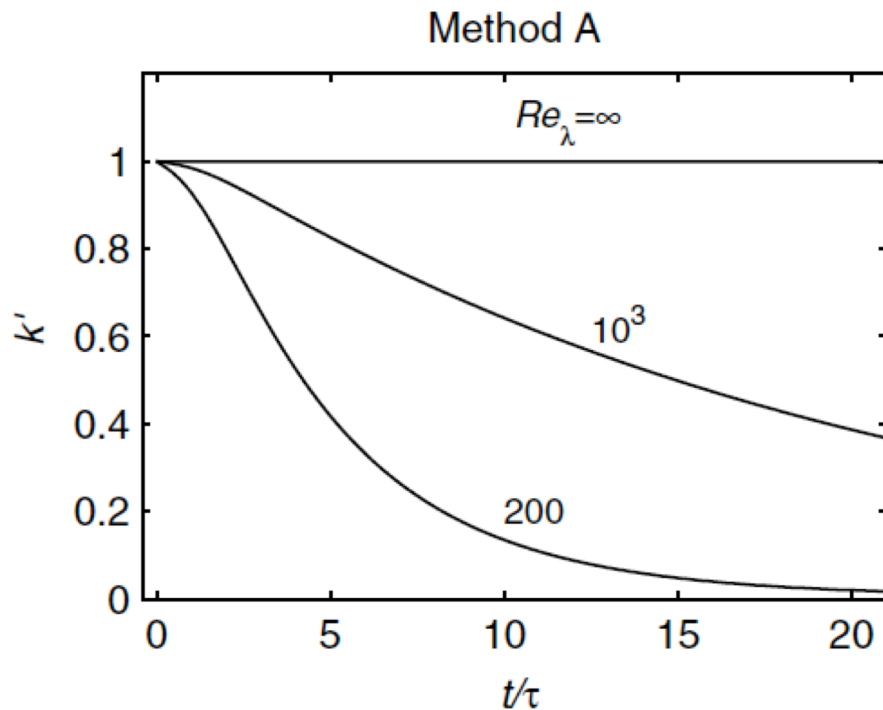
Other cubic nonlinear term splitting

$$\frac{\partial \rho a u_j}{\partial x_j} \longrightarrow \frac{1}{2} \left(\frac{\partial \rho a u_j}{\partial x_j} + \rho a \frac{\partial u_j}{\partial x_j} + u_j \frac{\partial \rho a}{\partial x_j} \right)$$

Honein's entropy-preserving method

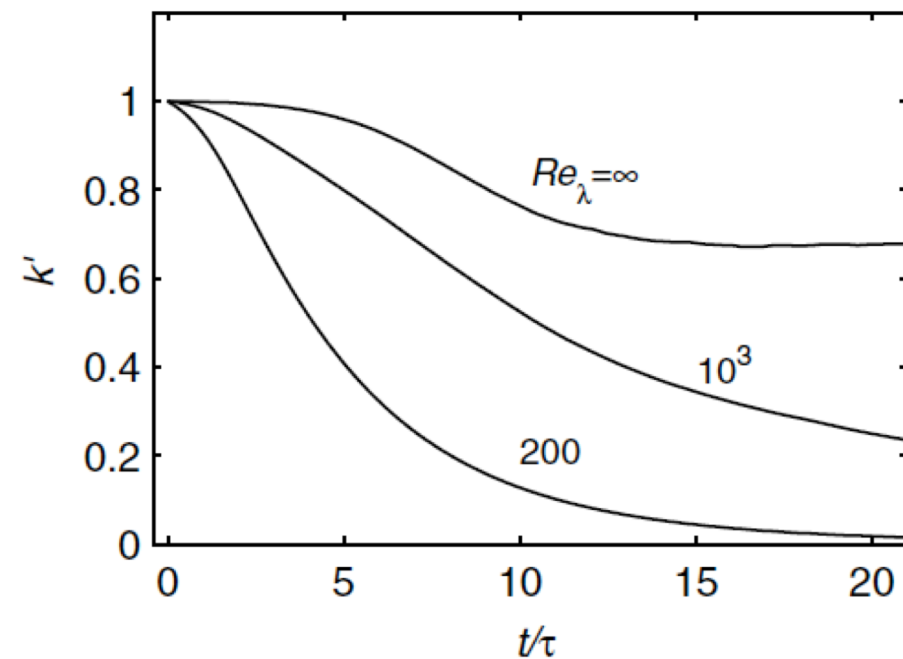
Internal energy entropy-preserving equation

$$\frac{\partial \rho e}{\partial t} + e \left(\gamma - \frac{s}{2C_v} \right) \frac{\partial \rho u_j}{\partial x_j} + \frac{e}{2C_v} \frac{\partial \rho s u_i}{\partial x_i} + \frac{\rho e u_i}{2C_v} \frac{\partial s}{\partial x_i} = \tau_{ij} \frac{\partial u_i}{\partial x_j} + \frac{\partial}{\partial x_i} \left(\kappa \frac{\partial T}{\partial x_i} \right)$$

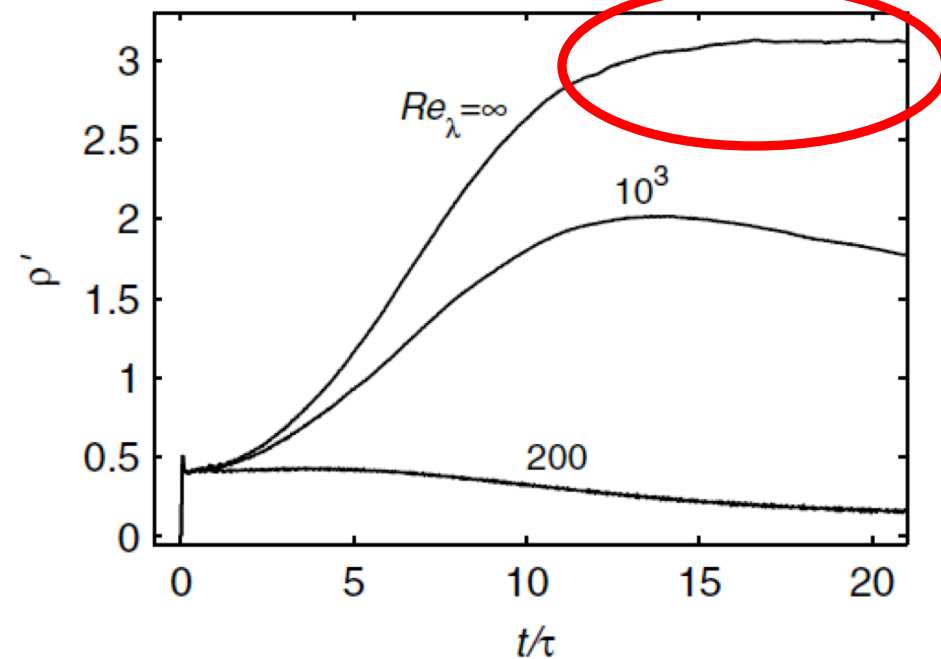


Yee's entropy-splitting method

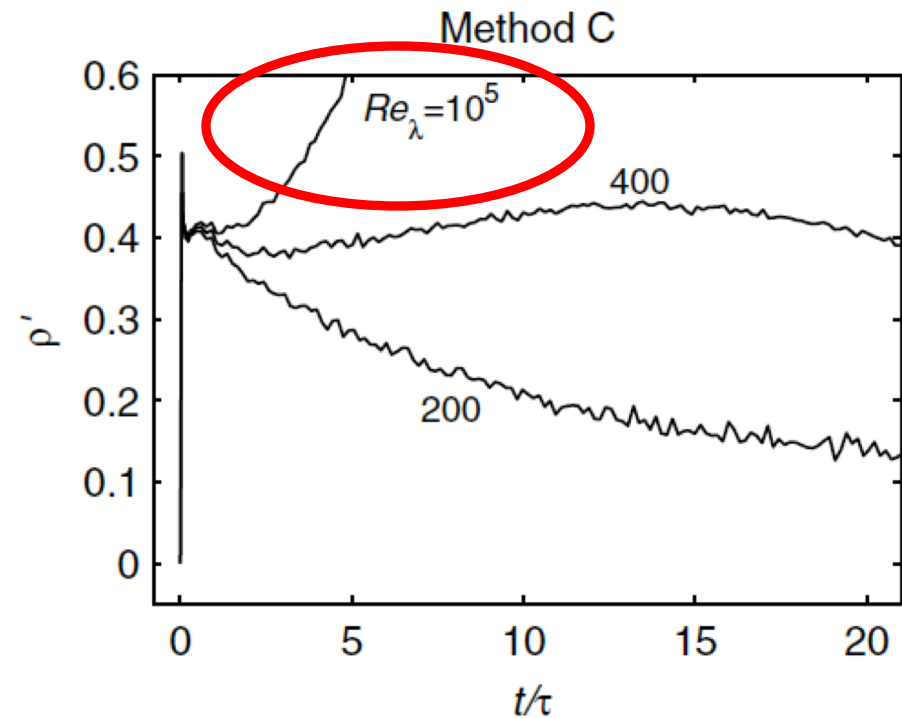
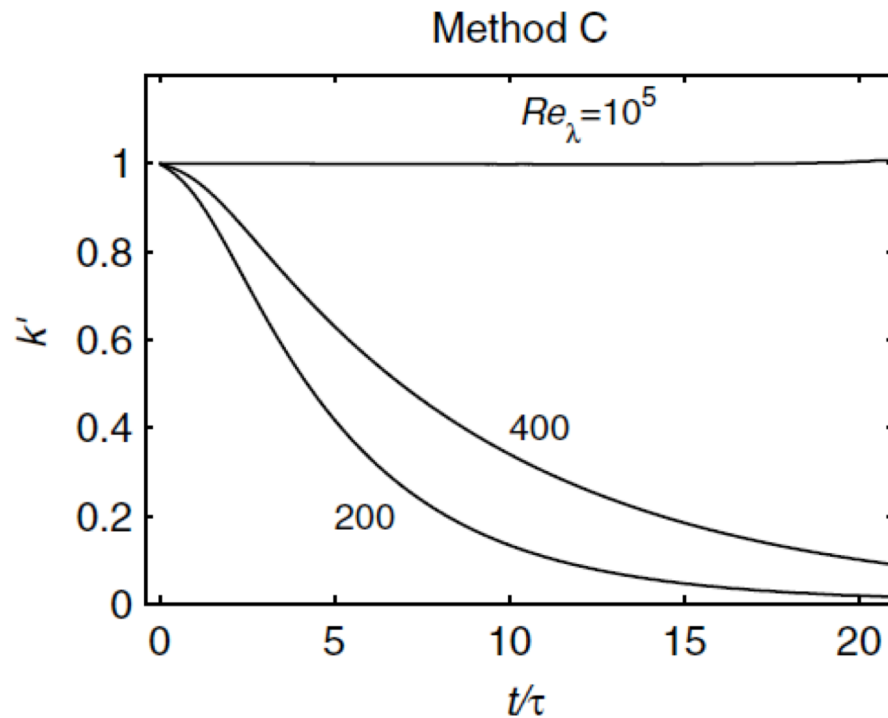
Method B



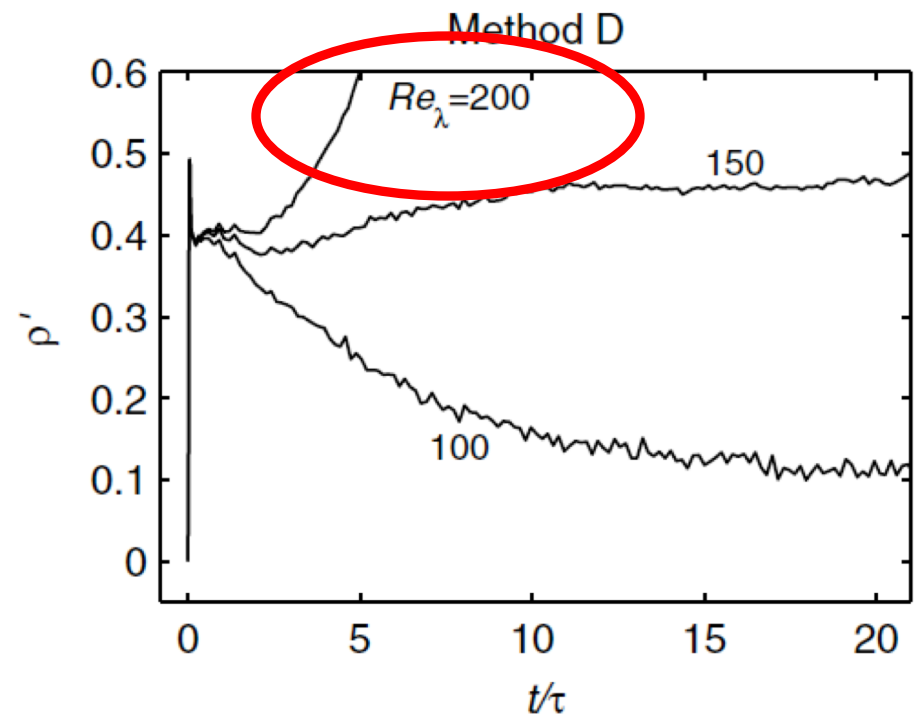
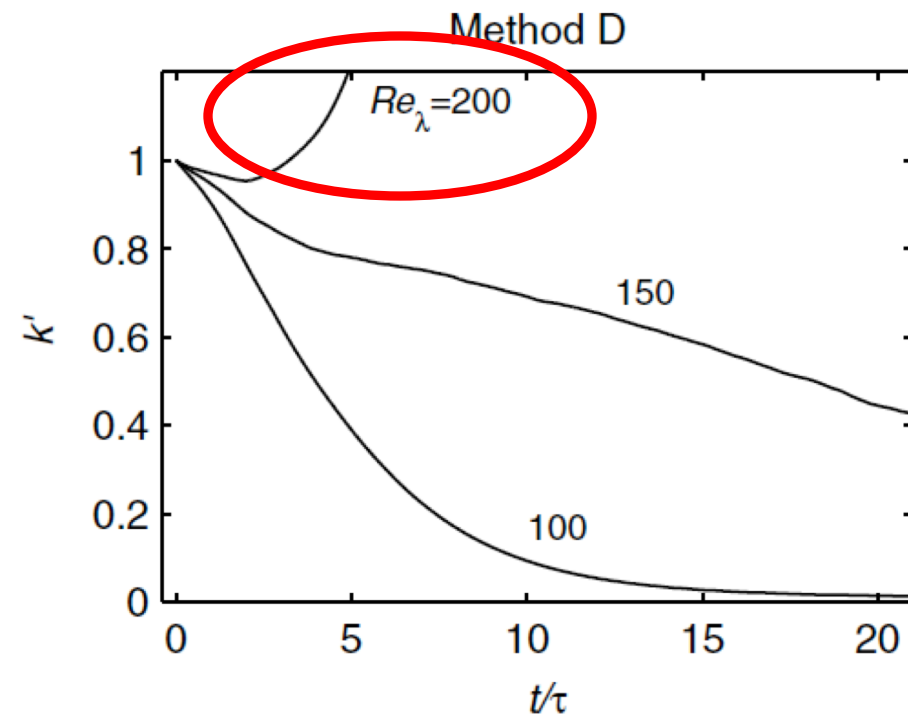
Method B



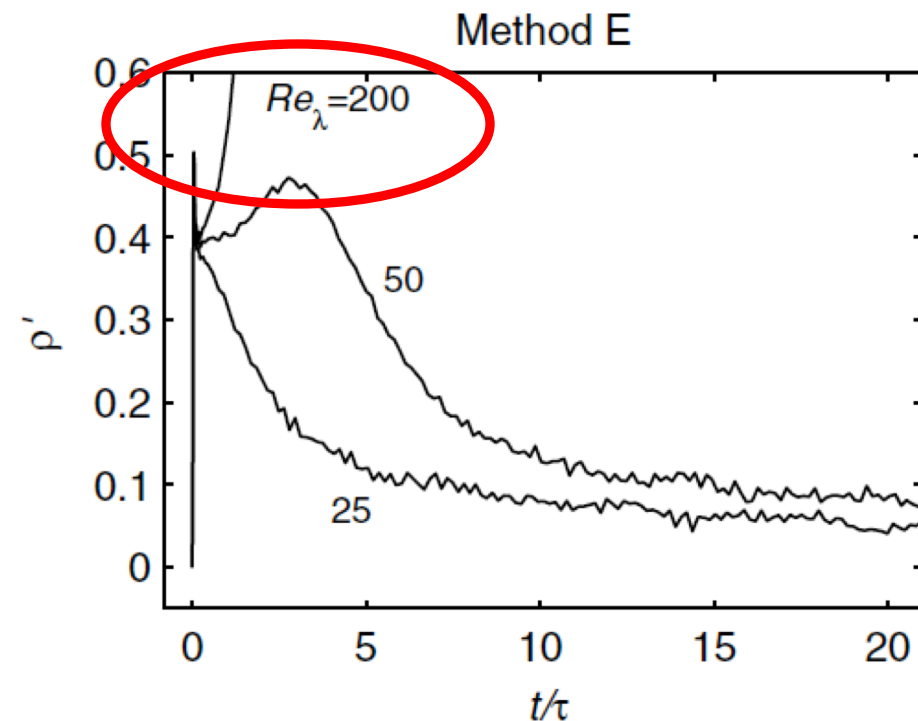
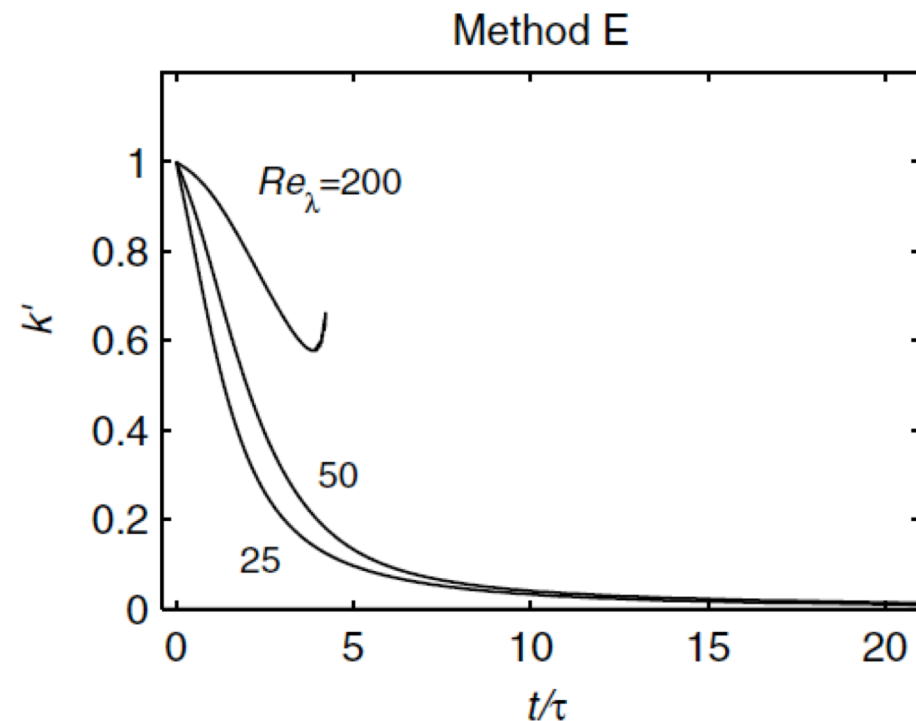
Honein's splitting & usual internal energy equation



