

## Collapse of a (non-)axisymmetric air cavity in water <br> Devaraj van der Meer • Oscar Enríquez • Ivo Peters • Stefan Gekle • Laura Schmidt • Utkarsh Jain • Anaïs Gauthier • Detlef Lohse



## UNIVERSITY OF TWENTE.

PHYSICS OF FLUIDS


October 17, 2017 - Paris

## Disclaimer

- singularity
- (controlled) instabilities
- air entrapment
- experiments, theory \& numerics

Talk by Utkarsh Jain - Wednesday at 16:30

## Try this at home


... in our lab



## Experimental setup



## Disk pulled through interface



$$
\begin{aligned}
& V_{\text {impact }}=1.0 \mathrm{~m} / \mathrm{s} \\
& R_{\text {disk }}=0.03 \mathrm{~m} \\
& \text { camera @ } 1000 \mathrm{fps}
\end{aligned}
$$

## Series of events



void creation

## Series of events



void creation

void collapse

## Series of events



void creation

void collapse
jet creation at singularity

## Series of events


void creation
void collapse
jet creation at singularity
air-entrainment "giant bubble"

## Dimensional Analysis

## Relevant parameters:

- disk radius $R_{0} \approx 2 \mathrm{~cm}$
- density $\rho=1000 \mathrm{~kg} / \mathrm{m}^{3}$
- disk velocity $V \approx 1 \mathrm{~m} / \mathrm{s}$
- viscosity $\eta=1.0 \mathrm{mPas}$
- gravity $g=9.8 \mathrm{~m} / \mathrm{s}^{2}$
- surface tension $\sigma=0.074 \mathrm{~N} / \mathrm{m}$

$$
\begin{aligned}
\mathrm{Fr} & =\frac{V^{2}}{g R_{0}} \approx 5 \\
\mathrm{We} & =\frac{\rho V^{2} R_{0}}{\sigma} \approx 300 \\
\mathrm{Re} & =\frac{\rho V R_{0}}{\eta} \approx 20,000
\end{aligned}
$$

(Froude number) inertia and gravity dominant
(Weber number)
(Revnalds number) viscosity unimportant $\Rightarrow$ potential flow
incompressible:
$\vec{\nabla} \cdot \vec{u}=0 \Rightarrow$

$$
\nabla^{2} \varphi=0
$$

## Boundary Integral simulations

Laplace equation for potential:

$$
\nabla^{2} \varphi=0
$$

is solved as a boundary integral:
$\varphi\left(\vec{x}_{0}, t\right)=\iint_{\partial V}\left[G\left(\left|\vec{x}-\overrightarrow{x_{0}}\right|\right) \vec{\nabla} \varphi(\vec{x}, t)-\varphi(\vec{x}, t) \vec{\nabla} G\left(\left|\vec{x}-\overrightarrow{x_{0}}\right|\right)\right] \cdot \overrightarrow{d A}$
(Green's third identity)
Unsteady Bernoulli equation provides time evolution:

$$
\frac{\partial \varphi}{\partial t}+\frac{1}{2}|\vec{\nabla} \varphi|^{2}=-g z-\frac{\sigma}{\rho} \kappa
$$

## BI simulation vs. experiments

$\mathrm{Fr}=3.4$

$\mathrm{Fr}=13.6$

No free parameter!


## Model: Slender cavity limit



Flow in horizontal layers:
$\rightarrow$ Assume potential flow
$\rightarrow$ Assume axisymmetry
$\rightarrow$ Neglect axial flow
needed: equation for 2D fluid flow in layers

## 2D Rayleigh-Besant equation

Euler equation in cylindrical coordinates

$$
\frac{\partial u_{r}}{\partial t}+u_{r} \frac{\partial u_{r}}{\partial r}=-\frac{1}{\rho} \frac{\partial p}{\partial r}
$$

Continuity equation

$$
r u_{r}=R \dot{R}
$$

Integrate with boundary conditions:

result $\frac{d}{d t}(R \dot{R}) \log \frac{R}{R_{\infty}}+\frac{1}{2} \dot{R}^{2}=g z$
2D Rayleigh equation
R. Bergmann et al, PRL 96, 154505 (2006).

## Void creation (at microscale)



Impact of a train of micro droplets

$$
\begin{aligned}
& (d=100 \mu \mathrm{~m} ; \\
& \mathrm{V}=12.6 \mathrm{~m} / \mathrm{s})
\end{aligned}
$$

$$
\begin{gathered}
\frac{d}{d t}(R \dot{R}) \log \frac{R}{R_{\infty}}+\frac{1}{2} \dot{R}^{2}=\underset{\text { large }}{\text { gravity }} \begin{array}{c}
\text { negligible }
\end{array} \\
\frac{d}{d t}(R \dot{R})=\frac{d^{2}}{d t^{2}}\left(\frac{1}{2} R^{2}\right)=0 \\
\Rightarrow R(t)=\sqrt{2 R_{0} \dot{R}_{0}\left(t+t_{0}\right)}
\end{gathered}
$$

Void created with constant velocity $V$ :

$$
t=\frac{z}{V}
$$

$$
R(z)=\sqrt{2 R_{0} \frac{\dot{R}_{0}}{V}\left(z+z_{0}\right)}
$$

Parabolic shape!