Staggered grid method



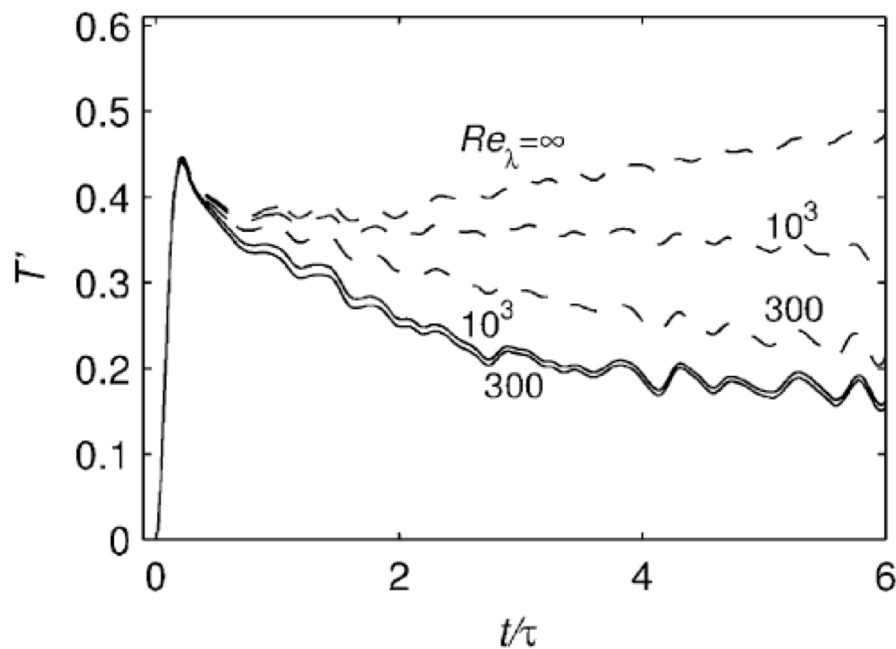
Conservative formulation & bad cubic splitting



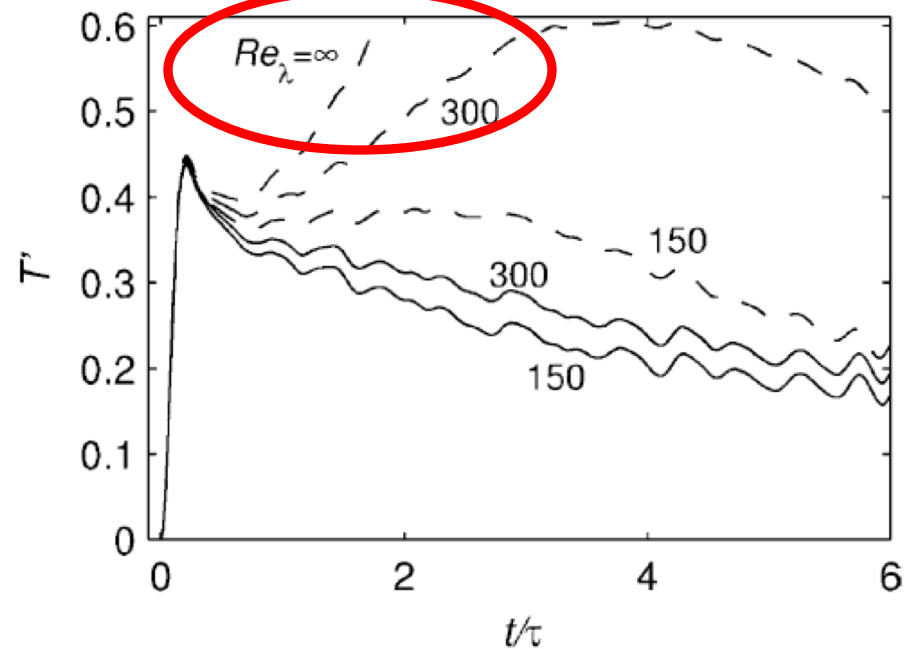
Honein's entropy-preserving method

Yee's entropy-splitting method

Method A

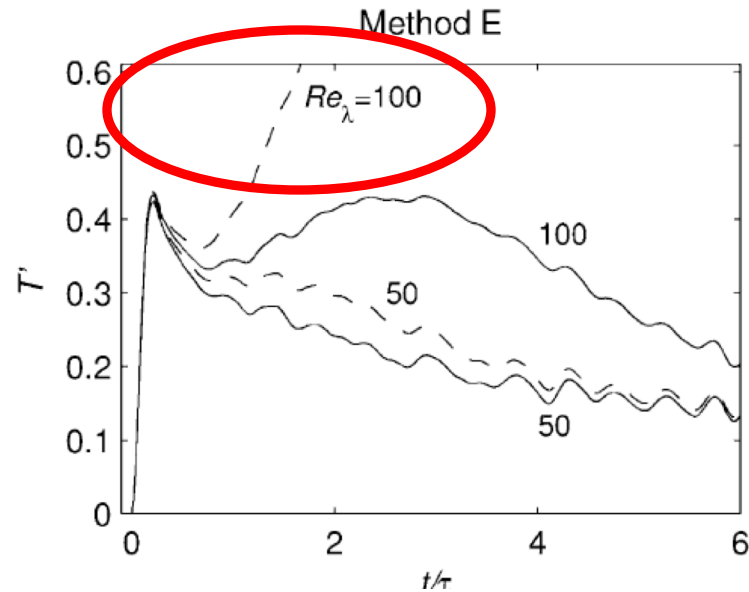


Method B

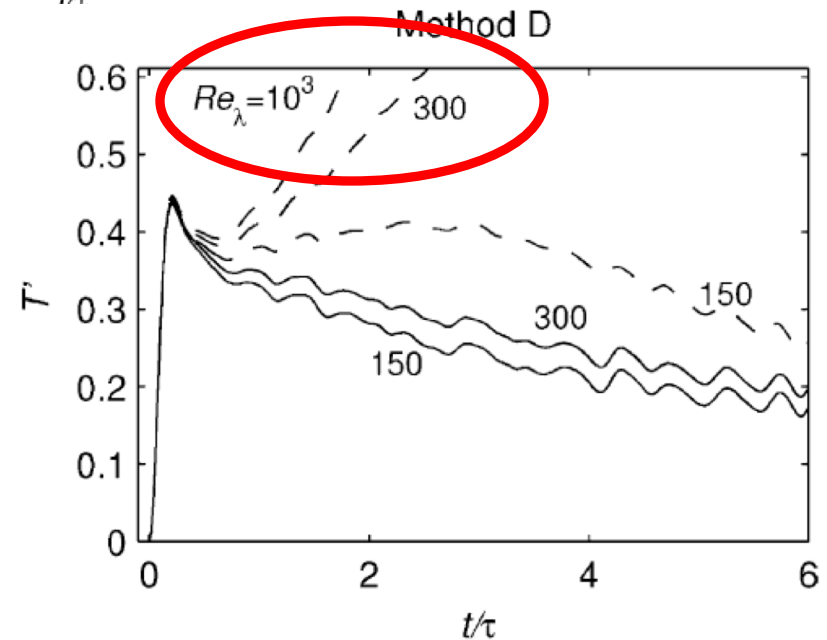
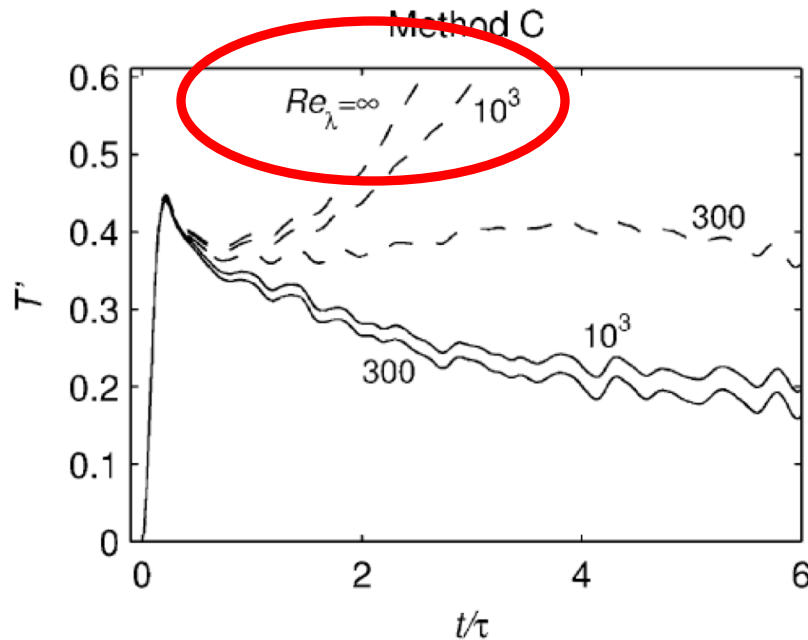


Solid line: LES Dashed: model-free

Honein's method: model-free vs. LES



Conservative formulation
& bad cubic splitting



Honein's splitting & usual internal energy equation

Staggered grid method

How to measure the accuracy of LES results ?

- need for a quantitative error norm
- need for reference data (accuracy $\neq$ sensitivity)
- sensitivity important for robustness evaluation
- several quantities predicted at the same time $\rightarrow$ unique error ?

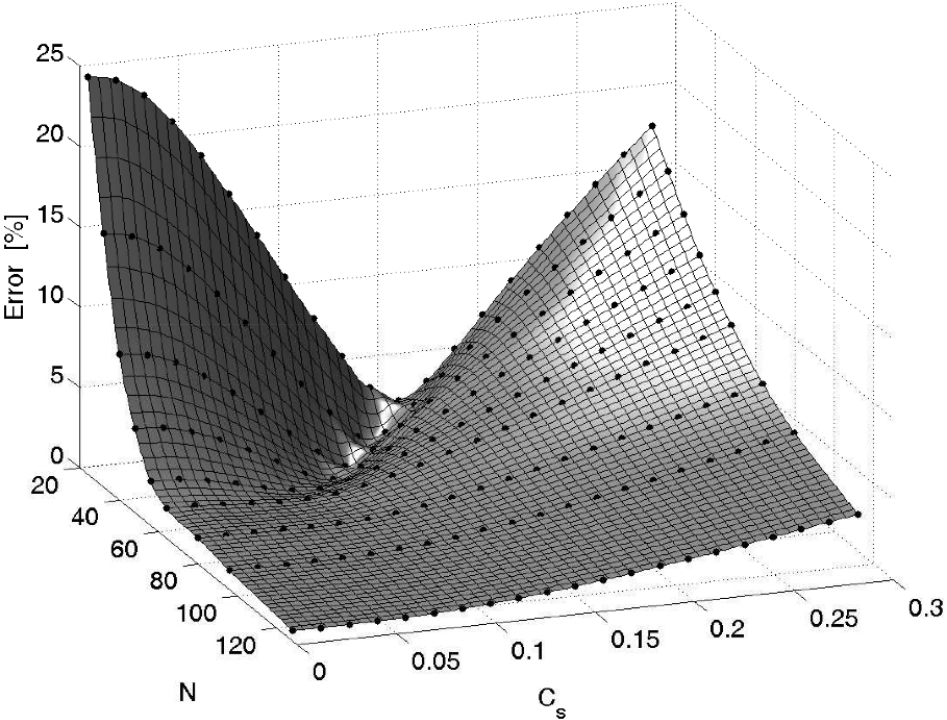
L2-type error norms

$$D_p(N, C_s) = \left[\frac{\int_0^T \left\{ \int_0^{k_c} k^p (E_{LES}(k, t) - E_{DNS}(k, t)) dk \right\}^2 dt}{\int_0^T \left\{ \int_0^{k_c} k^p E_{DNS}(k, t) dk \right\}^2 dt} \right]^{1/2}$$

Physical-quantity-based error norms

$$d_p(N, C_s) = \left[\frac{\int_0^T \int_0^{k_c} k^p (E_{LES}(k, t) - E_{DNS}(k, t))^2 dk dt}{\int_0^T \int_0^{k_c} k^p E_{DNS}^2(k, t) dk dt} \right]^{1/2}$$

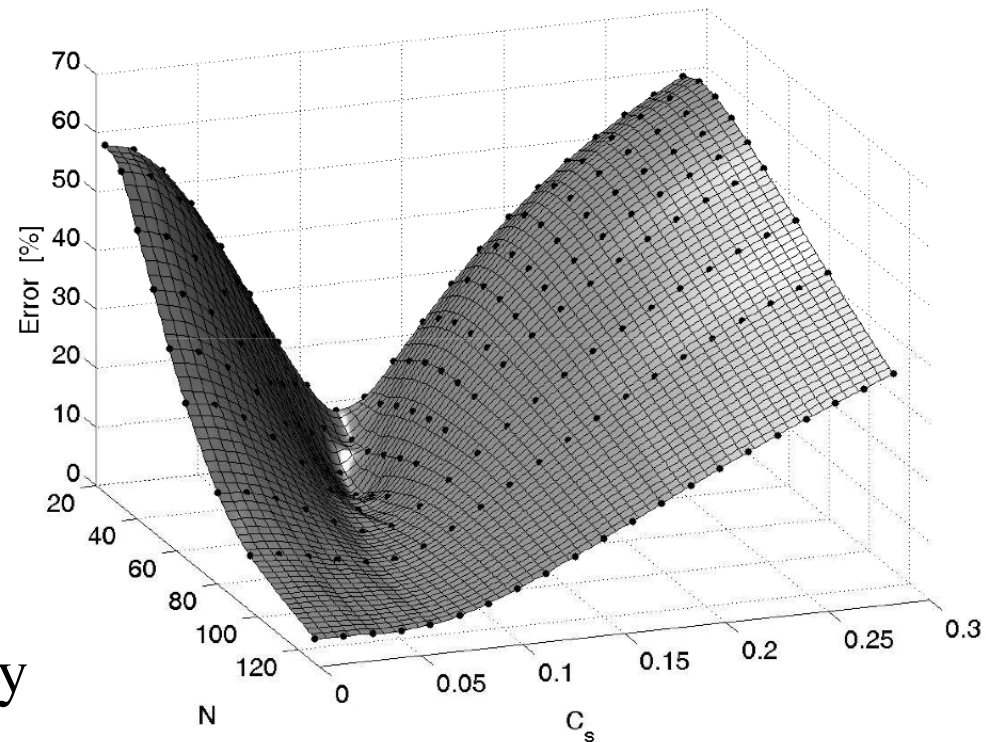
$$L = \int_0^{+\infty} k^{-1} E(k) dk, \quad K = \int_0^{+\infty} E(k) dk, \quad \varepsilon = 2\nu \int_0^{+\infty} k^2 E(k) dk$$



Resolved kinetic energy

Error surfaces

(Smagorinsky model, $Re_\lambda=100$)



Resolved enstrophy

L2-type error norms

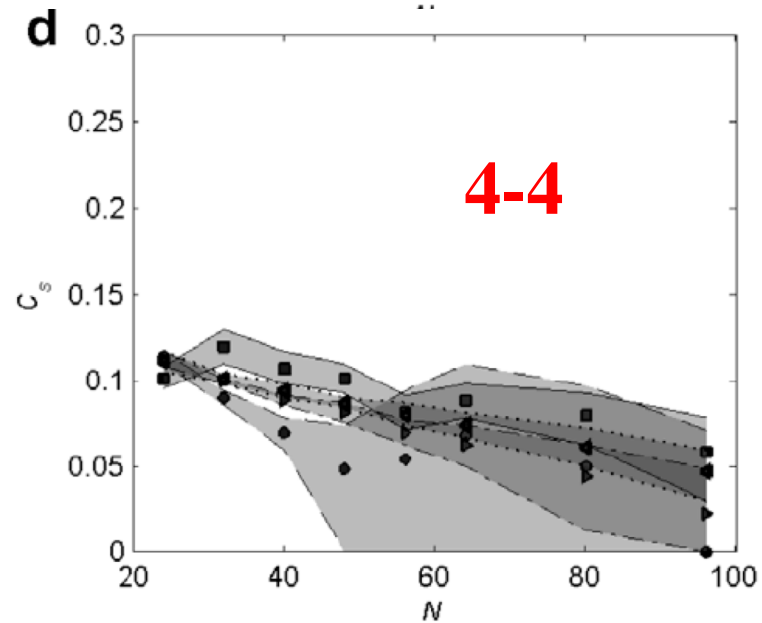
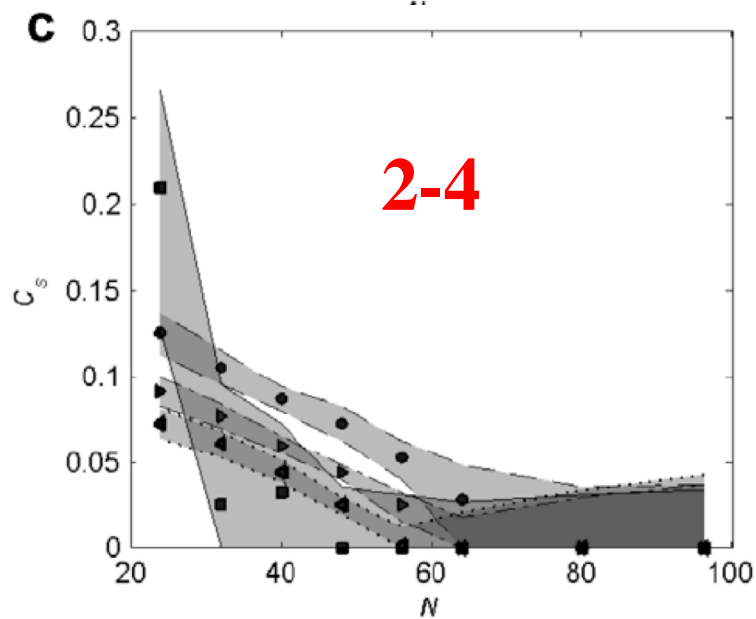
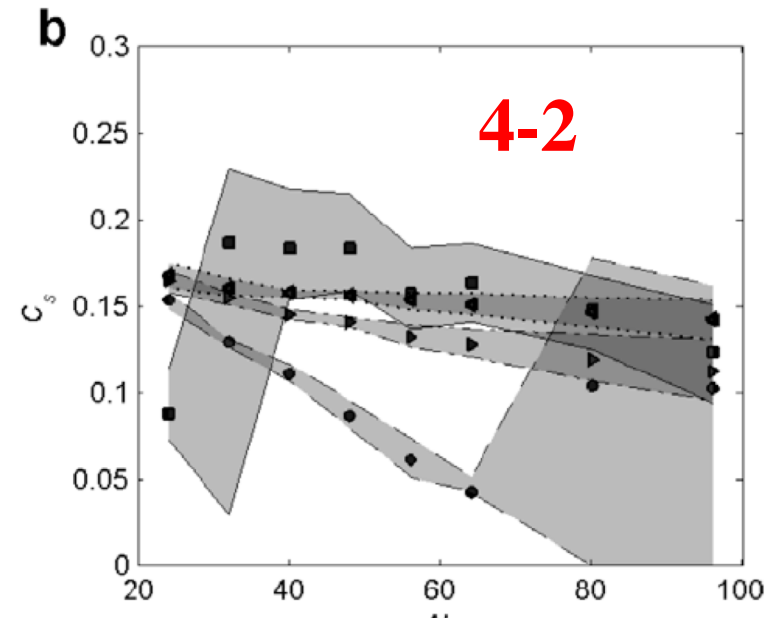
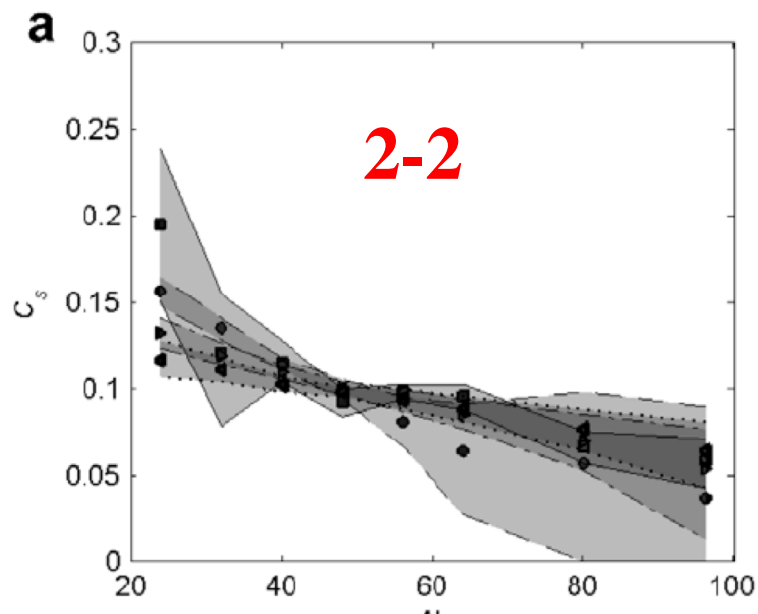
$$\tilde{D}(N, C_s) = \frac{\sum_p \left[D_p(N, C_s) / D_p(N, \hat{C}_s^{(p)}(N)) \right]}{\sum_p \left[1 / D_p(N, \hat{C}_s^{(p)}(N)) \right]}$$

Physical error norms

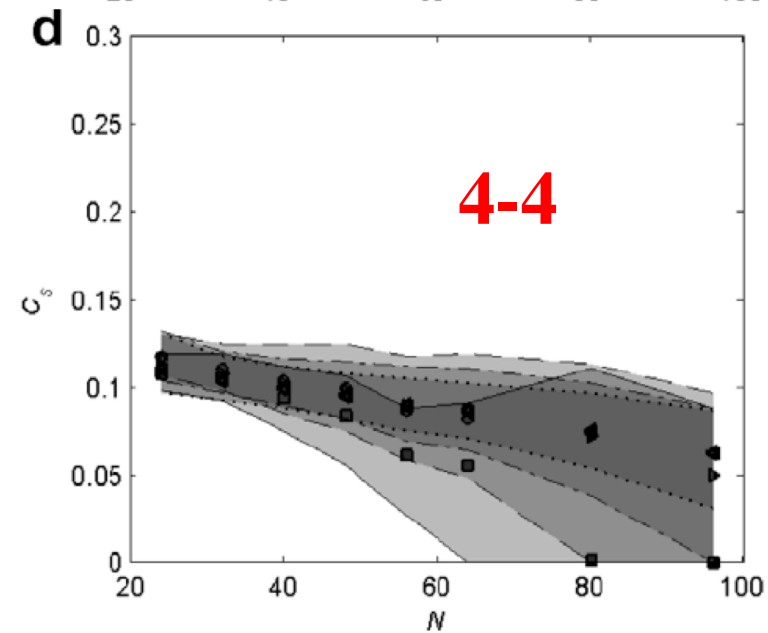
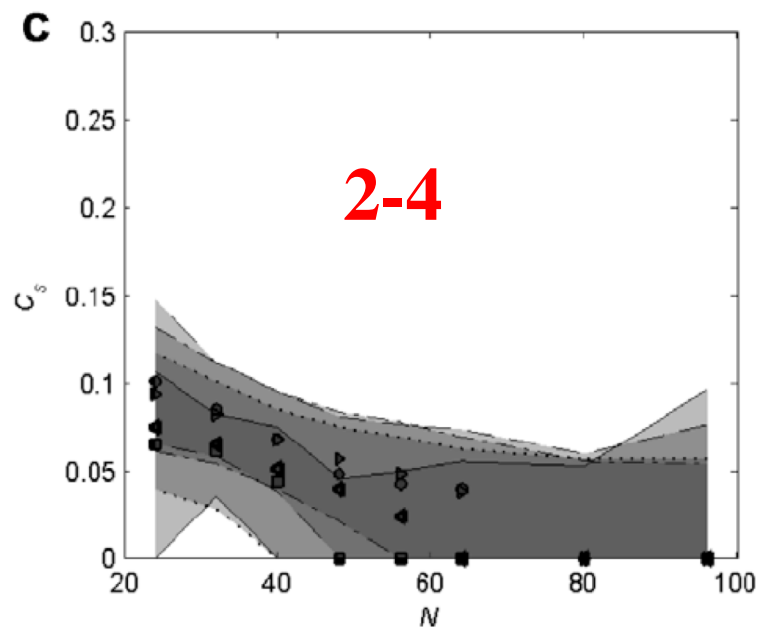
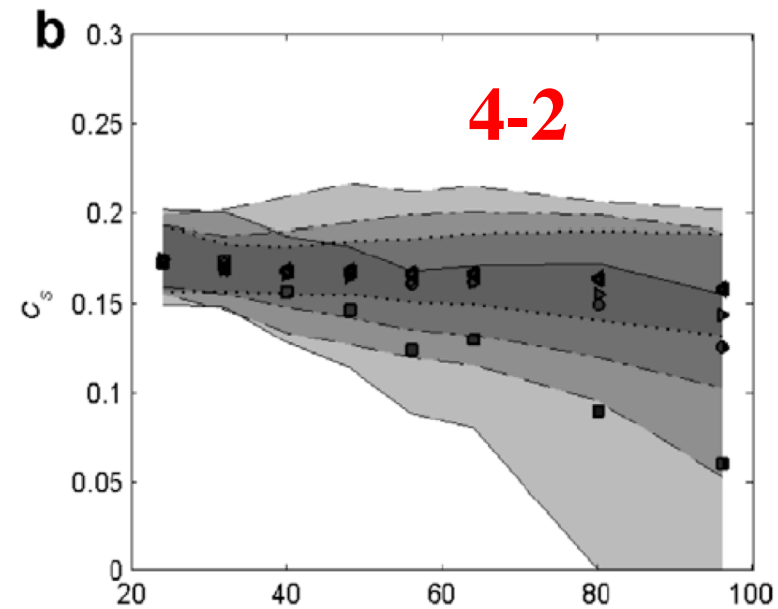
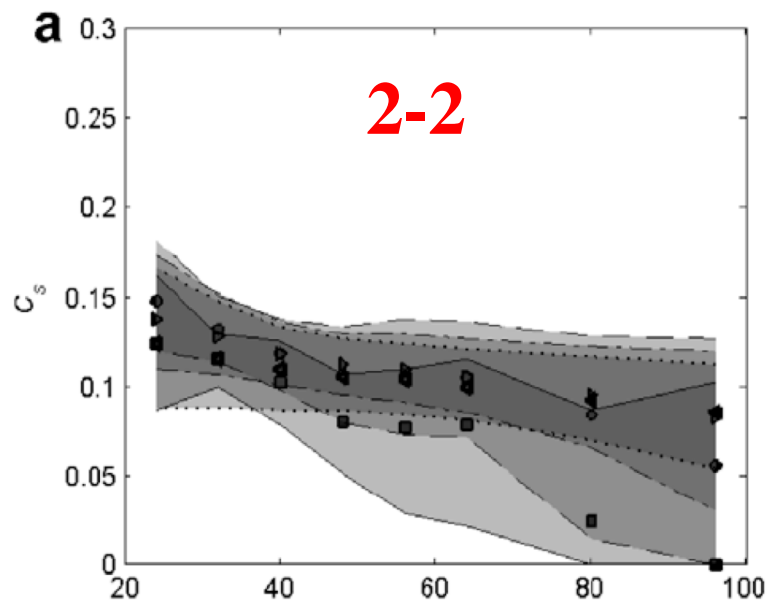
$$\tilde{d}(N, C_s) = \frac{\sum_p \left[d_p(N, C_s) / d_p(N, \hat{C}_s^{(p)}(N)) \right]}{\sum_p \left[1 / d_p(N, \hat{C}_s^{(p)}(N)) \right]}$$

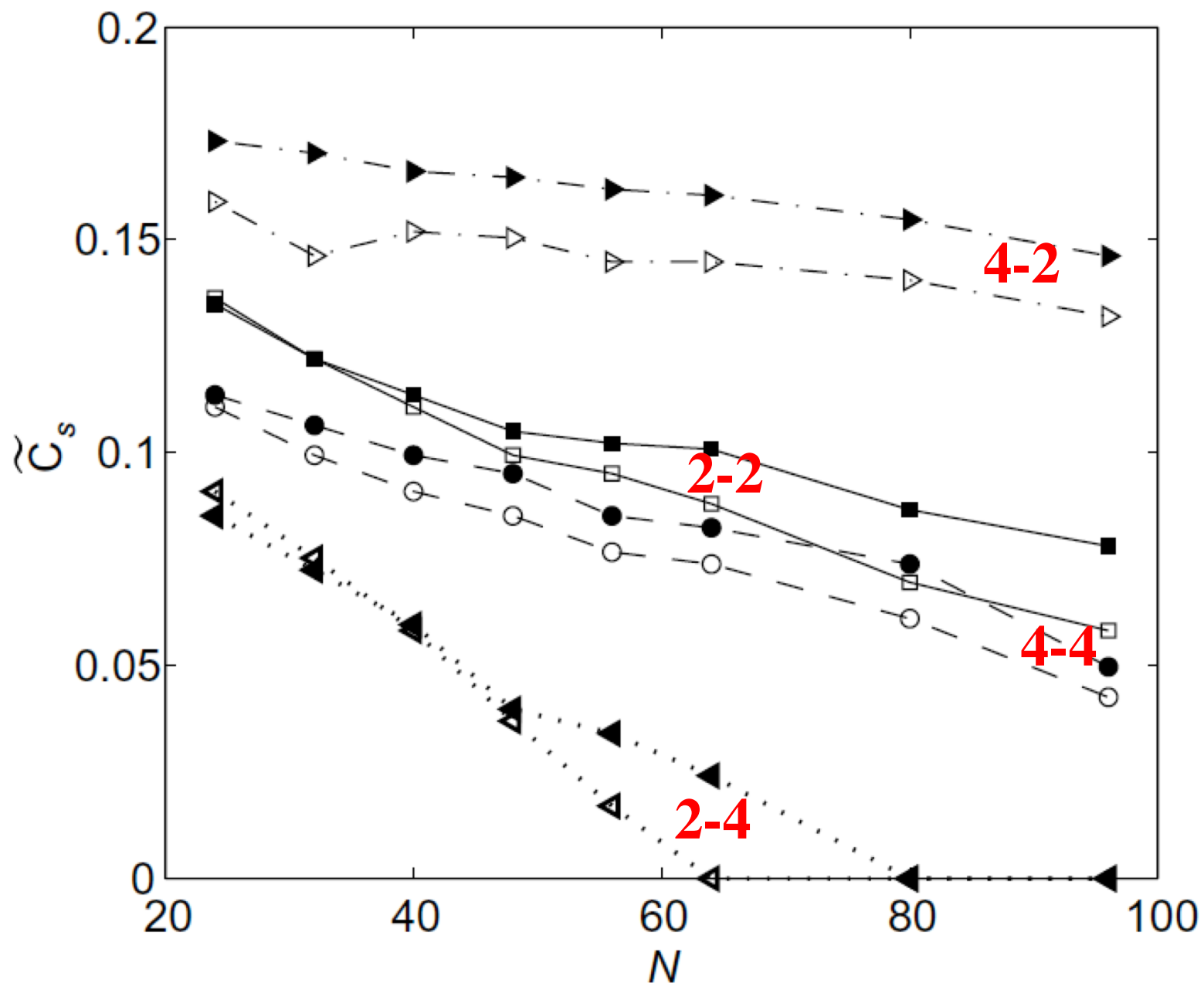
Best LES solution on the same grid (optimal C_s)