## Cavity shape (at microscale)


W. Bouwhuis et al, preprint (2014).

## Cavity shape (disc impact)

$\Delta t_{\text {reach }}=\frac{z}{V}$ determined by impact speed
$\Delta t_{\text {coll }}$
determined by hydrostatic pressure $\frac{d}{d t}(R \dot{R}) \log \frac{R}{R_{\infty}}$


## Neck radius $R(\tau)$ [experiment]



## Neck radius scaling



## Neck radius scaling



What is the expected scaling of the neck radius close to the singularity?

$$
\begin{array}{llll}
-2 & -4 & -3 & -2 \\
& & & -1 \\
\log _{10} \tau
\end{array}
$$

## Back to 2D Rayleigh

$$
\left.\begin{array}{l}
\frac{d}{d t}(R \dot{R}) \underset{\text { large }}{\log \frac{R}{R_{\infty}}}+\frac{1}{2} \dot{R}^{2}=\underset{\substack{\text { gravity } \\
\text { negligible }}}{ } \\
\frac{d}{d t}(R \dot{R})=\frac{d^{2}}{d t^{2}}\left(\frac{1}{2} R^{2}\right)=0 \\
\left.R\left(t_{c}\right)=0 \quad \text { (collapse time } t_{c}\right)
\end{array}\right\} \Rightarrow
$$

$$
R(t) \sim R_{0}\left(t_{c}-t\right)^{1 / 2}
$$

## Rayleigh scaling in experiment?



## Rayleigh scaling in experiment?

Scaling exponent $\alpha$ consistently higher than 2D Rayleigh:

$$
R \sim R_{0} \tau^{1 / 2}
$$

| 0 |  |
| :--- | :--- |
| $00-1$ | $\mathrm{Fr}=163$ |

More elaborate analysis of 2D Rayleigh:

$$
\alpha=\frac{\log \left(\frac{1}{4} \gamma^{2}\right)}{1+2 \log \left(\frac{1}{4} \gamma^{2}\right)}
$$

where $\gamma$ is the cavity aspect ratio

$$
\log _{10} \tau
$$

J. M. Gordillo \& M. Pérez-Saborid, JFM 562, 303 (2006); J. Eggers et al., PRL 98, 094502 (2007);
S. Gekle et al, PRE 80, 036305 (2009).

## How about the airflow in the cavity ?



crown splash \& cavity formation

## How about the airflow in the cavity ?



## Following smoke particles



## Air speed from smoke measurements



Numerical Modeling 1: Boundary-integral

- Potential flow for both liquid and air: irrotational,inviscid, incompressible



## Numerical Modeling 1: Boundary-integral

- Potential flow for both liquid and air: irrotational,inviscid, incompressible



## Numerical modeling 2: Multiscale



## Numerical modeling 2: Multiscale



## Numerical modeling 2: Multiscale



## Results: Air velocity profile



## Results: Air velocity profile



## Results: Air velocity profile



## Results: Air flow at the cavity neck




## Results: Air flow at the cavity neck


S. Gekle et al., PRL 104, 024501 (2010).

## Breaking axial symmetry

Instead of round disc:


## Breaking axial symmetry

## Instead of round disc:


use flowershaped discs with perturbation:

$m=2$


## Experiment for $\boldsymbol{m}=2$

## Top view



## Experiment for $\boldsymbol{m}=\mathbf{2}$

## Top view



## Experiment for $\boldsymbol{m}=\mathbf{2}$

## Top view


cavity shape reverses!

## Basic mechanism: circular cavity

Continuity argument: (cylindrical cavity)

$$
\left.\begin{array}{l}
U R=C \Rightarrow \\
\frac{d U}{d t}=\frac{U^{2}}{R} \\
\left(\dot{U}=-\frac{U \dot{R}}{R}=\frac{U^{2}}{R}\right. \\
\dot{R}=-U
\end{array}\right) .
$$



## Basic mechanism: elliptical cavity

Continuity argument: (elliptical cavity)

$$
\begin{aligned}
& U \mathcal{R}=C \\
& \mathcal{R}=\begin{array}{l}
\text { local radius } \\
\text { of curvature }
\end{array} \\
& \frac{d U}{d t}=\frac{U^{2}}{\mathcal{R}}
\end{aligned}
$$