Near-optimal LES: L2 error norm

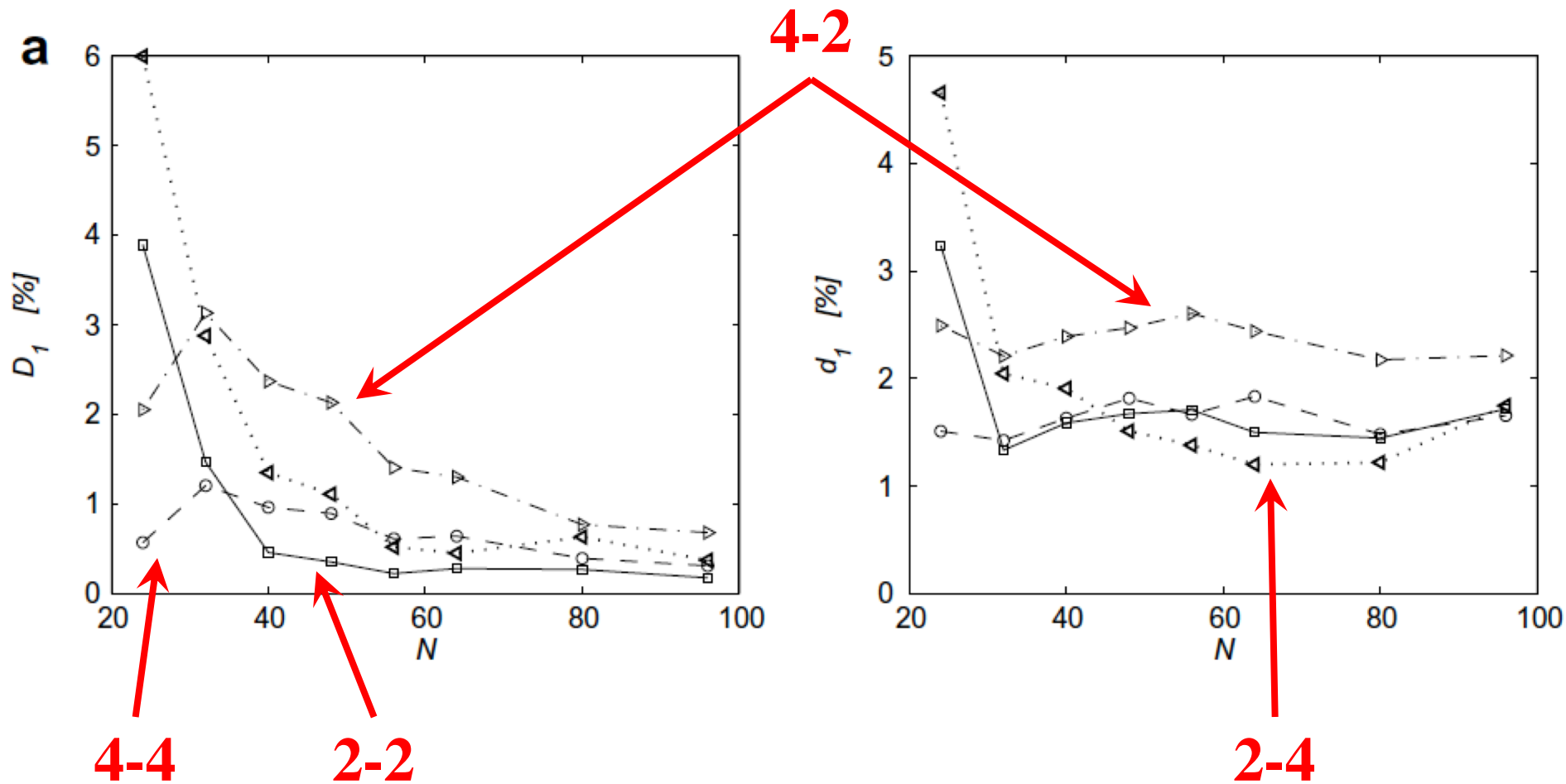


Near-optimal LES: physical error norm



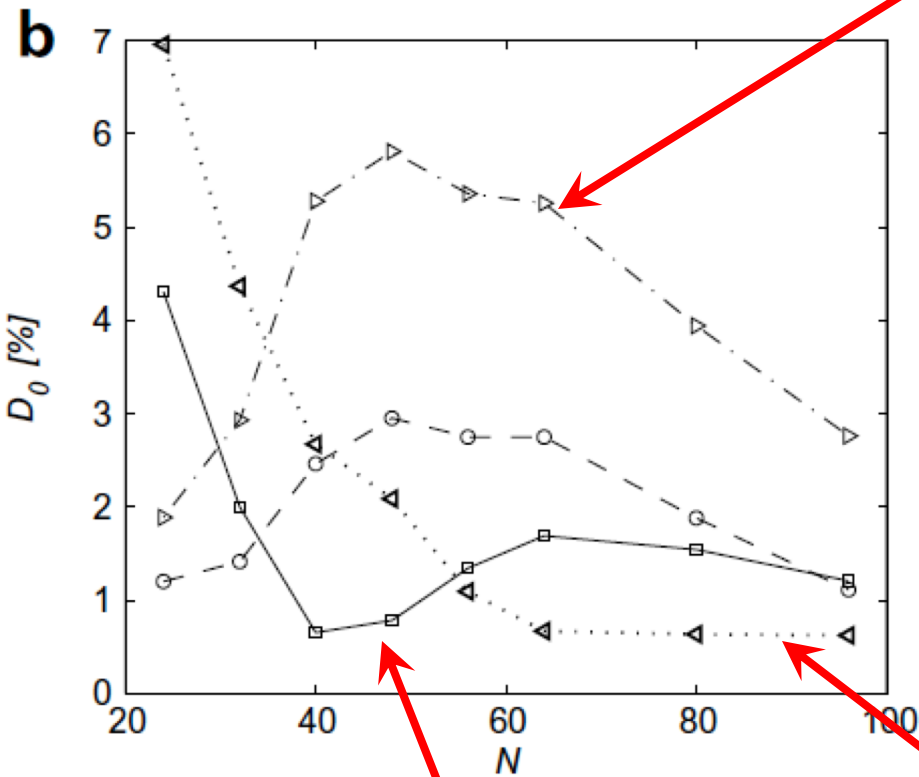


$p=-1$

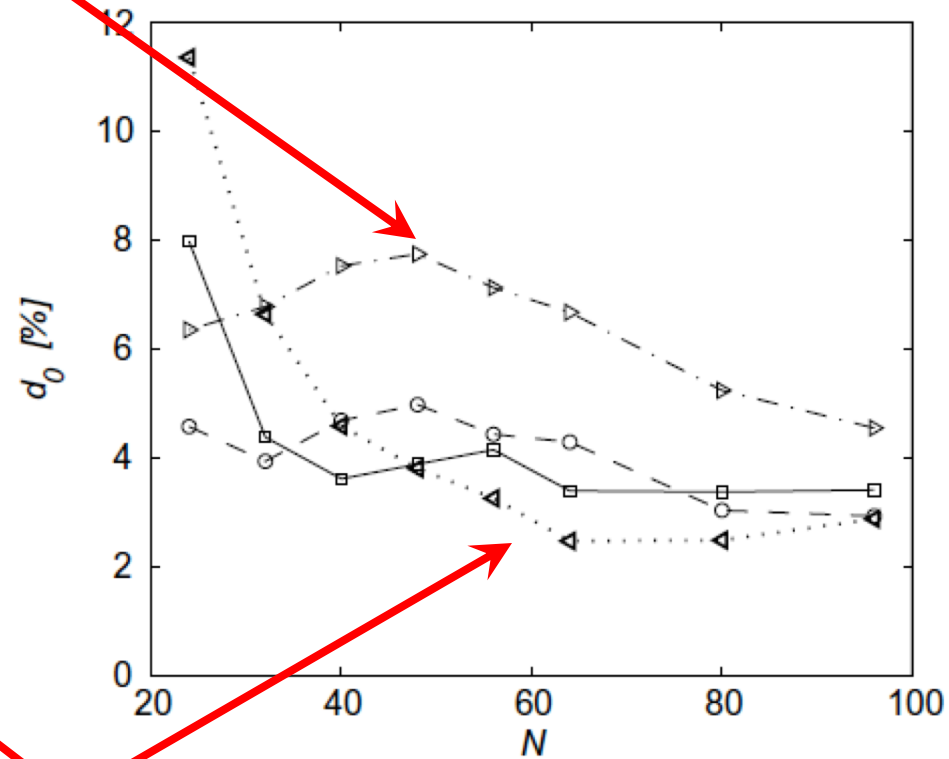


$p=0$

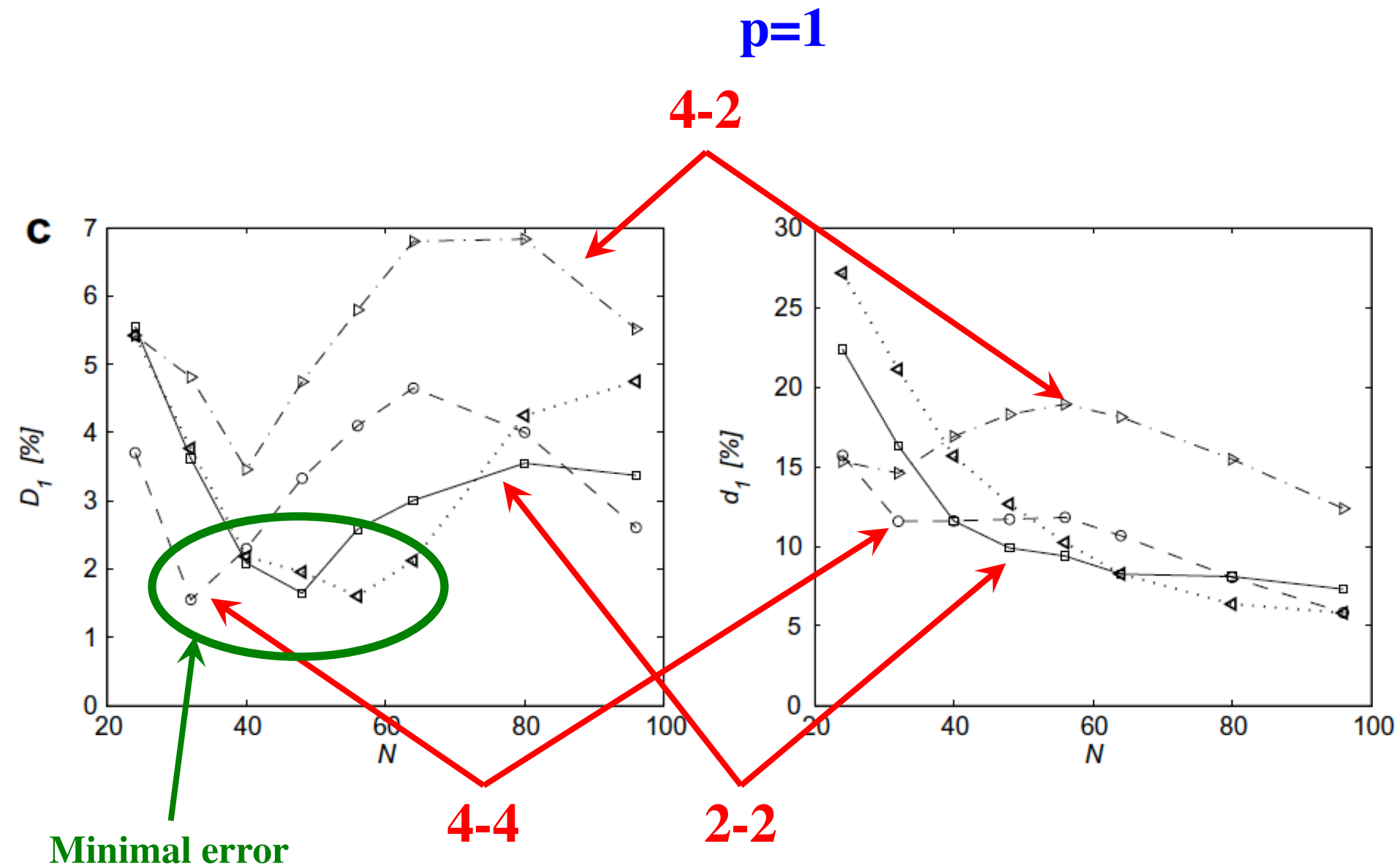
4-2



2-2



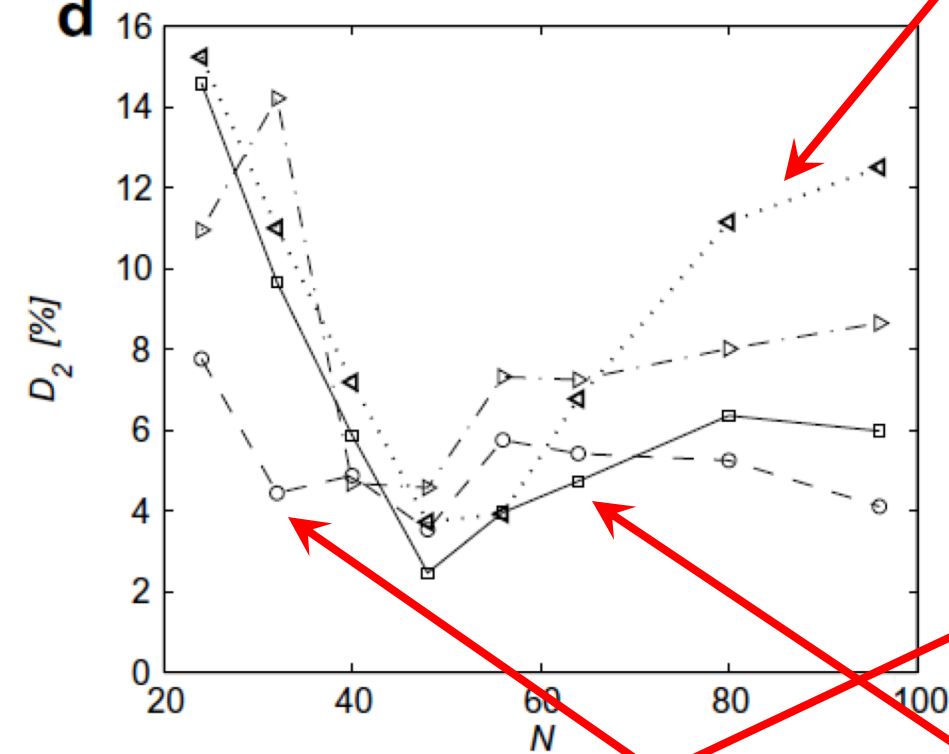
2-4



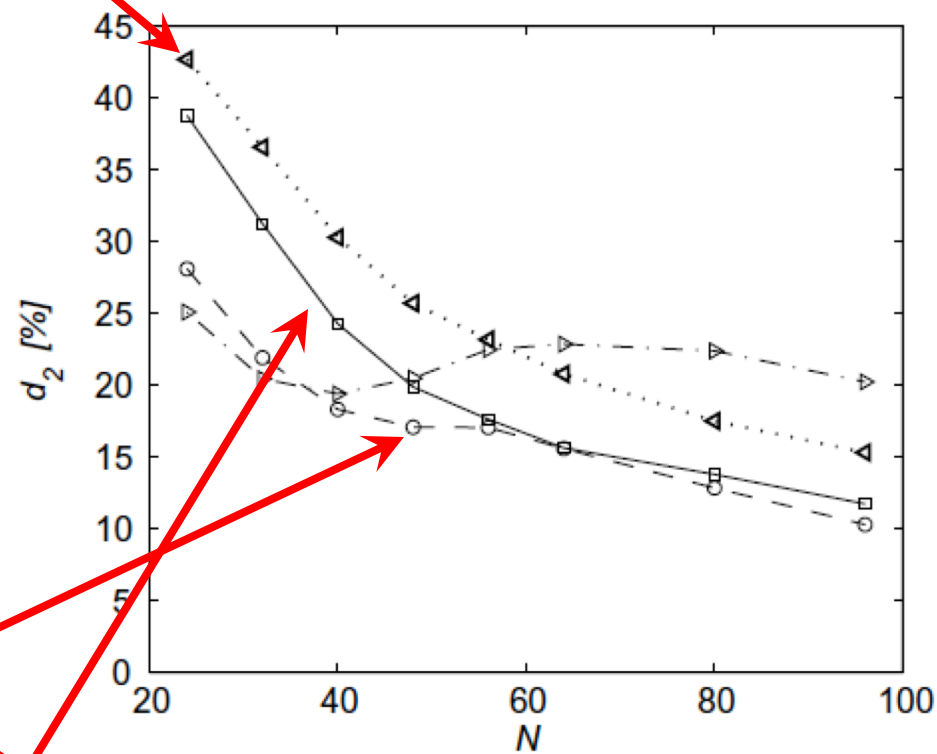
$p=2$

4-2

d



d_2 [%]



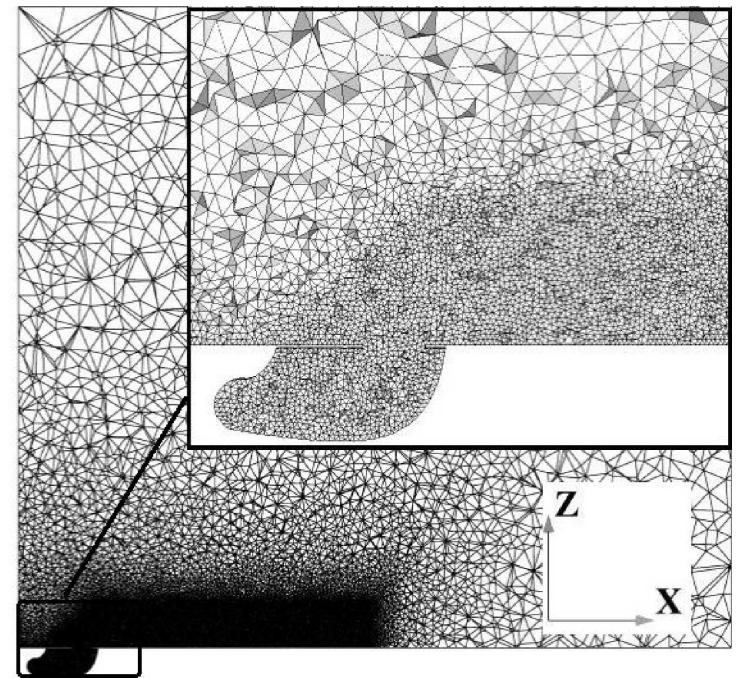
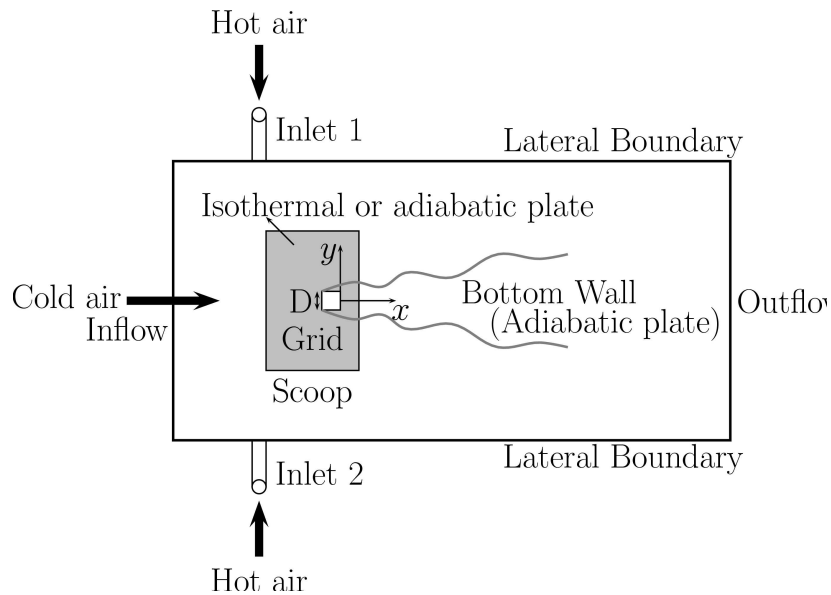
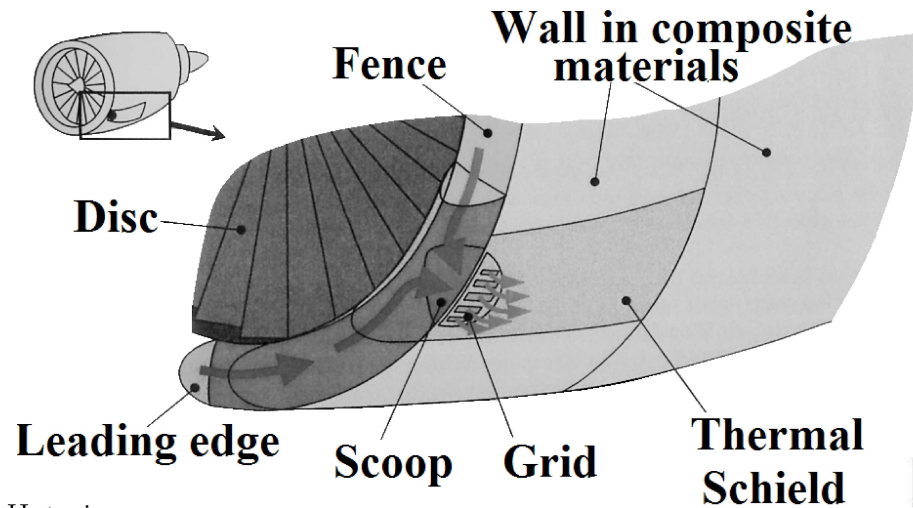
4-4

2-2

4. Response surface analysis & uncertainty quantification tools

Dealing with complex flow LES

(Jouhaud & Sagaut, *J. Fluid Engng*, 2008)



- In complex configurations:
 - optimal values of subgrid models are not known
 - best tuning of artificial viscosity parameter not known
- ⇒ these two parameters are considered as uncertain parameters
- ⇒ comparison with experimental data should account for possible numerical result variability

- Optimal linear unbiased statistical predictor
- Based on sampling points (1 sample = 1 usual simulation)
- Several variants have been developed (cokriging, ...)
- Kriging methods also provide an estimation of the interpolation error
- Sampling points can be generated dynamically to minimize the interpolation error (adaptive refinement)

Estimator at position x

Estimated function at position x_s

$$\hat{f}(x) = c^T(x, x_s) C^{-1}(x, x_s) f(x_s)$$

Covariance vector

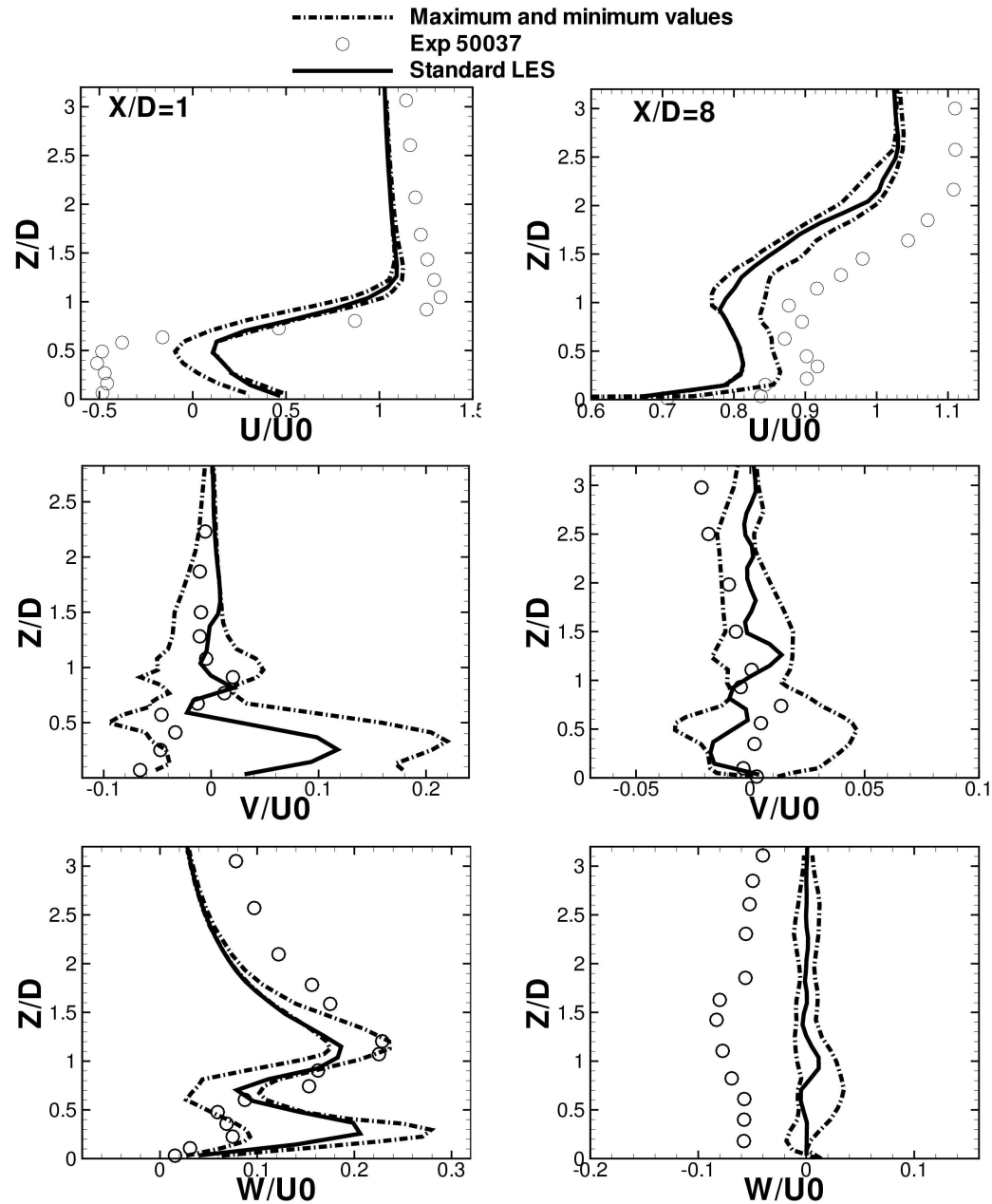
Covariance matrix

$$c(x, x_s) = \begin{pmatrix} Cov(x, x_1) \\ \dots \\ Cov(x, x_n) \end{pmatrix}$$

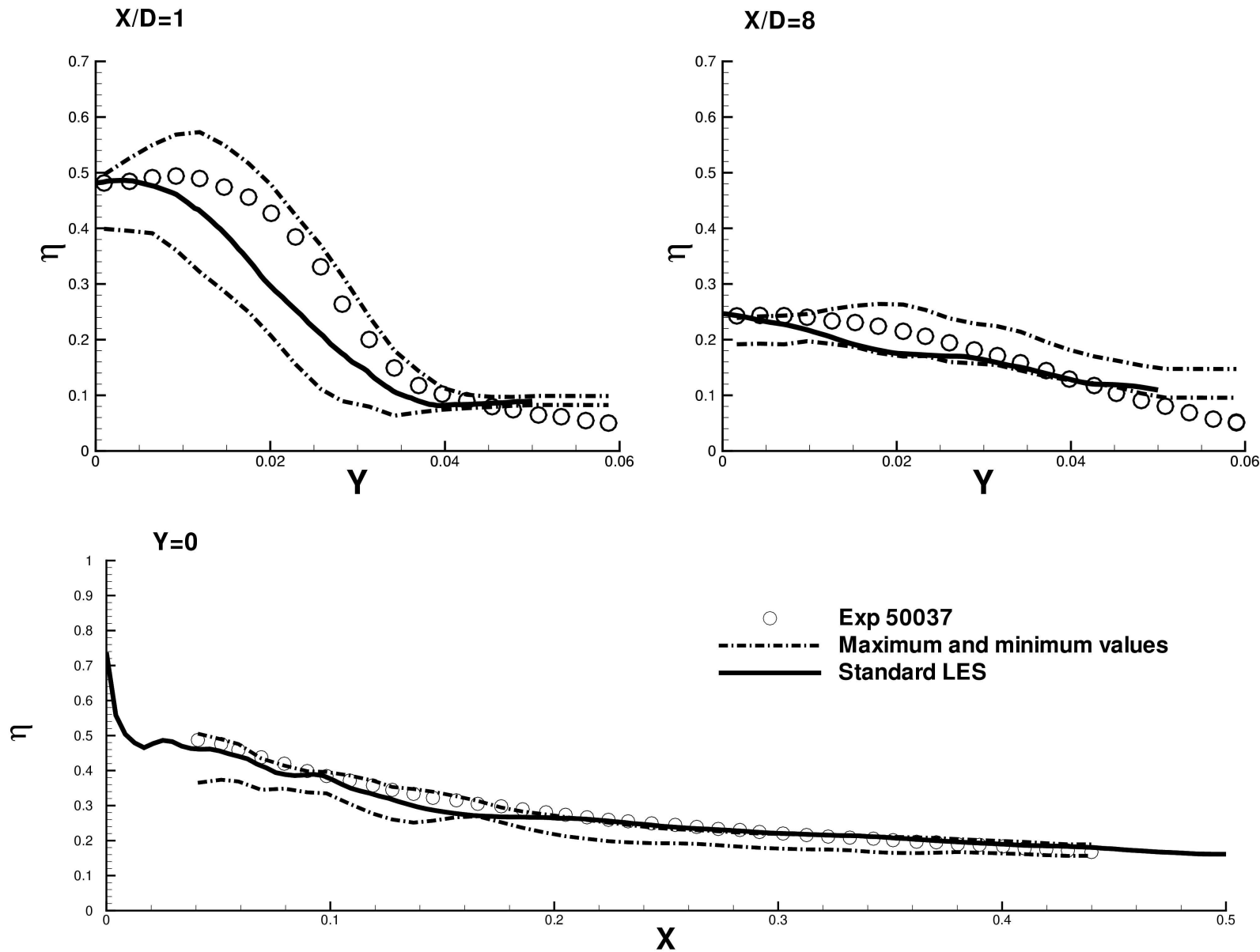
$$C(x_s, x_s) = \begin{pmatrix} \sigma^2 & Cov(x_1, x_2) \dots Cov(x_1, x_n) \\ Cov(x_2, x_1) & \sigma^2 & \dots Cov(x_2, x_n) \\ \dots & \dots & \dots & \dots \\ Cov(x_n, x_1) Cov(x_n, x_2) \dots & \sigma^2 \end{pmatrix}$$