## Experiments in linear regime

- Linear behavior: $a \ll R$
- Disks with 1\% perturbation
- Water + powdered milk for visualization



## A bit of theory:

Collapsing cavity with small azimuthal perturbation mode $m$ around average $\bar{R}(t)$ :

$$
R(\theta, t)=\bar{R}(t)+a_{m}(t) \cos (m \theta)
$$

corresponds to the flow potential:

$$
\phi(r, t)=Q(t) \log r+d_{m}(t) r^{-m} \cos (m \theta)
$$

The kinematic boundary condition at the cavity wall gives:

$$
\left.\frac{\partial \phi}{\partial r}\right|_{r=\bar{R}(t)}=\frac{\partial R}{\partial t} \Rightarrow Q(t)=\bar{R} \dot{\bar{R}} ; \quad d_{m}(t)=\frac{-\bar{R}^{m+1}}{m}\left[\dot{a}_{m}+a_{m} \frac{\dot{\bar{R}}}{\bar{R}}\right]
$$

Combining with Bernoulli between $\boldsymbol{R}_{\infty}$ and cavity wall $\boldsymbol{R}(\boldsymbol{t})$ (dynamic b.c.):

$$
\rho\left[\frac{\partial \phi}{\partial t}+\frac{1}{2}|\vec{\nabla} \phi|^{2}\right]_{R_{\infty}}^{R(t)}=\left(P_{\infty}-P_{0}\right)+\ngtr
$$

after linearizing in $a_{\mathrm{m}}$ provides the amplitude equation:

$$
\ddot{a}_{m}+\left(\frac{2 \dot{\bar{R}}}{\bar{R}}\right) \dot{a}_{m}+\left(\frac{\ddot{\bar{R}}}{\bar{R}}(1-m)\right) a_{m}=0
$$

L.E. Schmidt, N.C. Keim, W.W. Zhang, and S.R. Nagel, Nature Phys. 5, 343 (2009)

## Inserting 2D Rayleigh scaling

Using the 2D Rayleigh scaling for the average radius $\bar{R}(t)$

$$
\bar{R}(t)=C \sqrt{R_{0} V}\left(t_{c}-t\right)^{1 / 2}
$$

in the linear amplitude equation

$$
\ddot{a}_{m}+\left(\frac{2 \dot{\bar{R}}}{\bar{R}}\right) \dot{a}_{m}+\left(\frac{\ddot{\bar{R}}}{\overline{\bar{R}}}(1-m)\right) a_{m}=0
$$

we find the solution:

$$
a_{m}(t)=a_{m}(0) \cos \left(\frac{1}{2} \sqrt{m-1} \log \left(t_{c}-t\right)+\widetilde{\delta}\right)
$$

or:

$$
a_{m}(\bar{R})=a_{m}(0) \cos \left(\sqrt{m-1} \log \left(\bar{R} / R_{0}\right)+\delta\right)
$$

constant amplitude $a_{\mathrm{m}} \Rightarrow$ $a_{\mathrm{m}} / \overline{\boldsymbol{R}}$ diverges !
frequency chirps !

## Results (linear regime)



$$
R(\theta, t)=\bar{R}(t)+a_{m}(t) \cos (m \theta)
$$

## Results (linear regime)



$$
a_{m}(\bar{R})=a_{m}(0) \cos \left(\sqrt{m-1} \log \left(\bar{R} / R_{0}\right)+\delta\right)
$$

$\bar{R} \approx R_{0}$

## Linear regime, 1\% perturbation






## Linear regime, 1\% perturbation



## Experiments in the non-linear regime

 (perturbations $a$ comparable to disc radius $R_{0}$ )- jet formation



## Experiments in the non-linear regime

 (perturbations $a$ comparable to disc radius $R_{0}$ )- jet formation
- cusp formation ( $m \geq 3$ )

$$
m=5
$$



## Experiments in the non-linear regime

 (perturbations $a$ comparable to disc radius $R_{0}$ )- jet formation
- cusp formation ( $m \geq 3$ )
- subcavities

$$
m=5
$$

A side view ( $m=20,10 \%$ perturbation)


A side view (m=20, 10\% perturbation)


## A side view (m=20, 10\% perturbation)



## Understanding the side view pattern



Also described by theory?

- use axisymmetric BI result $\bar{R}(z, t)$
- in linear amplituo equation

O. Enríquez et al., JFM 701, 40-58 (2012).


## Conclusions

## Void creation and collapse:

- experiment \& BI numerics agree wonderfully
- 2D Rayleigh equation captures essential dynamics
- Airflow becomes supersonic in the neck


## Breaking axisymmetry:

- perfect agreement small oscillations with Schmidt's theory
- large perturbations show cusp, jet \& subcavity formation
- axisymmetric BI result + Schmidt's theory captures pineapple shape cavity