a priori covariogram function: $Cov(y, z) = \sigma^2 \exp(-|y - z|)$

Mean flow predicted by LES



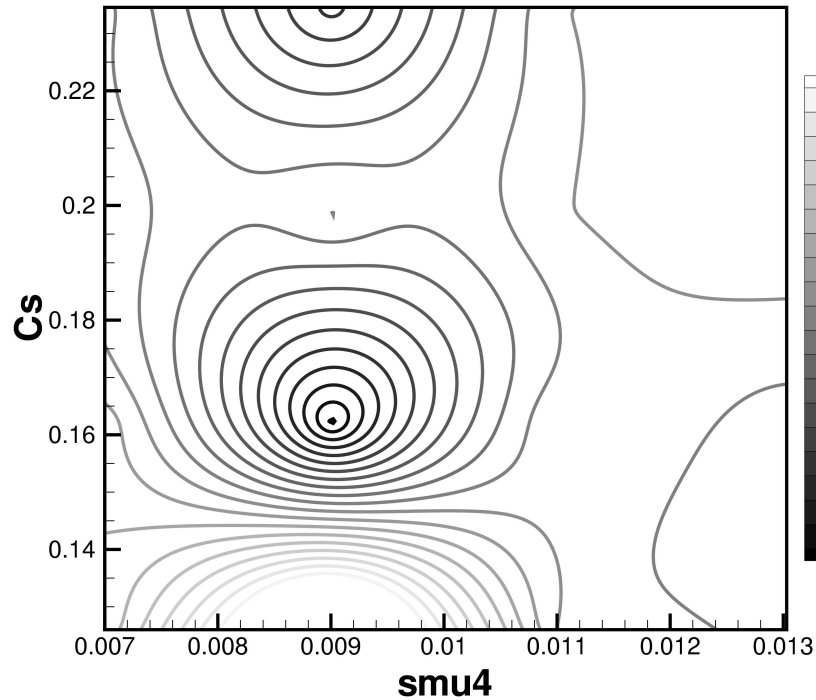
Mean temperature field



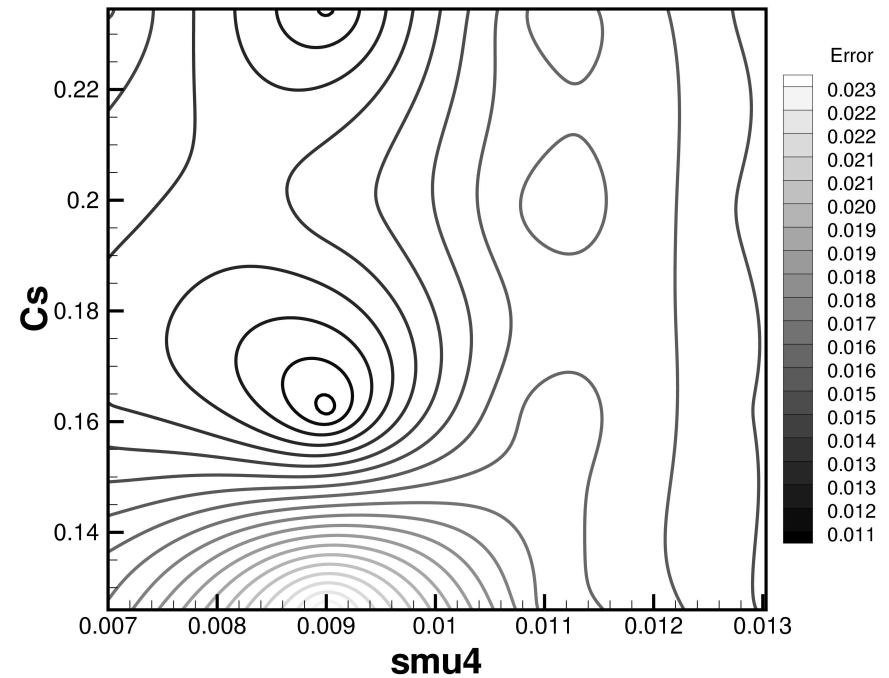
Which solution is the best LES solution ?

- If some experimental data are available:
 - some error functions can be defined
 - solutions with the lowest error norm can be identified
 - ⇒ « clean » definition of the best/optimal LES solution(s)
- « best » LES a priori depends on the error measure

Kriging-based response surface of error at $X/D=8$

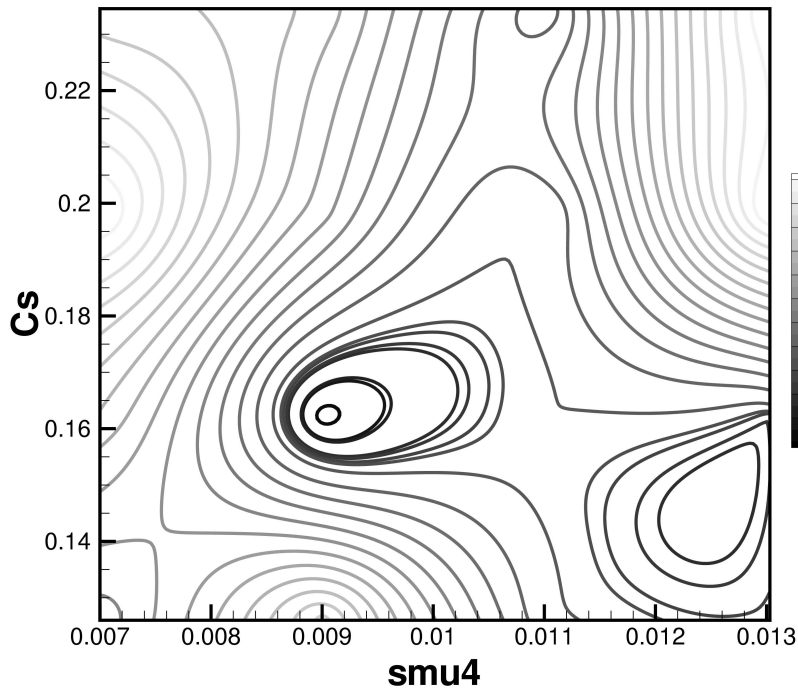


L_1 norm

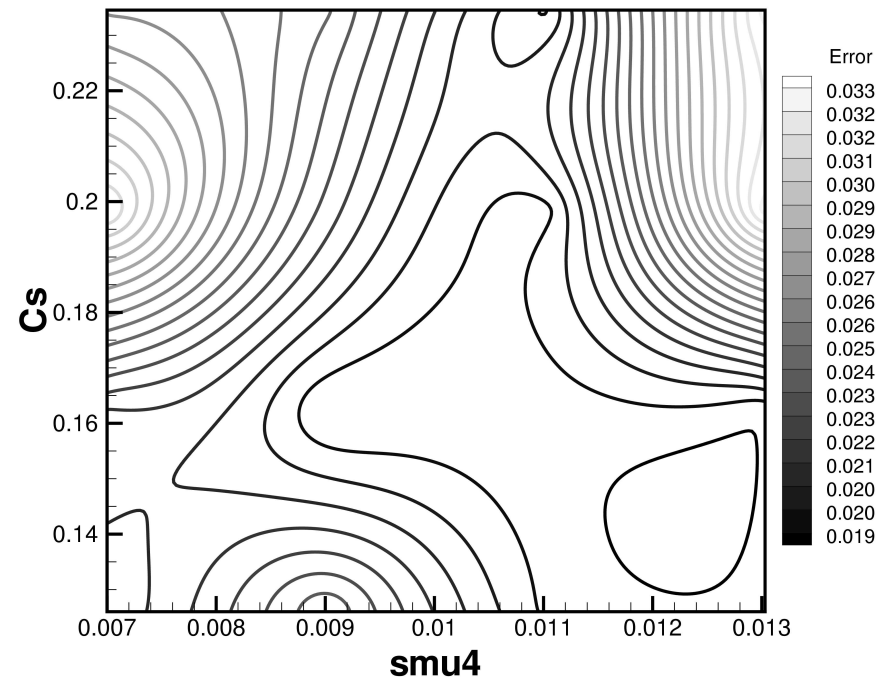


L_2 norm

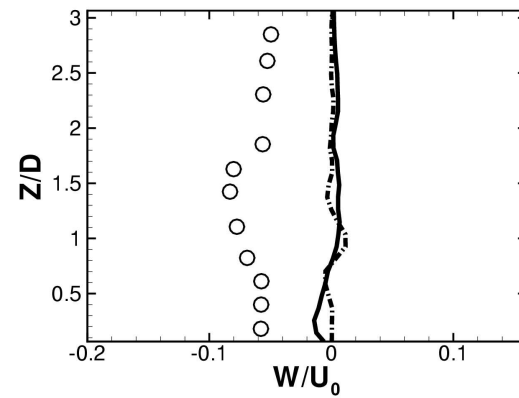
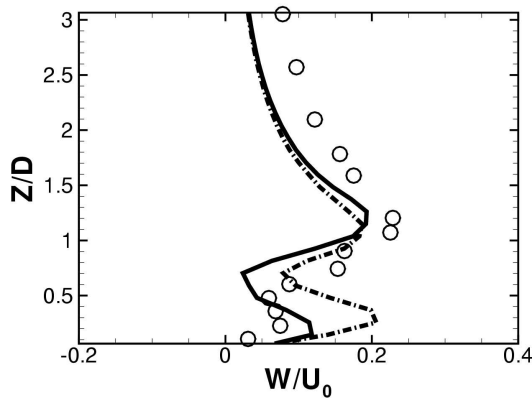
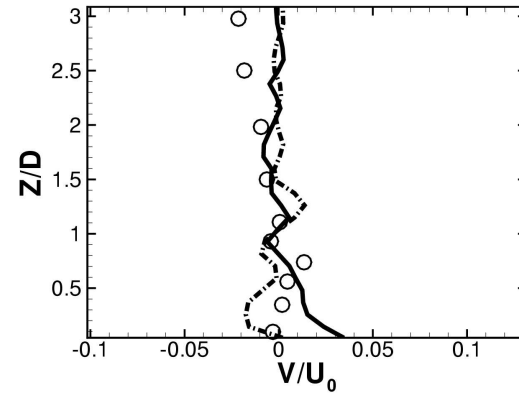
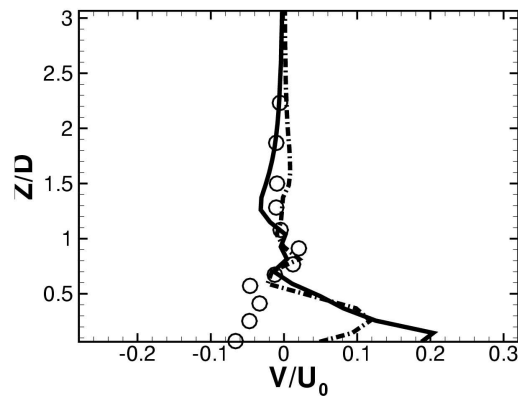
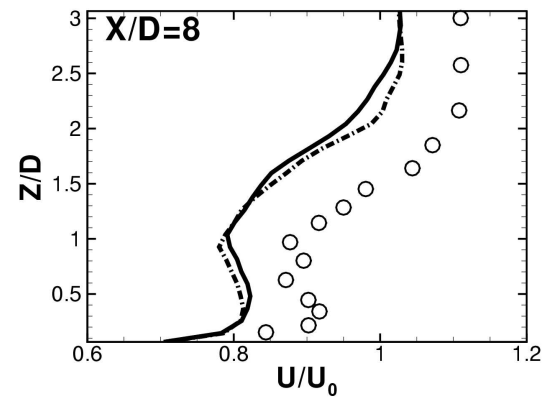
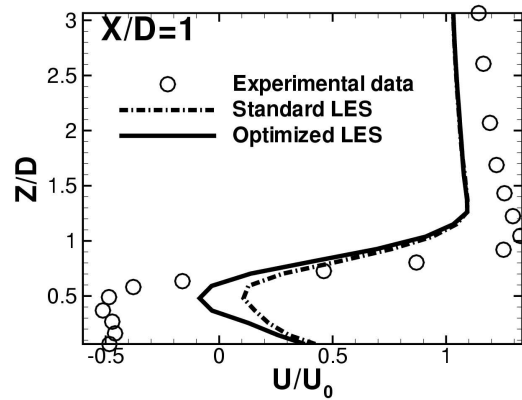
Kriging-based response surface of global error
at both locations $X/D=8$ and $X/D=1$



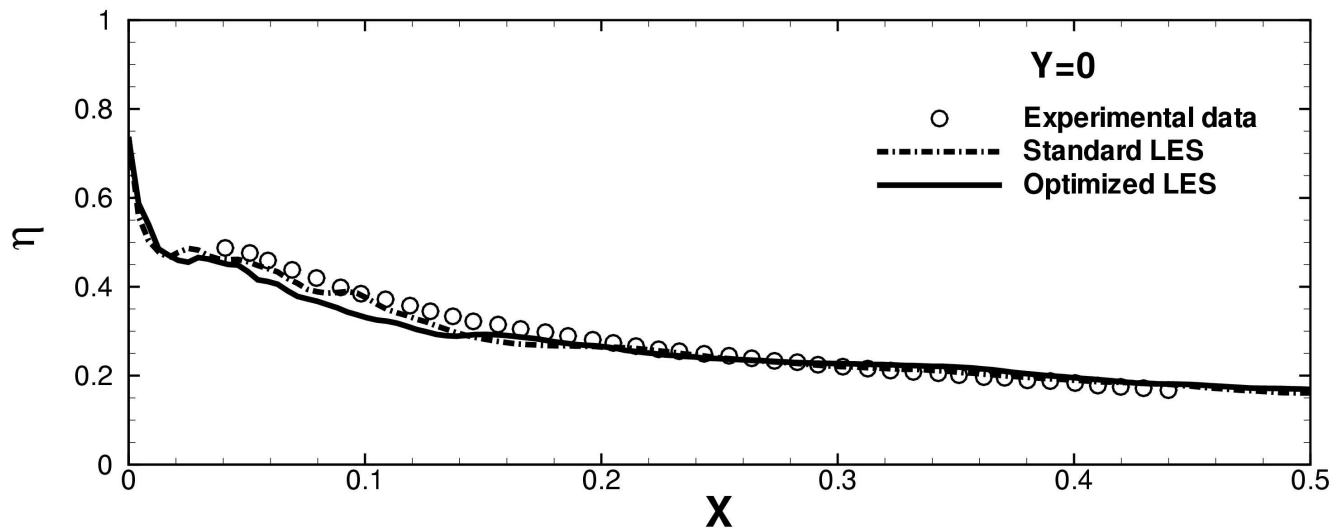
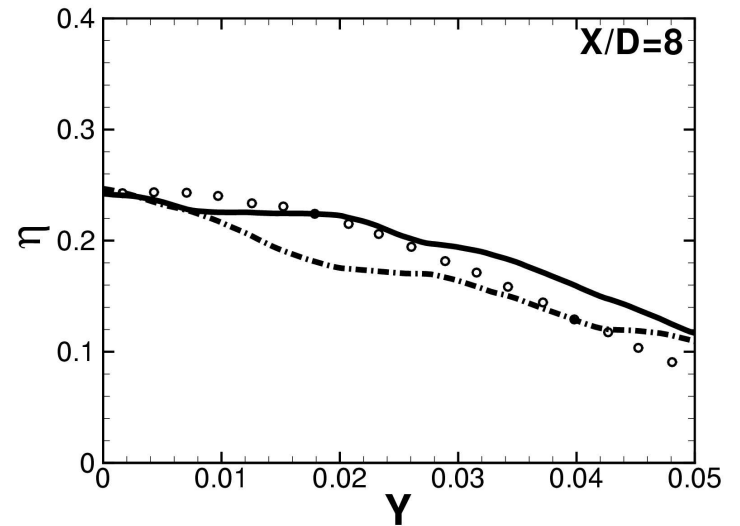
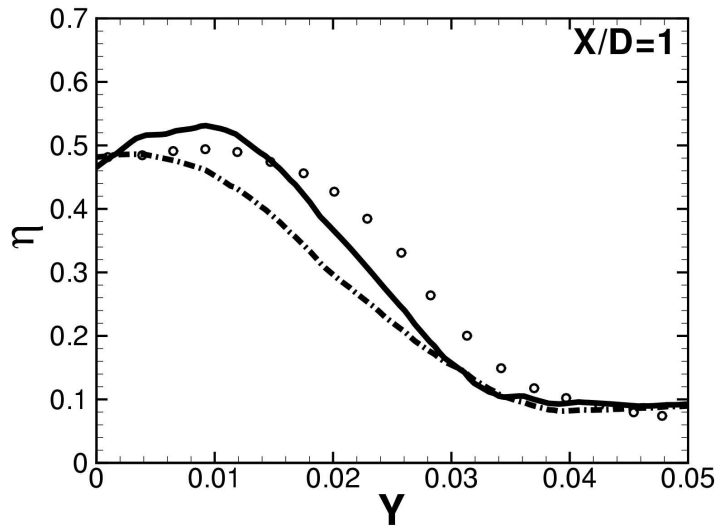
L_1 norm



L_2 norm



Best LES solution



2. Uncertainty in inflow BC

[Ko, J., Lucor, D., Sagaut, P. (2008) *Phys. Fluids* 20, 077102]

- Spatially developing shear flows very sensitive to inflow BC
- Dispersion in CFD results (jets, ...)
- Case study: 2D mixing layer with stochastic inflow BC

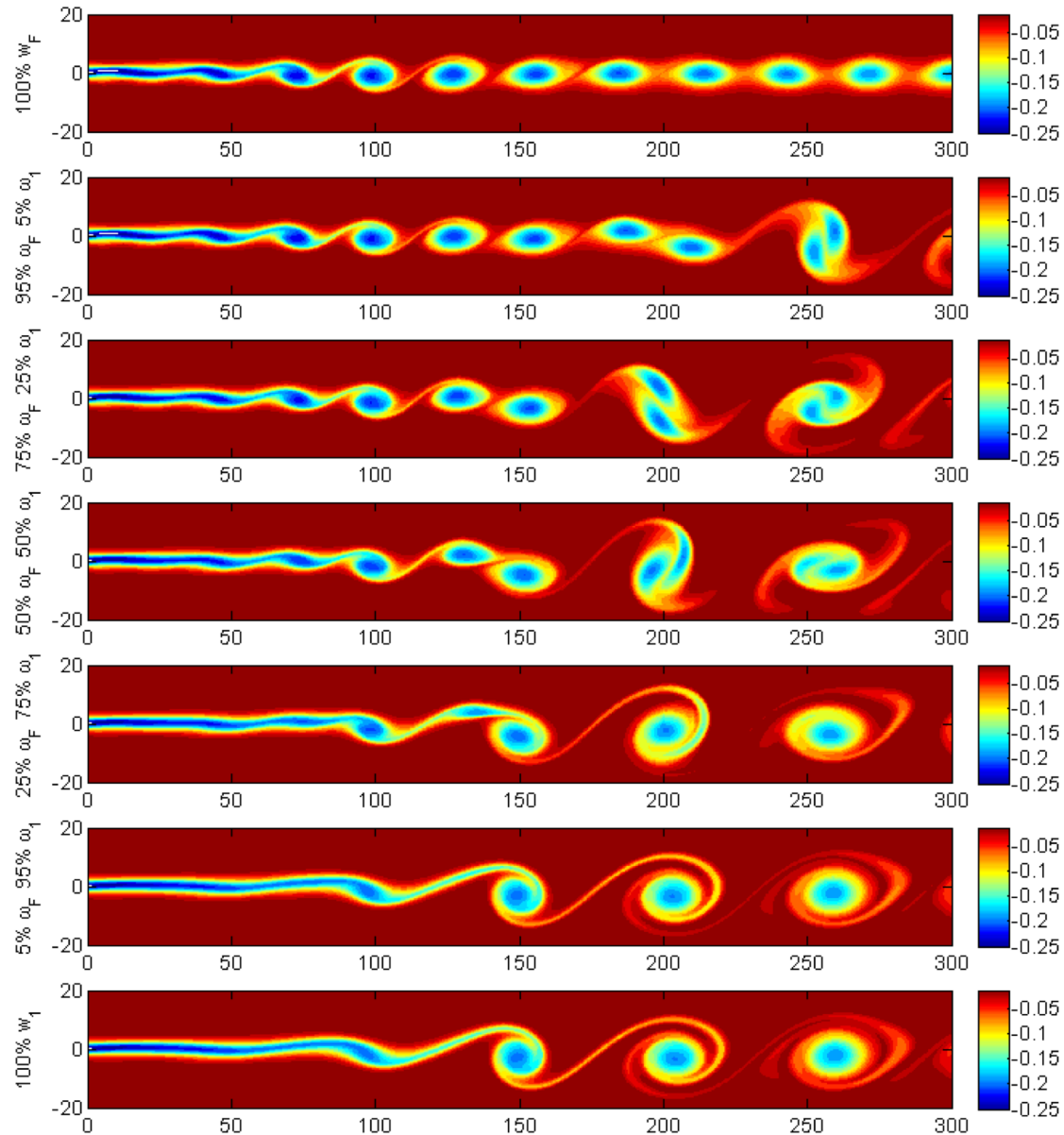
$$u_{\text{in}}(y, t) = \bar{u}_{\text{in}}(y) + \sum_{i=1}^N \epsilon_i [f(y) \sin(\omega_i t) + \gamma_i]$$

$$v_{\text{in}}(y, t) = \bar{v}_{\text{in}}(y) + \sum_{i=1}^N \frac{\epsilon_i}{i} [g(y) \sin(\omega_i t) + \gamma_i]$$

Present results: $N = 2$ $\epsilon_1 + \epsilon_2 = U/10$

2. Deterministic views

Fundamental
mode forcing

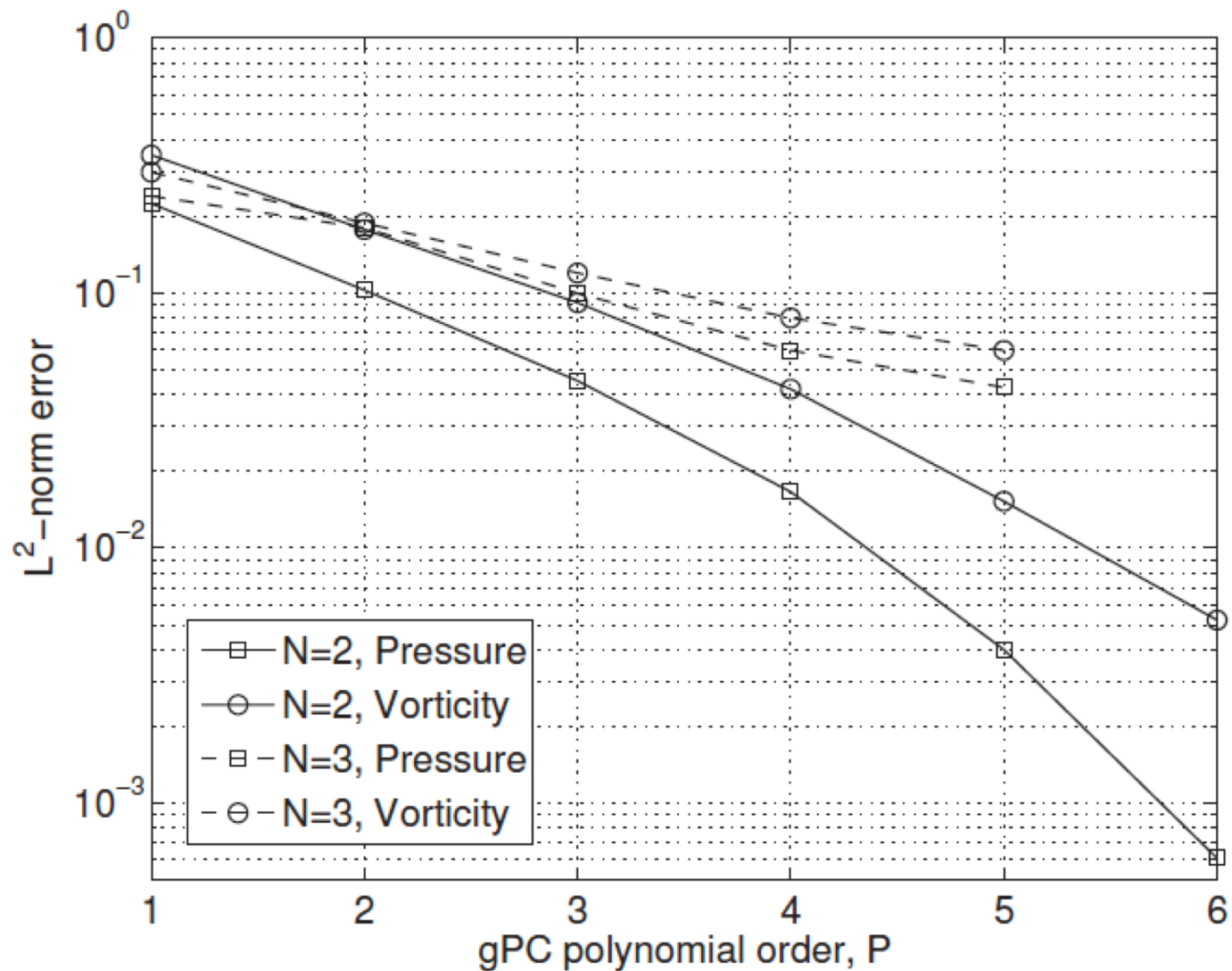


$$\varepsilon_2 = 0$$

$$\varepsilon_1 = 0$$

1st subharmonic
mode forcing

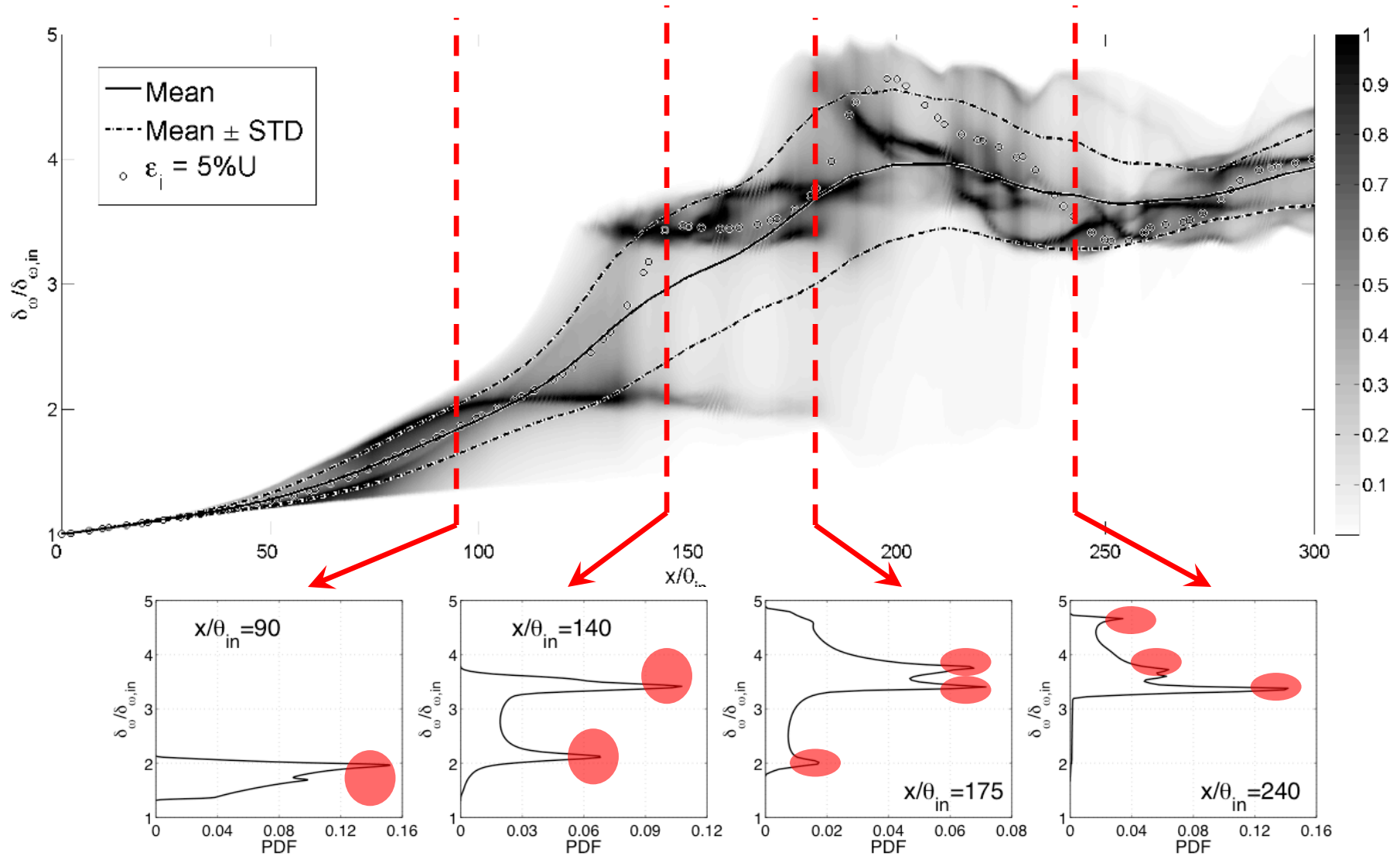
NI-gPC expansion convergence

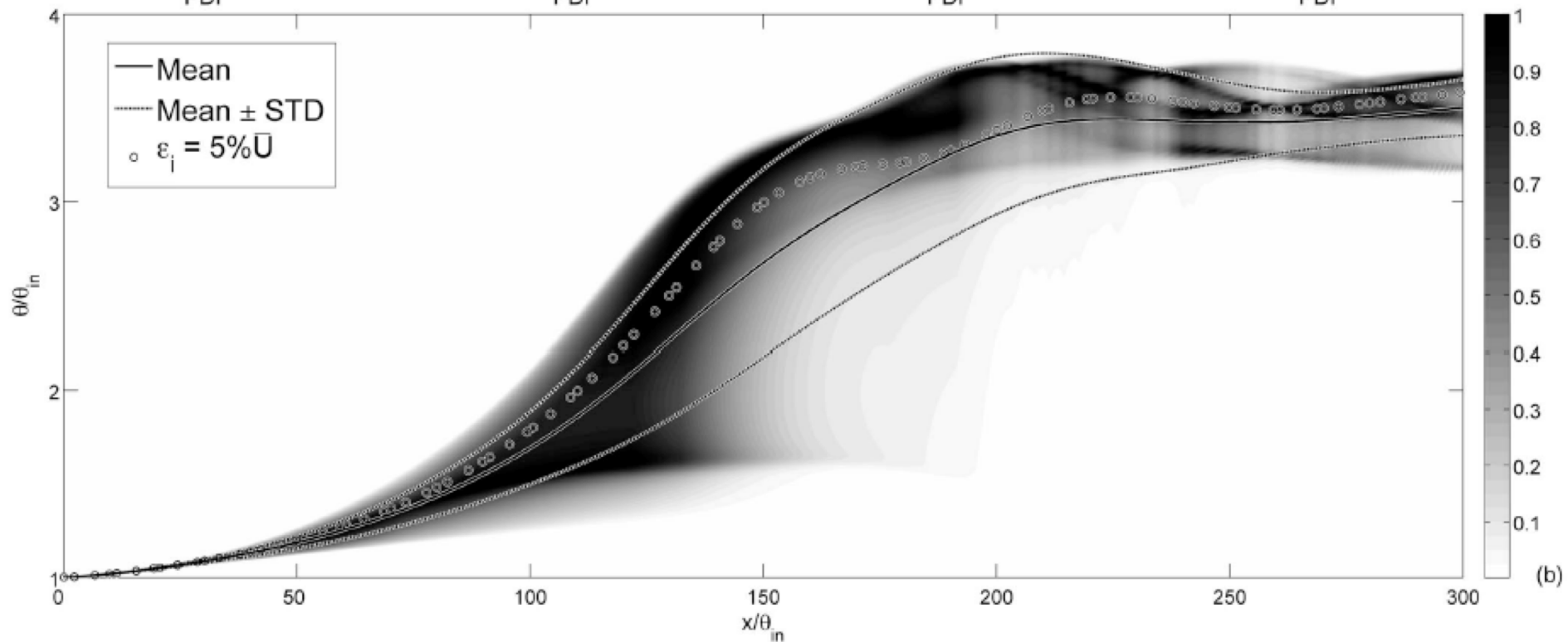
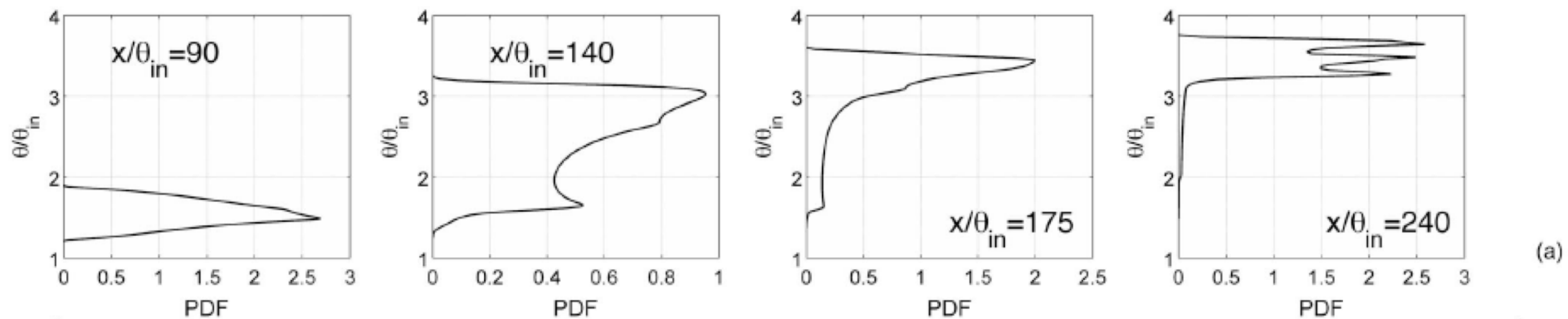


→ *Pseudo-spectral convergence obtained*

2. Stochastic analysis

PDF of vorticity thickness of time-averaged fl





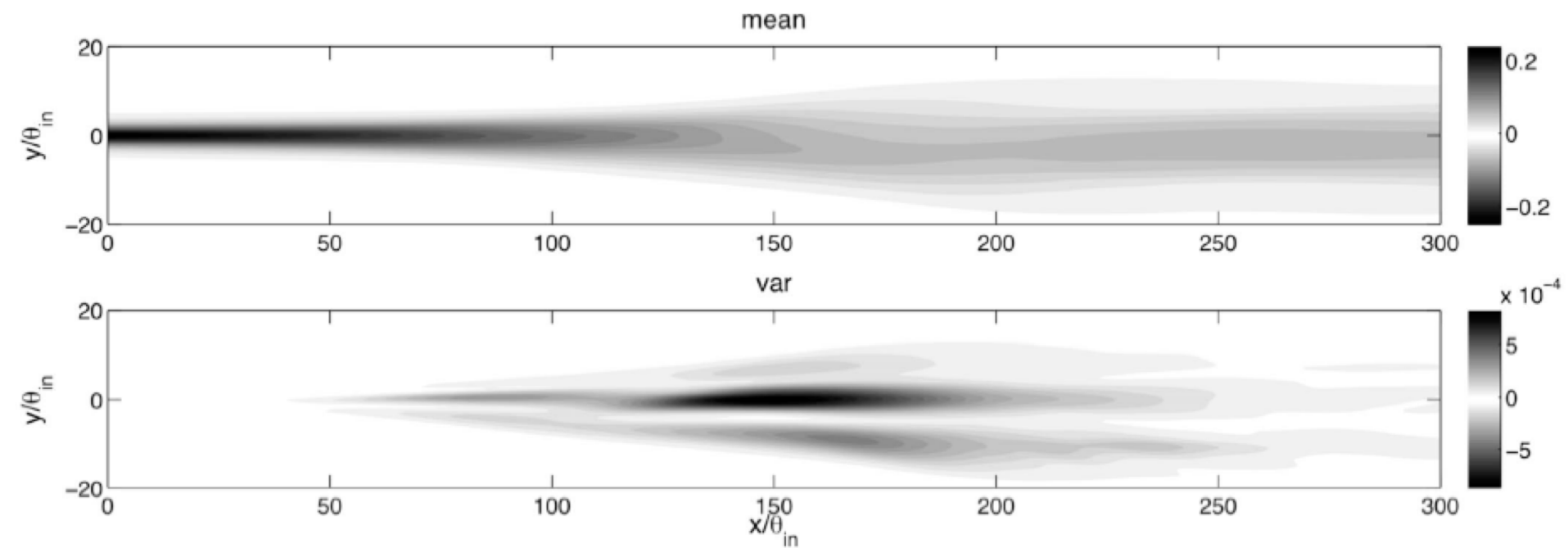


FIG. 6. Spatial distributions of the mean solution and variance of the time-averaged vorticity in the mixing layer with bimodal perturbation forcing.

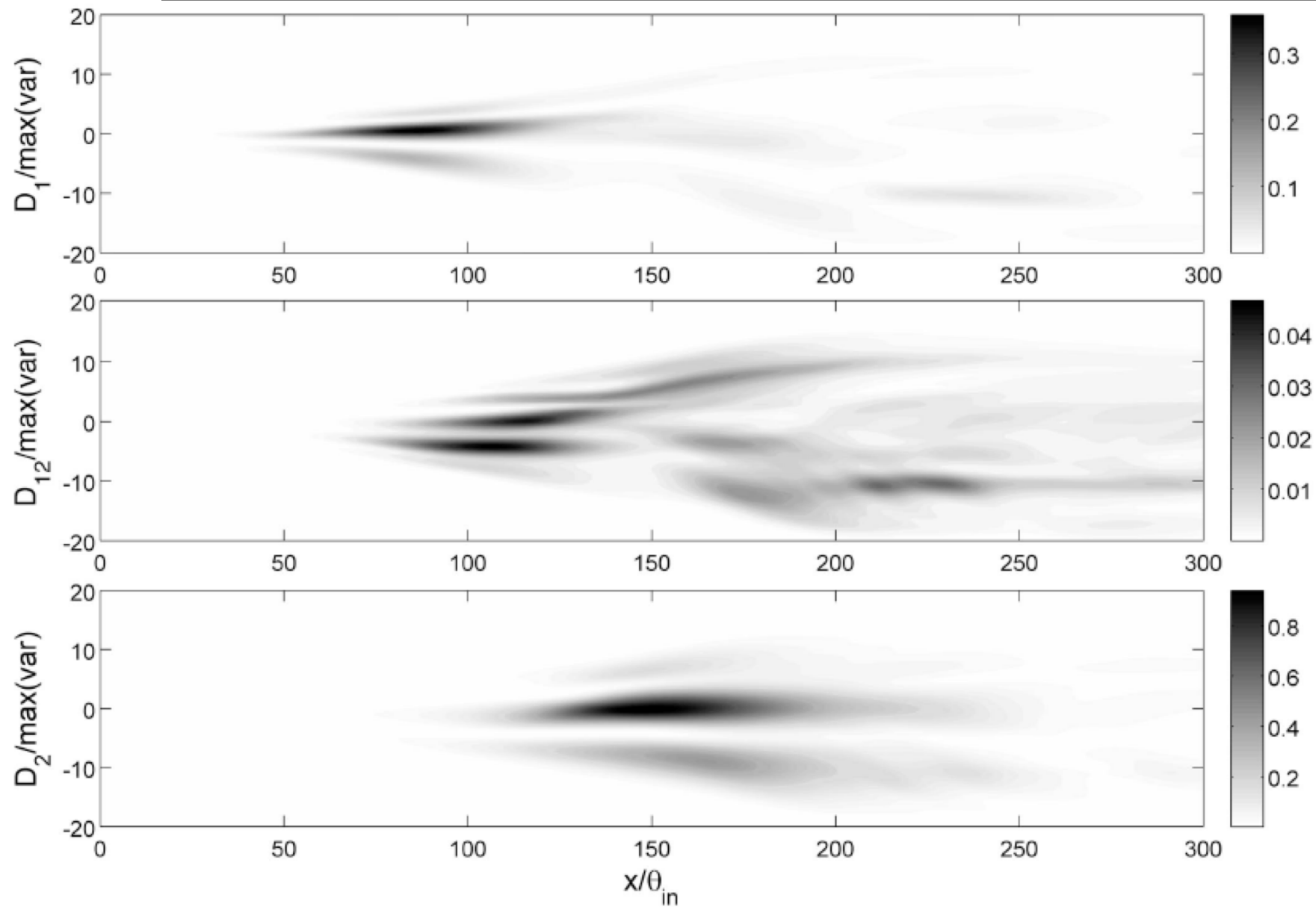
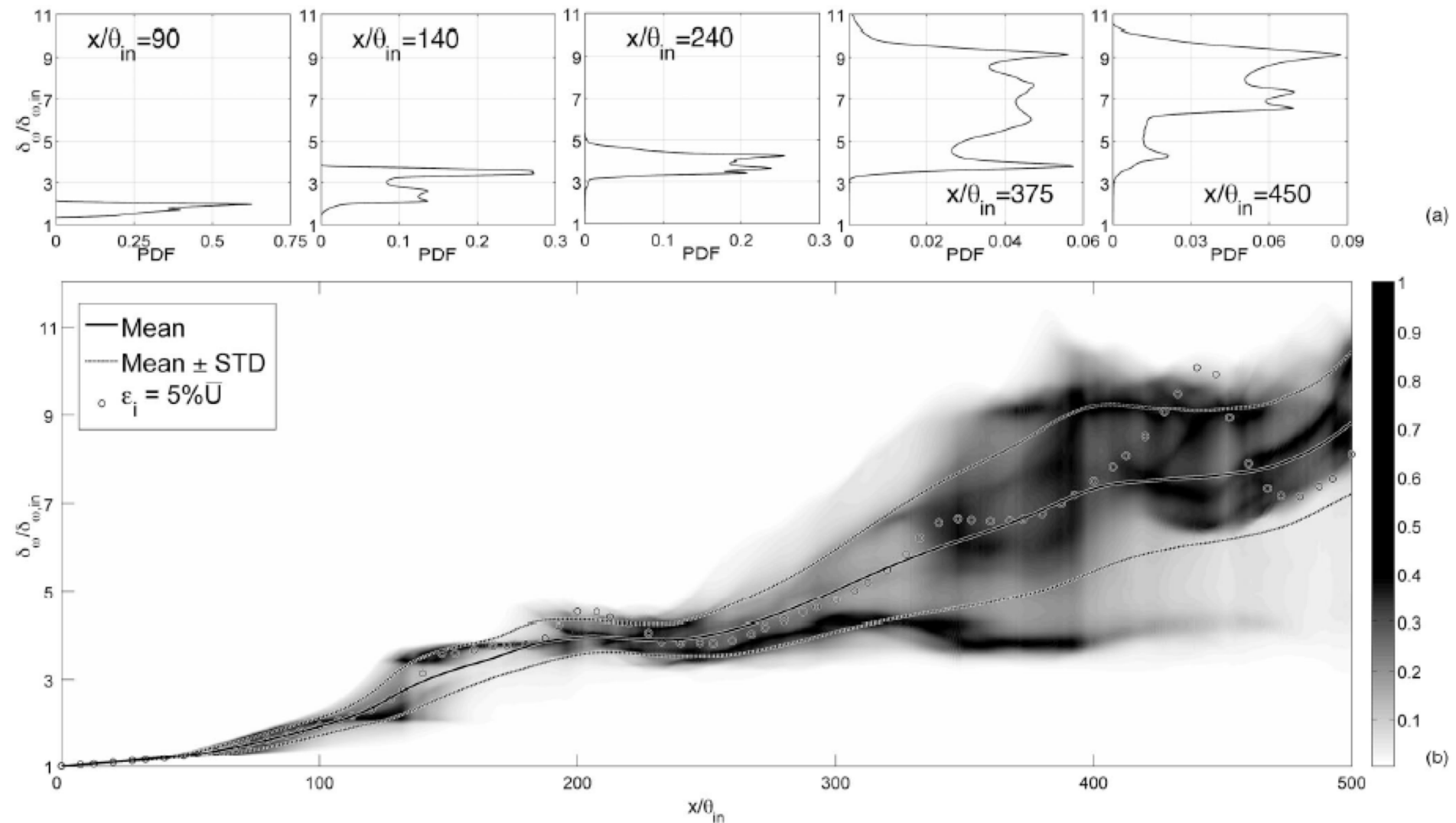
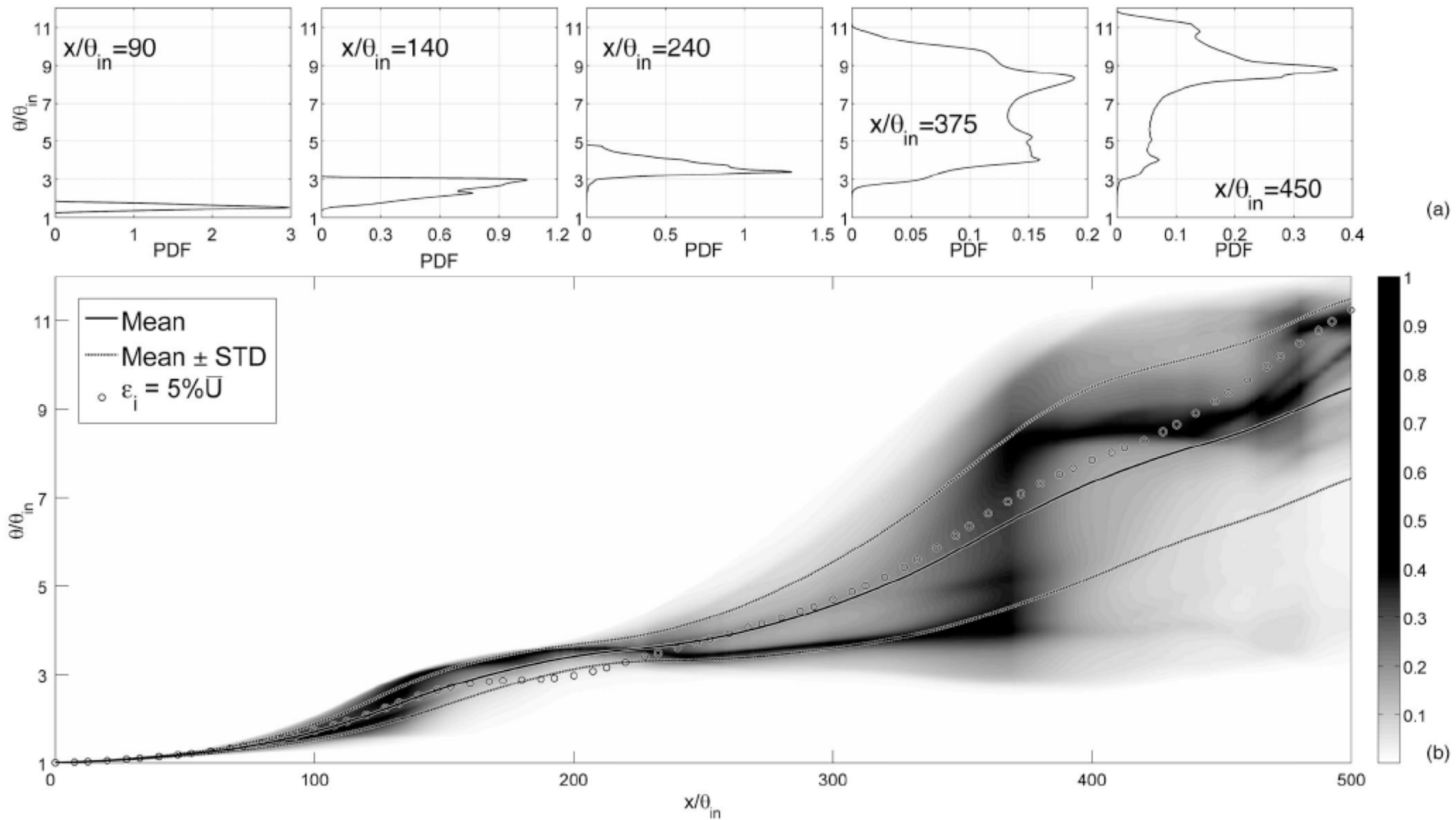


FIG. 10. Spatial distributions of the sensitivity indices (partial variances) of the time-averaged vorticity in the mixing layer with bimodal perturbation forcing.



Trimodal case



Trimodal case

UNCERTAINTY IN THE COEFFICIENTS

k- ϵ

Launder-Sharma low-Re



$$C_{\mu} = 0.09$$

$$C_{\epsilon 1} = 1.44$$

$$C_{\epsilon 2} = 1.92$$

$$\sigma_{\epsilon} = 1.30 \quad \sigma_{\kappa} = 1$$

RANS

MODELS

MODEL

PARAMETERS

k- ω

Menter - SST



$$\beta_1 = 0.0750$$

$$\beta_2 = 0.0828$$

$$\gamma_1 = 0.55 \quad \gamma_2 = 0.44$$

$$\sigma_{\kappa 1} = 0.85 \quad \sigma_{\kappa 2} = 1$$

$$\sigma_{\omega 1} = 0.5$$

$$\sigma_{\omega 2} = 0.856$$

$$\beta^* = 0.09$$

DATA FROM LITERATURE

USED IN INDUSTRIAL

APPLICATIONS

Considering THREE CANONICAL TURBULENT FLOWS

- freely decaying homogeneous isotropic turbulence
- homogeneous shear flow in the asymptotic regime
- boundary layer in the logarithmic region.

ANALYTICAL EXPRESSIONS

$$\begin{aligned} C_\mu &= \left(\frac{\kappa_{log}}{u_\tau^2} \right)^{-2} \\ C_{\varepsilon_1} &= \frac{C_{\varepsilon_2} - 1}{\frac{\mathcal{P}}{\varepsilon}} \\ C_{\varepsilon_2} &= 1 - \frac{1}{n} \\ \sigma_\varepsilon &= \frac{\kappa_{VK}^2}{\sqrt{C_\mu (C_{\varepsilon_2} - C_{\varepsilon_1})}} \\ \sigma_K &= 1 \end{aligned}$$

$$\begin{aligned} C_\mu &= f \left(\frac{\kappa_{log}}{u_\tau^2} \right) \\ C_{\varepsilon_1} &= f \left(n, \frac{\mathcal{P}}{\varepsilon} \right) \\ C_{\varepsilon_2} &= f(n) \\ C_\mu &= f \left(n, \frac{\mathcal{P}}{\varepsilon}, \kappa_{VK}, \frac{\kappa_{log}}{u_\tau^2} \right) \\ - & \end{aligned}$$

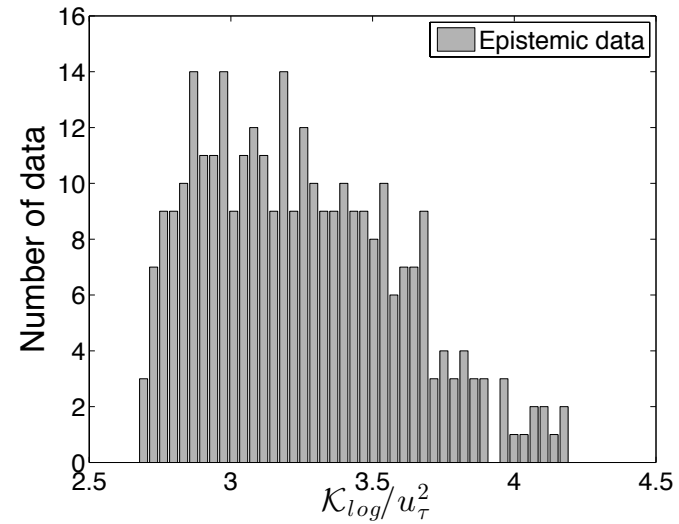
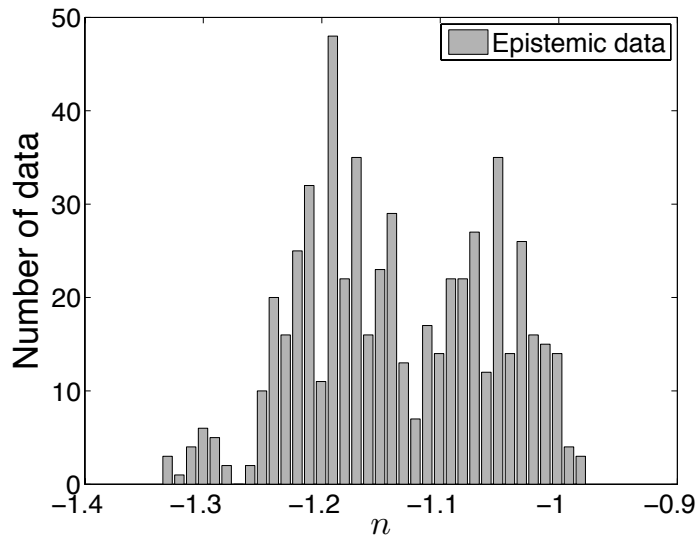
$$\begin{aligned} \beta^* &= \left(\frac{\kappa}{u_\tau^2} \right)^{-2} \\ \beta_1 &= -\frac{\beta^*}{n} \\ \gamma_1 &= \frac{\beta_1}{\beta^* \frac{\mathcal{P}}{\varepsilon}} \\ \sigma_{\omega_1} &= \frac{\sqrt{\beta^* (\beta_1 \beta^* - \gamma_1)}}{\kappa_{VK}^2} \\ \sigma_{K_1} &= \frac{\beta_1 \sigma_{\omega_1}}{\beta^*} \end{aligned}$$

$$\begin{aligned} \beta^* &= f \left(\frac{\kappa}{u_\tau^2} \right) \\ \beta_1 &= f \left(n, \frac{\kappa_{log}}{u_\tau^2} \right) \\ \gamma_1 &= f \left(n, \frac{\mathcal{P}}{\varepsilon}, \frac{\kappa_{log}}{u_\tau^2} \right) \\ \sigma_{\omega_1} &= f \left(n, \frac{\mathcal{P}}{\varepsilon}, \kappa_{VK}, \frac{\kappa_{log}}{u_\tau^2} \right) \\ \sigma_{K_1} &= f \left(n, \frac{\mathcal{P}}{\varepsilon}, \kappa_{VK}, \frac{\kappa_{log}}{u_\tau^2} \right) \end{aligned}$$

MODEL COEFFICIENTS → FLOW PROPERTIES

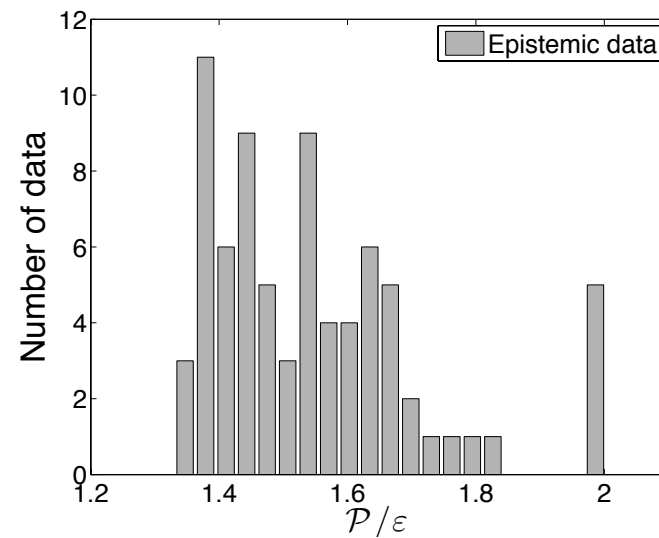
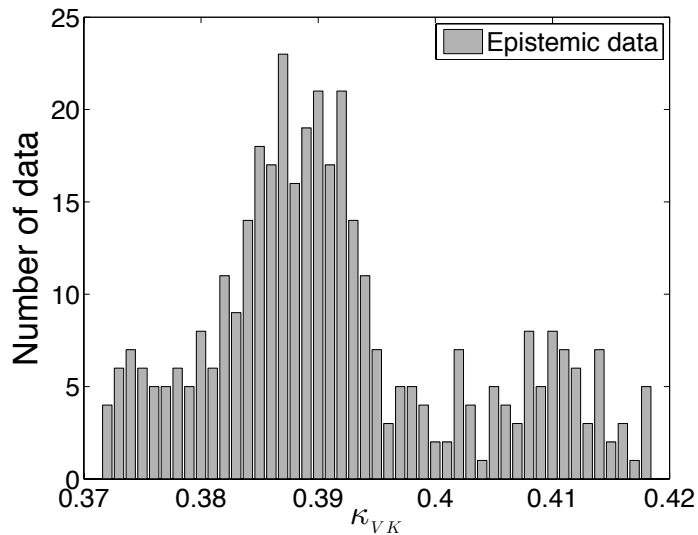
LITERATURE SURVEY

n



$\frac{K_{\log}}{u_{\tau}^2}$

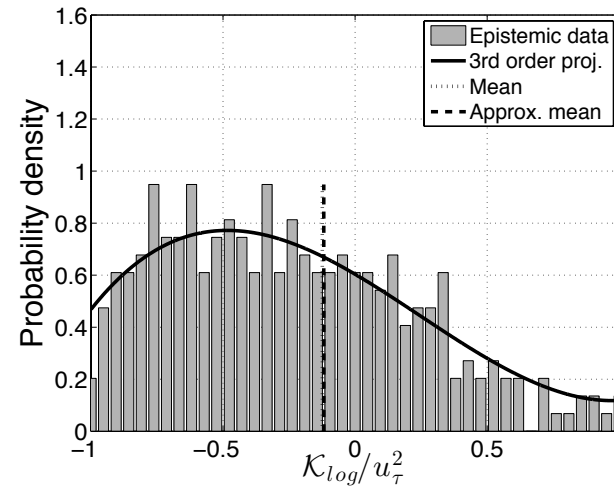
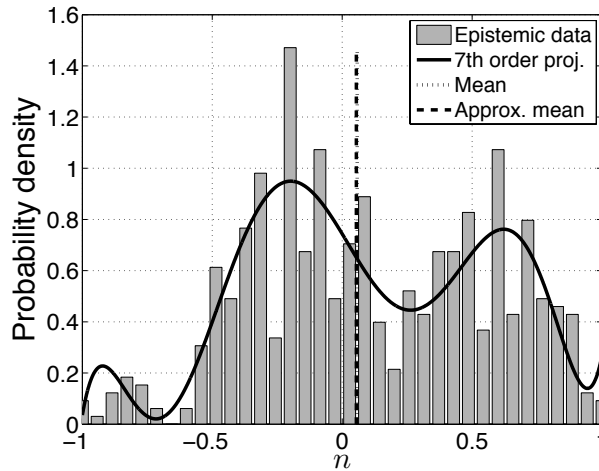
κ_{VK}



$\frac{P}{\epsilon}$

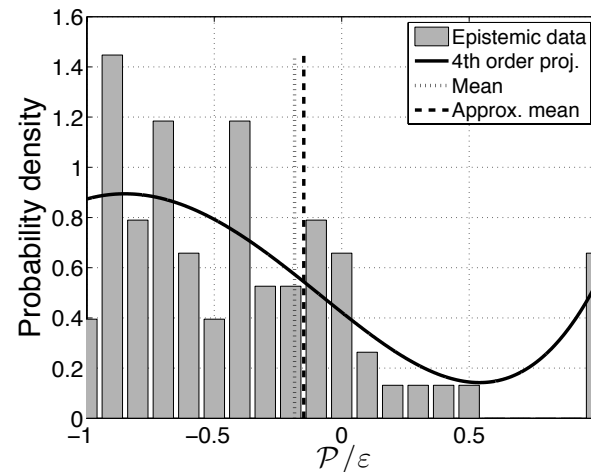
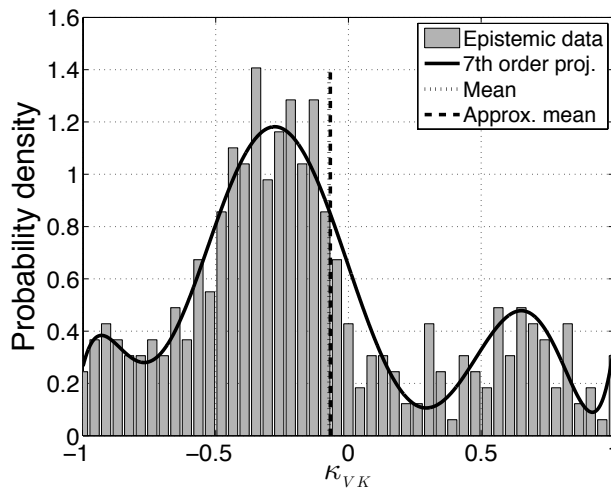
CONTINUOUS PDF APPROXIMATION FOR THE EPISTEMIC INPUT PARAMETERS

n



$\frac{K_{log}}{u_{\tau}^2}$

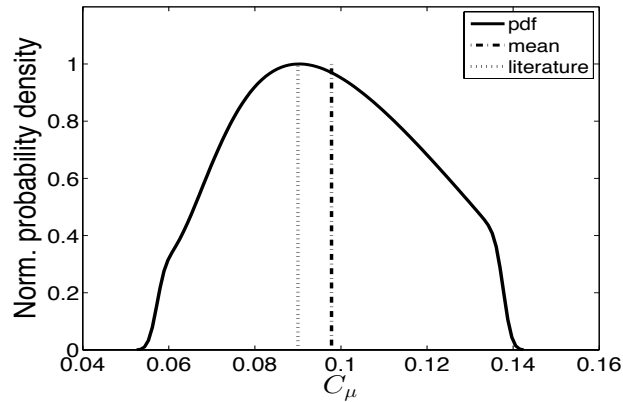
K_{vk}



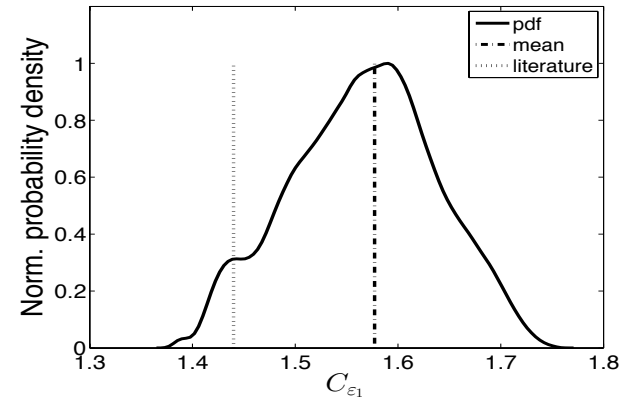
$\frac{P}{\epsilon}$

UNCERTAINTY PROPAGATION IN THE **k-ε** LS MODEL COEFFICIENTS

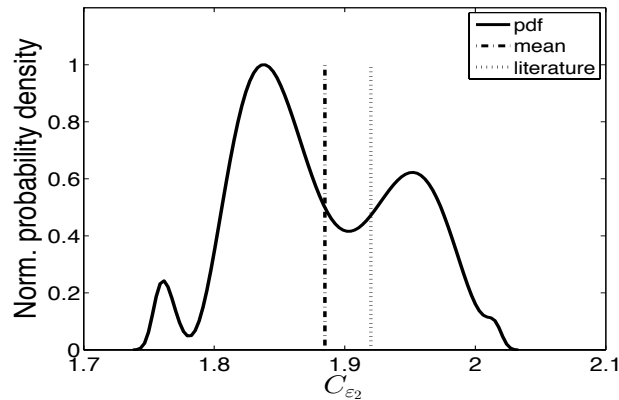
C_μ



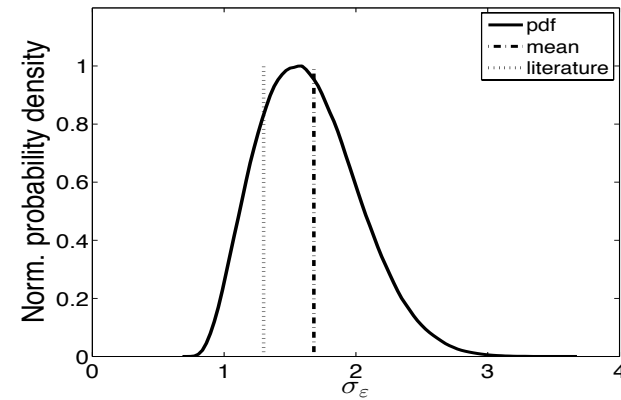
$C_{\epsilon 1}$



$C_{\epsilon 2}$



σ_ϵ

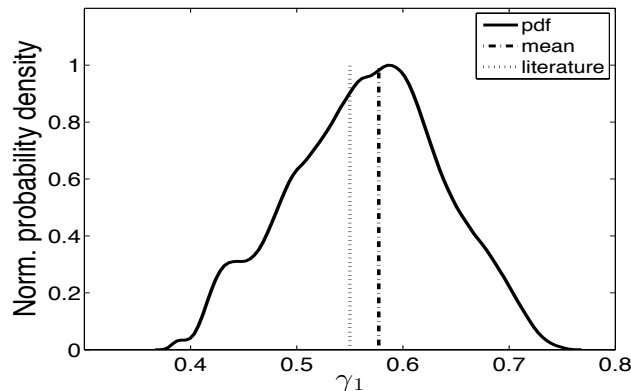


Mean value for
 $C_\mu \approx 0.09$

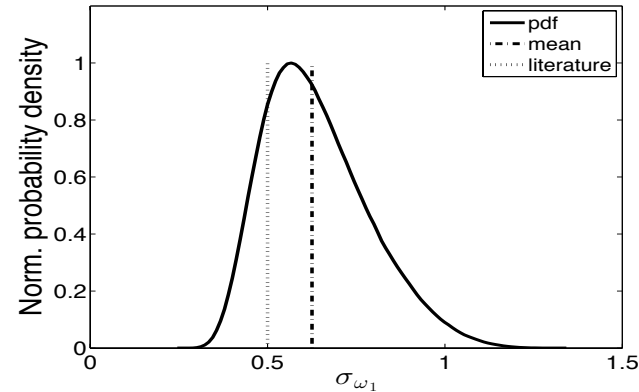
High sensitivity to
 P/ϵ

UNCERTAINTY PROPAGATION IN THE **k- ω SST** MODEL COEFFICIENTS

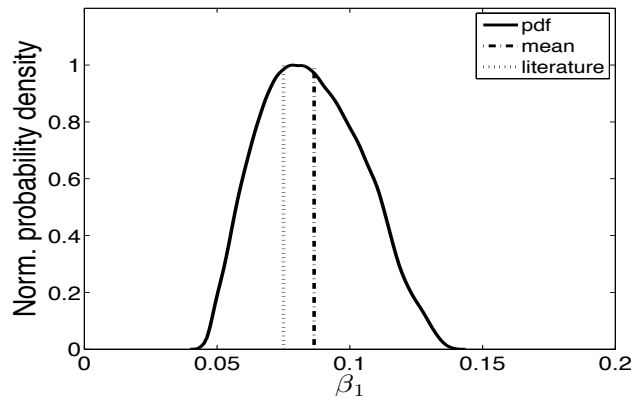
γ_1



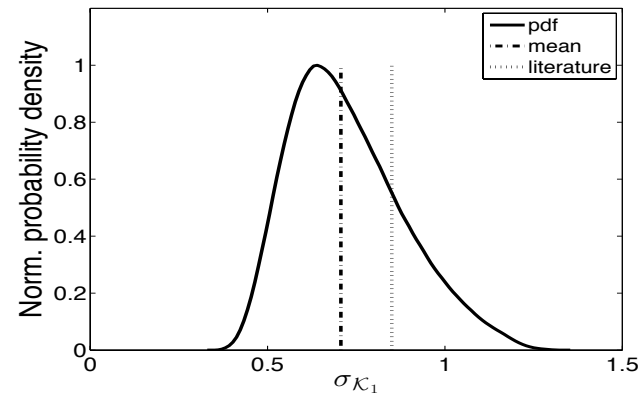
$\sigma_{\omega 1}$



β_1



σ_{k1}



**Mean values $\approx$
literature**

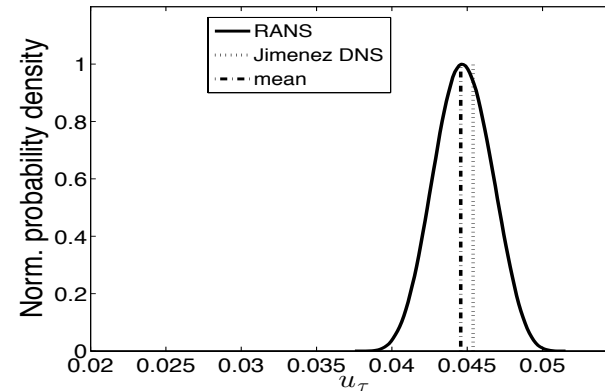
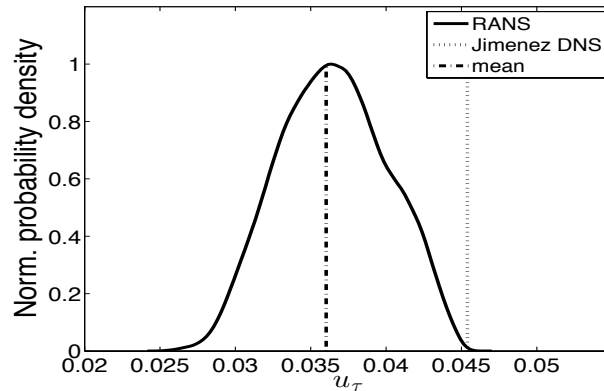
**High sensitivity to
 P/ε**

SOME RESULTS: FRICTION VELOCITY

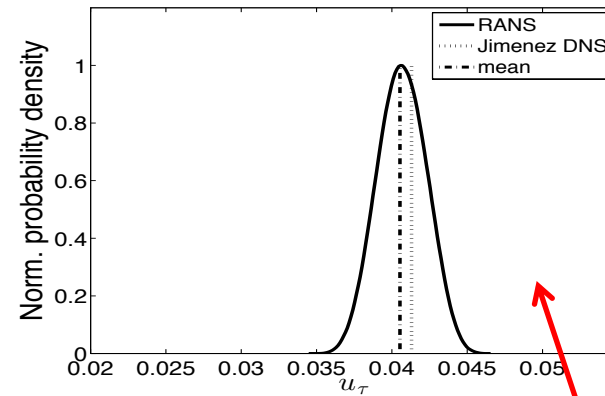
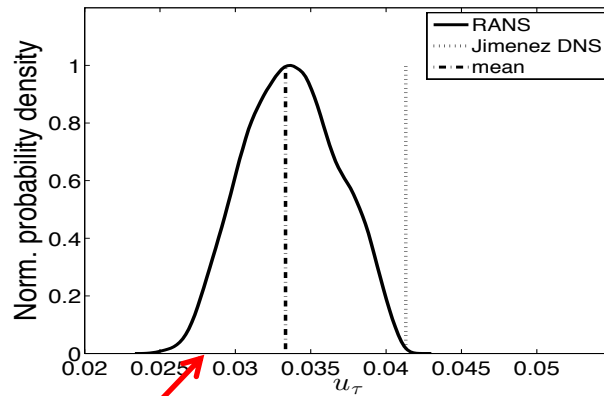
k- ϵ LS

k- ω SST

Re $_{\tau}$ =950



Re $_{\tau}$ =2000



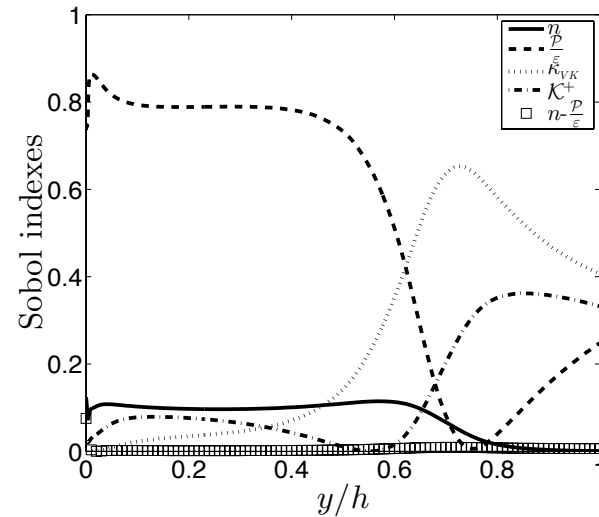
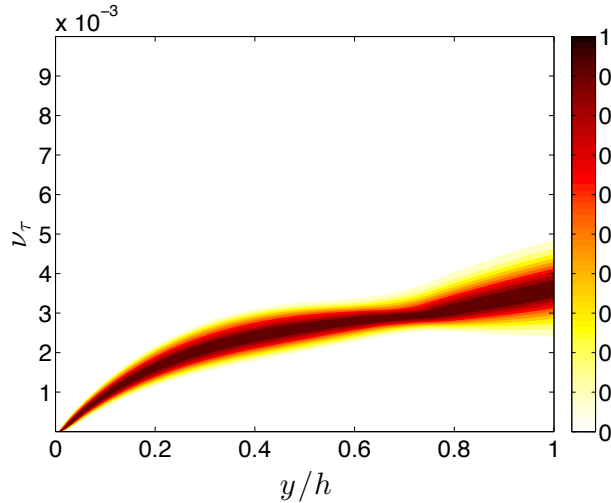
**Underprediction and
high variance**

**High sensitivity to
P/ ϵ**

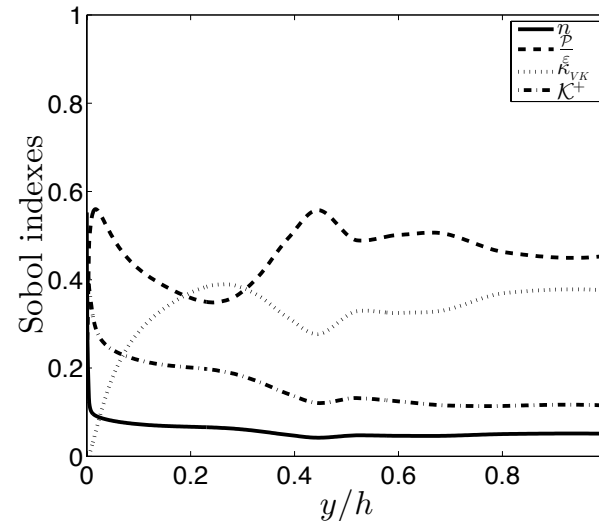
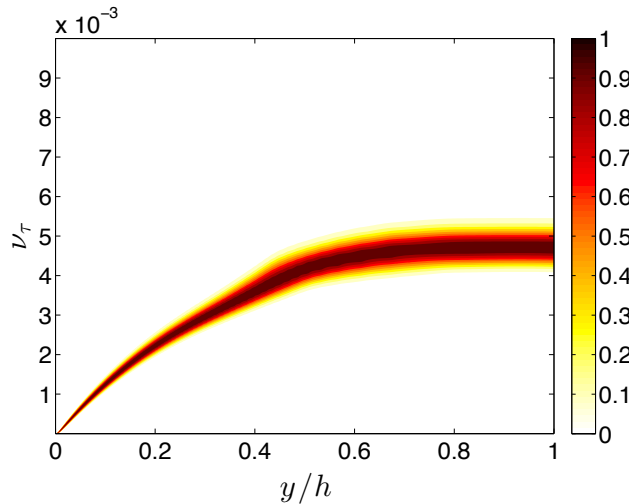
**Less
sensitive**

SOME RESULTS: EDDY VISCOSITY PROFILE

k- ϵ
LS
Re $_{\tau}$ =2000



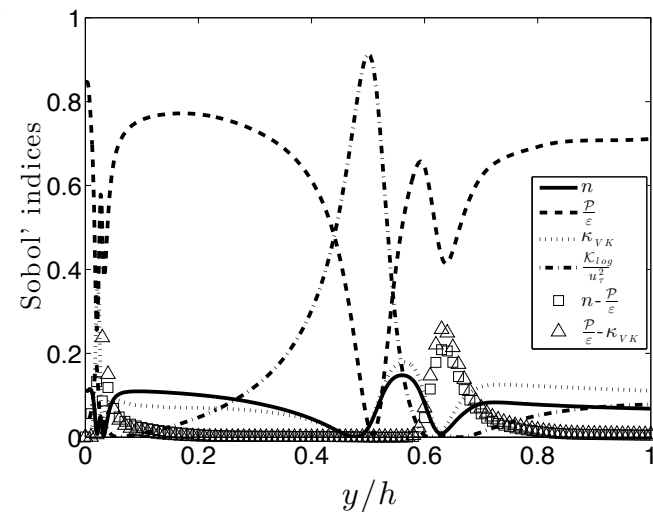
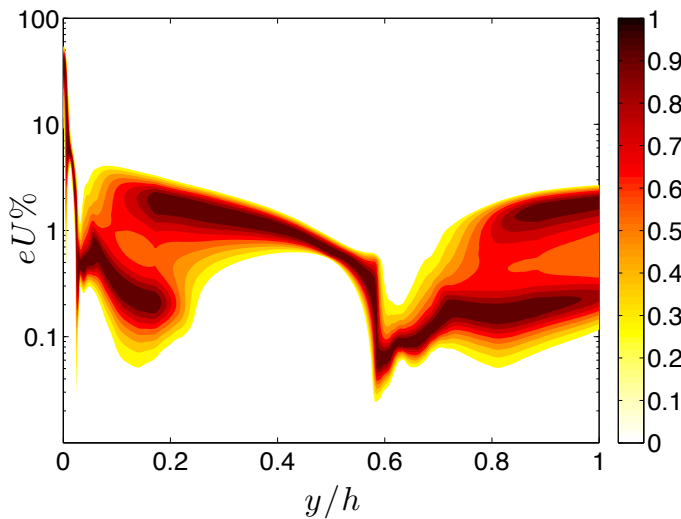
k- ω
SST
Re $_{\tau}$ =2000



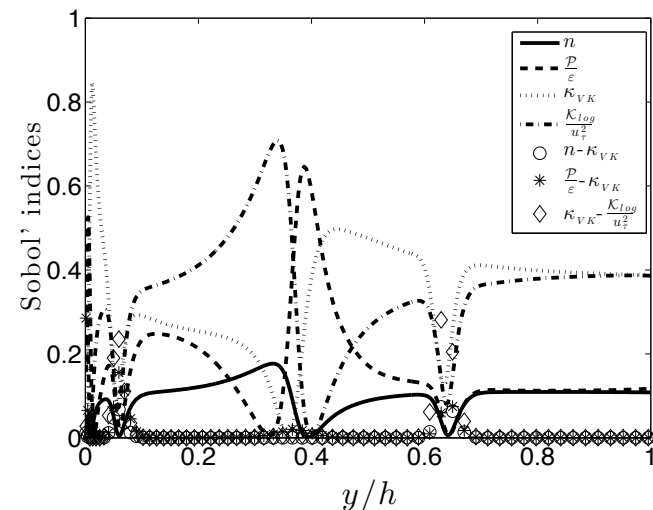
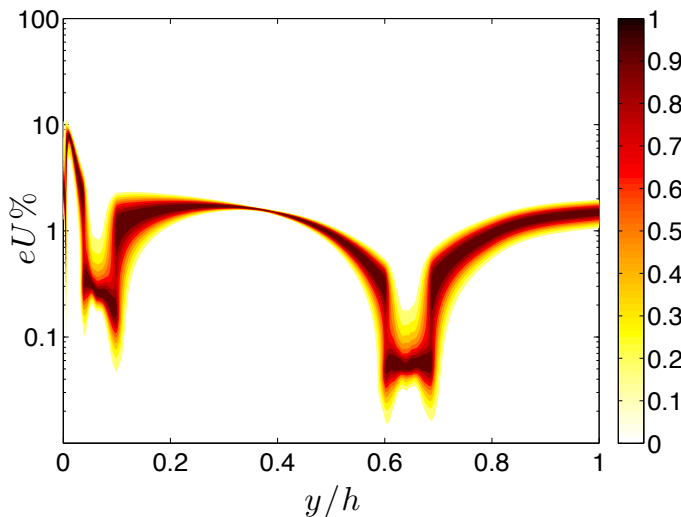
SOME RESULTS: LOCAL L_2 ERROR ON THE VELOCITY

$$eU(y) = \sqrt{\frac{(U(y) - U_{DNS}(y))^2}{U_{DNS}^2(y)}}$$

k- ϵ
LS
Re $_{\tau}$ =2000



k- ω
SST
Re $_{\tau}$ =2000



Thank you !